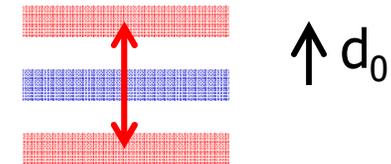


Modell für den Hertzschen Dipol

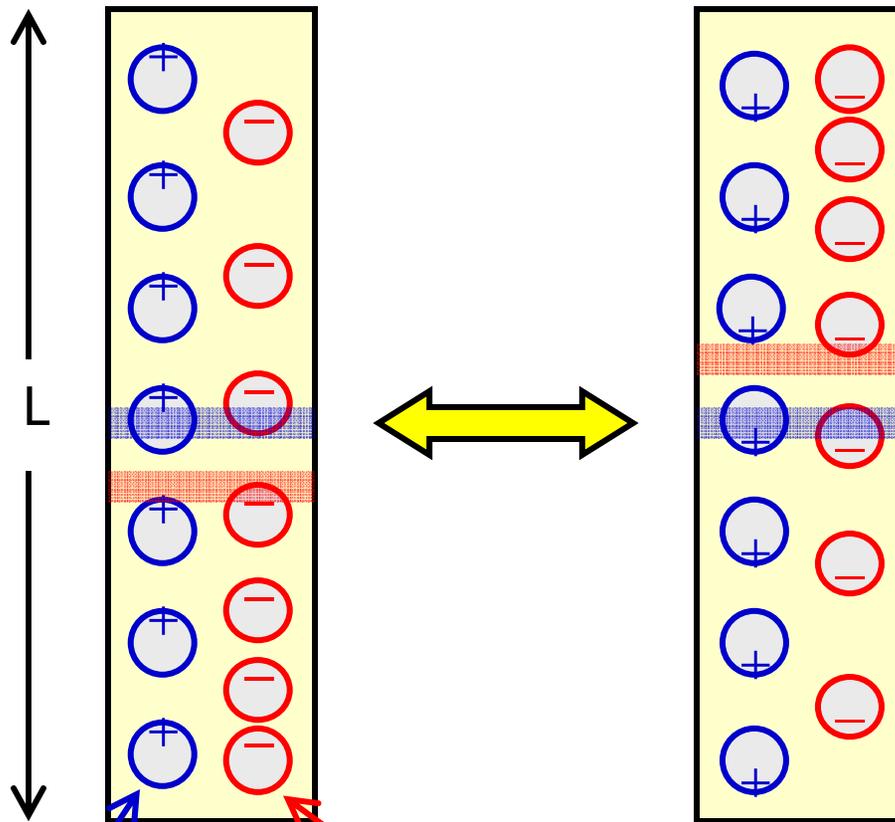


Dipolnäherung \Leftrightarrow Bewegung der Ladungsschwerpunkte

$$p(t) = Qd_0 e^{i\omega t}$$



d_0 ist sehr viel kleiner als L



feste Ionenrümpfe
(Gesamtladung Q)

frei bewegliche Elektronen
(Gesamtladung $-Q$)

Abstrahlung der Felder



$$\text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$



$$\text{rot } \vec{B} = \mu_0 \vec{j} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$$

$$\mu_0 \vec{j}$$

⇒ Dynamik des Stromflusses
(Quasistatik) ⇒

$$\text{Nahfelder: } \vec{E}, \vec{B} \propto \frac{1}{r^3}$$

E- und B-Feld 90° phasenverschoben

$$\frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$$

⇒ Eigendynamik der Felder ⇒

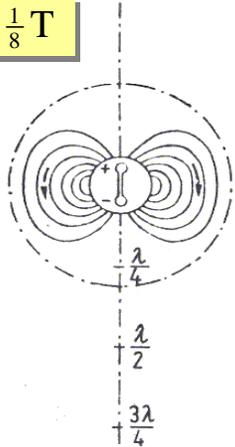
$$\text{Fernfelder: } \vec{E}, \vec{B} \propto \frac{1}{r}$$

E- und B-Feld phasengleich
dominant für $r \gg d_0$

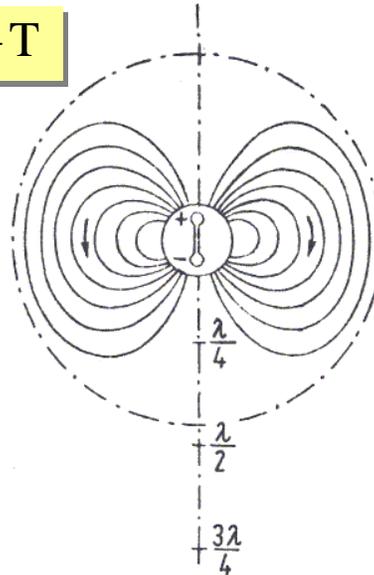
Zeitentwicklung des E-Feldes



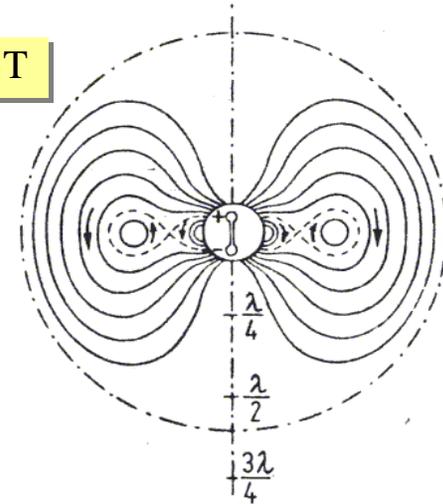
$$t = \frac{1}{8}T$$



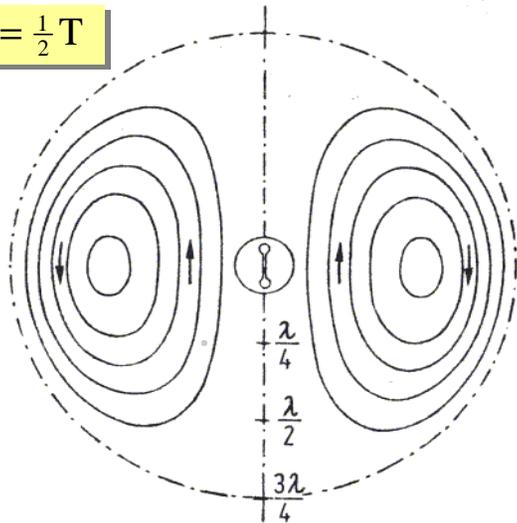
$$t = \frac{1}{4}T$$



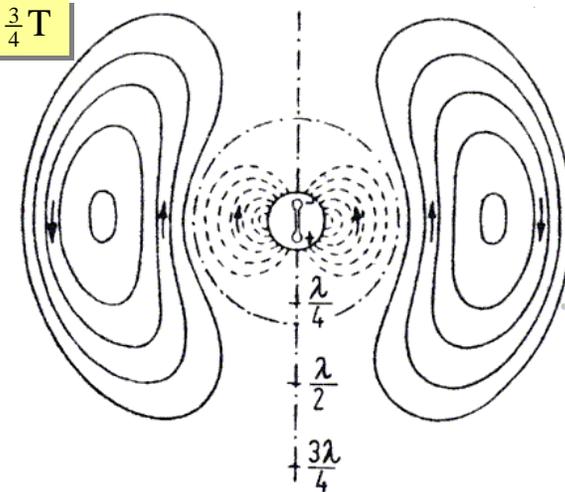
$$t = \frac{3}{8}T$$



$$t = \frac{1}{2}T$$



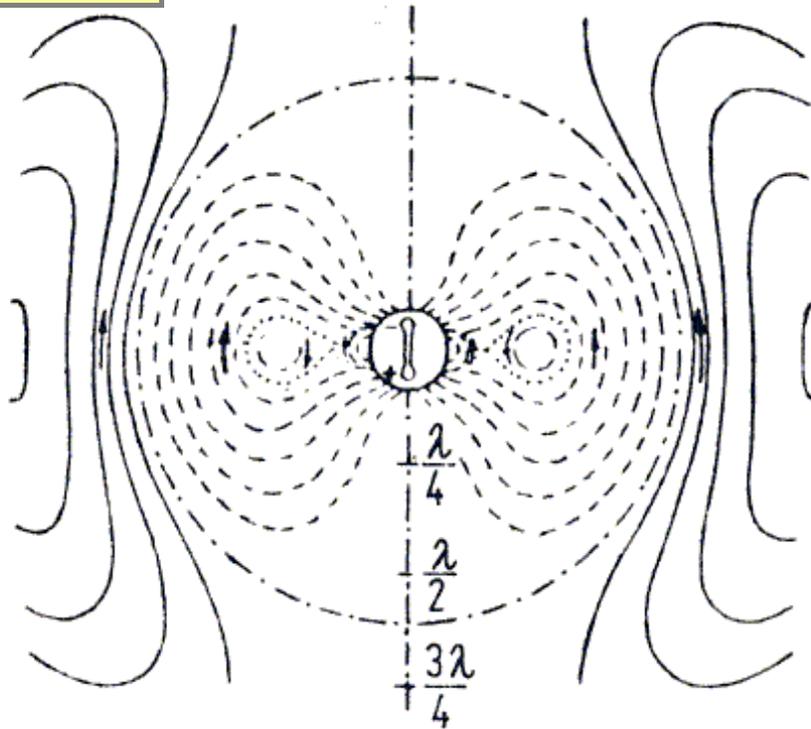
$$t = \frac{3}{4}T$$



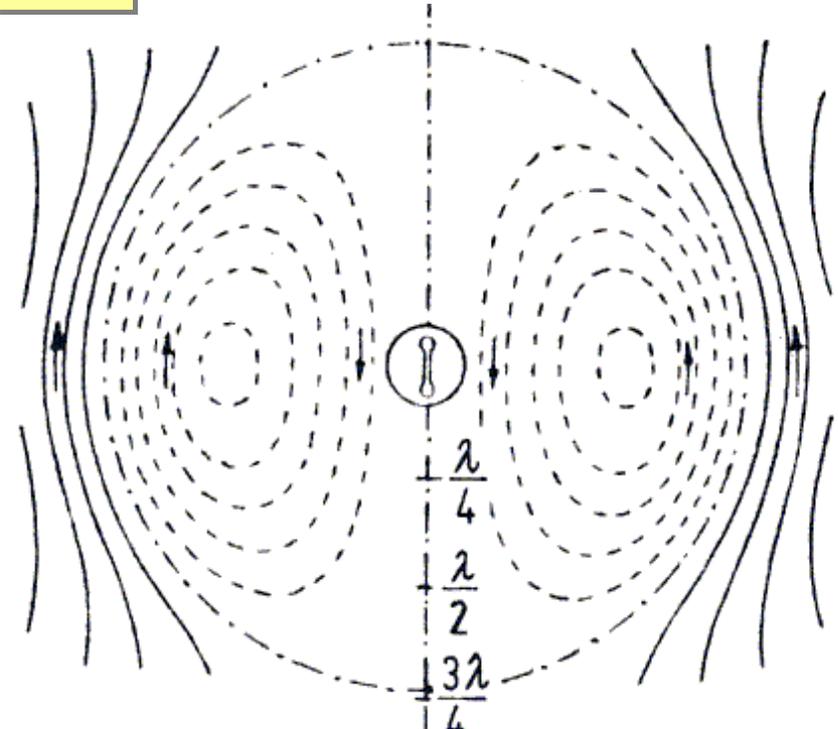
Zeitentwicklung des E-Feldes



$$t = \frac{7}{8}T$$



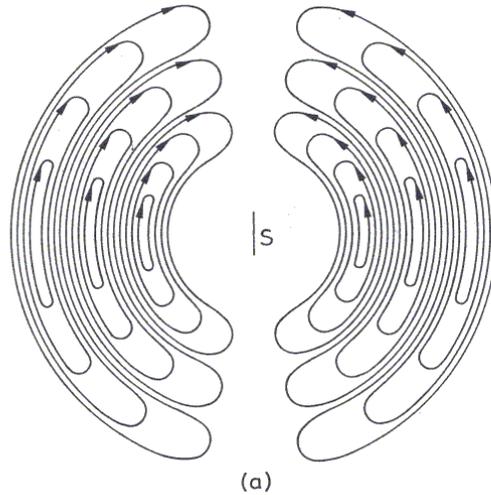
$$t = T$$



E- und B-Fernfelder

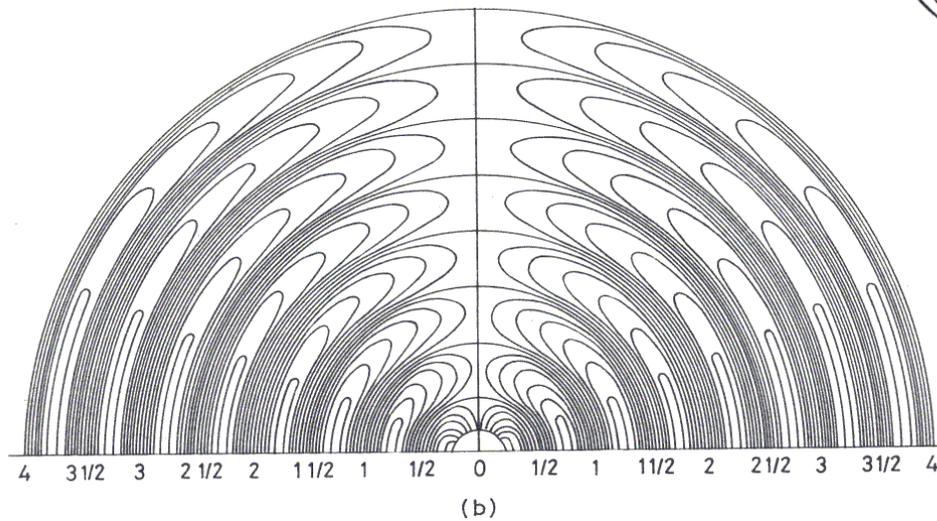
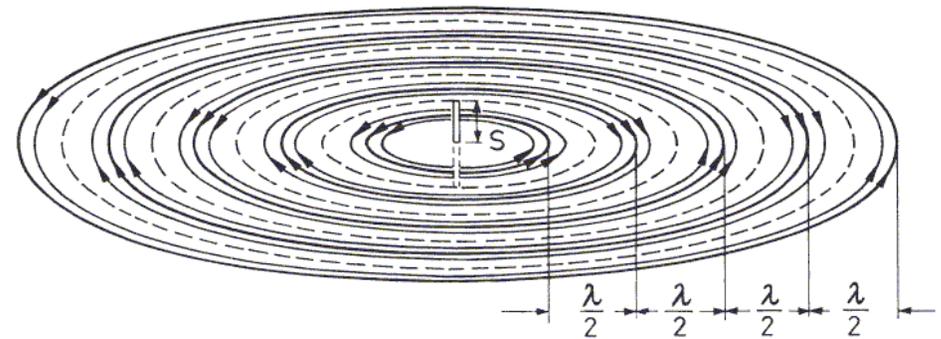


\vec{E}



(a)

\vec{B}

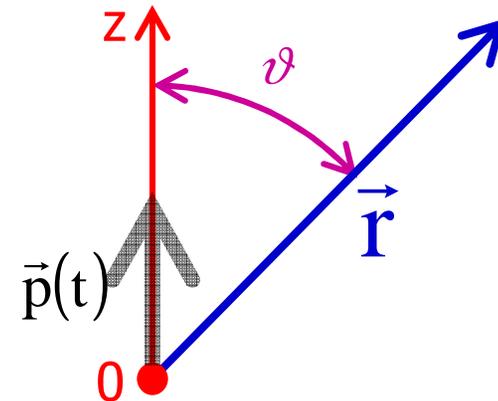


(b)

Quantitative Eigenschaften der Fernfelder



- \vec{B} -Feld konzentrisch um Dipolachse
- $|\vec{B}| \propto \sin \vartheta$ (max. in Äquatorialebene)
- $\vec{E} \perp \vec{B}$, $\vec{E} \perp \vec{r}$, $|\vec{E}| \propto \sin \vartheta$
(max. in Äquatorialebene)



- mittlere Energiestromdichte

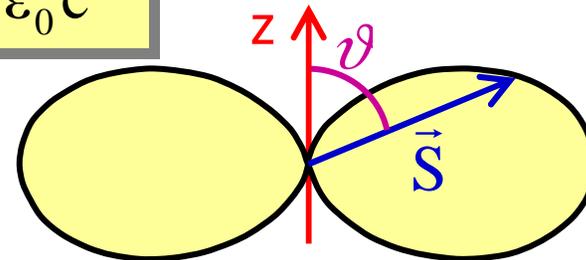
$$\bar{S}(\vartheta, \varphi) = \bar{S}(\vartheta) = \frac{p_0^2 \omega^4}{16 \pi^2 \epsilon_0 c^3} \cdot \frac{\sin^2 \vartheta}{r^2}$$

- mittlere abgestrahlte Leistung

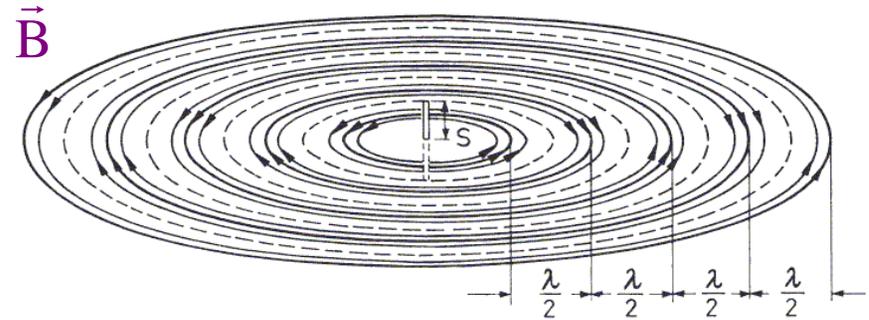
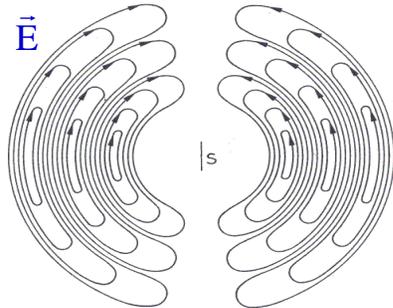
$$\bar{P} = \frac{p_0^2 \omega^4}{6 \pi \epsilon_0 c^3}$$

- Abstrahlcharakteristik

Abstrahlung $\propto \omega^4$
senkrecht zur Dipolachse



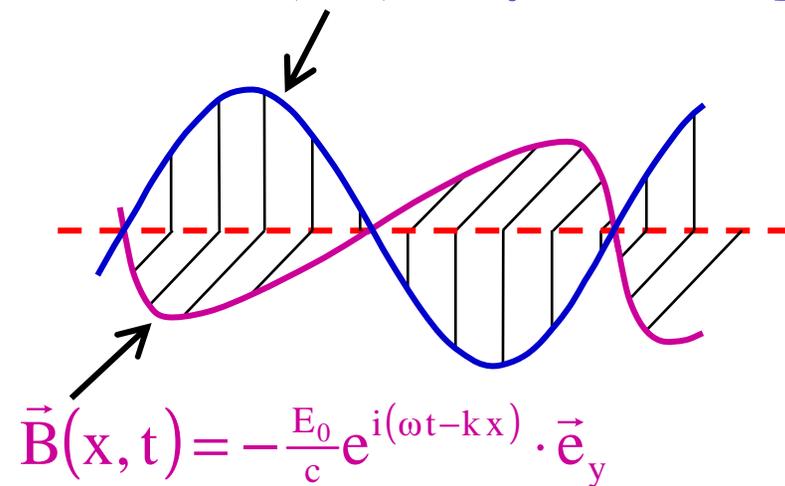
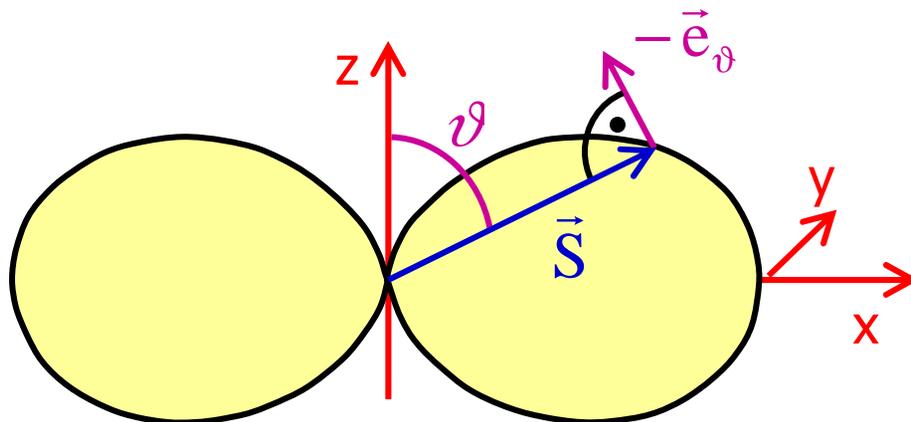
Strahlung in großem Abstand ($r \gg d$)



Krümmung der Phasenflächen zu vernachlässigen

⇒ Ebene Wellen, Polarisation $\parallel \vec{e}_\vartheta$

Strahlung senkrecht zur Antenne (x-Richtung): $\vec{E}(x, t) = E_0 e^{i(\omega t - kx)} \cdot \vec{e}_z$



$$\vec{B}(x, t) = -\frac{E_0}{c} e^{i(\omega t - kx)} \cdot \vec{e}_y$$