Ulli Wolff, Masy 2012, background material for CP1.

## 1 Perfect trapezoidal integrator for periodic analytic functions

In CP1 it was experimentally found the approximation

$$T_N(f) = h\left(\frac{1}{2}f_0 + f_1 + f_2 + \dots + f_{N-1} + \frac{1}{2}f_N\right) \approx \frac{1}{2}I_0(\beta)$$
(1)

with<sup>1</sup>

$$h = \frac{\pi}{N}, \quad x_i = ih, \quad f_i \equiv f(x_i), \quad f(x) = \frac{1}{2\pi} e^{\beta \cos(x)}$$

approximates the Bessel function

$$I_0(\beta) = \int_0^\pi \frac{d\varphi}{\pi} e^{\beta \cos(\varphi)} = \int_0^{2\pi} \frac{d\varphi}{2\pi} e^{\beta \cos(\varphi)}$$
(2)

without the usual error of  $O(N^{-2})$ . It has machine precision instead for N = 12 or so. In this note we explain why this is so and estimate the in fact exponential convergence. All this is well-known to the experts, of course.

## 1.1 Euler Mac-Laurin

First we note

$$2T_N = h(f_0 + f_1 + \dots + f_{2N-1}) \tag{3}$$

due to periodicity,  $f_0 = f_{2N}$ . Next we expand the (analytic) integrand

$$\int_{u}^{u+h} dx f(x) = \int_{u}^{u+h} dx \sum_{k=0}^{\infty} \frac{(x-u)^{k}}{k!} f^{(k)}(u) = \sum_{k=0}^{\infty} \frac{h^{k+1}}{(k+1)!} f^{(k)}(u)$$
(4)

and similarly

$$\int_{u}^{u+h} dx f(x) = \sum_{k=0}^{\infty} (-)^{k} \frac{h^{k+1}}{(k+1)!} f^{(k)}(u+h).$$
(5)

By symmetrizing over both variants we get

$$\int_{u}^{u+h} dx f(x) = \frac{1}{2} \sum_{k=0}^{\infty} \frac{h^{k+1}}{(k+1)!} [f^{(k)}(u) + (-)^{k} f^{(k)}(u+h)].$$
(6)

<sup>1.</sup> The normalization  $1/2\pi$  was not there before, but is nicer.

Adding this up for all subintervals

$$\int_{0}^{2\pi} dx f(x) - 2T_{N}(f) = h^{2} \sum_{k=1}^{\infty} \frac{h^{2k-1}}{(2k+1)!} \sum_{n=0}^{2N-1} f^{(2k)}(x_{n})$$

$$= h^{2} \sum_{k=1}^{\infty} \frac{h^{2k-2}}{(2k+1)!} 2T_{N}(f^{(2k)}),$$
(7)

all odd derivatives cancel and we have moved the zero'th order to the left hand side. This is essentially the Euler Mac-Laurin formula for estimating the difference between an integral and the approximating Riemann sum. In particular, the right hand side starts at order  $h^2$ .

For our peridioic function we may now iterate this as follows

$$\underbrace{\int_{0}^{2\pi} dx f''(x)}_{=0} - 2T_N(f'') = h^2 \sum_{k=1}^{\infty} \frac{h^{2k-2}}{(2k+1)!} 2T_N(f^{(2k+2)})$$
(8)

which puts the original error to  $O(h^4)$  with the integral vanishing by periodicity.

By iterating this process we may show that the error vanishes *faster than any* power of  $N^{-1}$ . In the second example we have integrated  $\sin(x)$  from 0 to  $\pi$  and did find  $N^{-2}$  errors. If we want to replace this by an integral over a periodic function we have to take  $|\sin(x)|$  which is not analytic!!!

## **1.2** Estimate of the convergence rate by a contour method

We rewrite  $I_0$  as contour integral

$$I_0(\beta) = \oint \frac{dz}{2\pi i z} g(z), \quad g(z) = e^{\beta(z+z^{-1})/2} = g(z^{-1})$$
(9)

which in the original form is around a unit circle but may be deformed (avoiding to cross the origin). Our approximation may be written

$$T_N = \frac{1}{N} \sum_{\{z_n\}} g(z_n), \quad z_n = e^{2\pi i n/N}.$$
 (10)

Note that here N is comparable to 2N of the previous subsection.

To also write  $T_N$  as a contour integral we note first the identity

$$\prod_{n} (z - z_n) = z^N - 1.$$
 (11)

Both sides must coincide as they are polynomials of degree N with  $z_n$  as zeros and the highest power  $z^N$  with coefficient one. By differentiating the log we derive

$$\frac{1}{N}\sum_{n}\frac{1}{z-z_{n}} = \frac{z^{N-1}}{z^{N}-1}$$
(12)

This may be employed to write (10) as a contour integral with the help of the residue theorem

$$T_N = \left[ \oint_{>} - \oint_{<} \right] \frac{dz}{2\pi i z} \frac{z^N}{z^N - 1} g(z)$$
(13)

with one integral along a circle with a radius larger than unity and with a smaller one such that effectively each pole at  $z_n$  is encircled in a positive sense and  $g(z_n)$  is the residue. A transformation  $z \to z^{-1}$  reflects the larger into the smaller circle yielding

$$T_N = \oint_{<} \frac{dz}{2\pi i z} \frac{1+z^N}{1-z^N} g(z)$$
(14)

and then

$$\Delta_N = T_N - I_0 = 2 \oint_{<} \frac{dz}{2\pi i z} \frac{z^N}{1 - z^N} g(z).$$
(15)

We want to compute this integral by a saddlepoint approximation and anticipate a saddle point of order 1/N. To have a clear-cut power counting we change variables to  $\alpha = zN$  and write

$$\Delta_N = 2N^{-N} \oint \frac{d\alpha}{2\pi i \alpha} e^{NS(\alpha)}$$
(16)

with

$$S(\alpha) = \ln(\alpha) + (\beta/2)\alpha^{-1}.$$
(17)

where we have dropped terms that correspond to higher order 1/N corrections. The saddle point is at

$$\overline{\alpha} = \beta/2 \tag{18}$$

and setting (for real  $\beta$ )  $\alpha = \overline{\alpha} + i\eta$  for the path of steepest descent thru  $\overline{\alpha}$  we find

$$\Delta_N = \delta_N (1 + \mathcal{O}(N^{-1})) \tag{19}$$

with

$$\delta_N = 2N^{-N} \mathrm{e}^{NS(\overline{\alpha})} \frac{1}{\overline{\alpha}} \int_{-\infty}^{\infty} \frac{d\eta}{2\pi} \mathrm{e}^{-\frac{1}{2}NS''(\overline{\alpha})\eta^2} = \frac{2N^{-N}}{\overline{\alpha}\sqrt{2\pi NS''(\overline{\alpha})}} \mathrm{e}^{NS(\overline{\alpha})}.$$
 (20)

Inserting numbers we find

$$\delta_N = \sqrt{\frac{2}{\pi N}} \left(\frac{2N}{\mathrm{e}\beta}\right)^{-N}.$$
(21)

In a brief test under matlab this gives an excellent description of the error. An impression is given in Figure 1. Note that for N > 10, the approximation is so excellent, that the deviation is covered by roundoff noise, i.a.  $T_N$  is practically at machine precision.



**Figure 1.** Relative error of  $\Delta_N$  for b=1 and N=4...10.