

# $t\bar{t}$ Production: $\sigma_{t\bar{t}}$ and $m_t$ @ ATLAS/LHC

Peter Galler

Humboldt-Universität zu Berlin, Institut für Physik

25.05.2012

# Theory Introduction

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$



# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

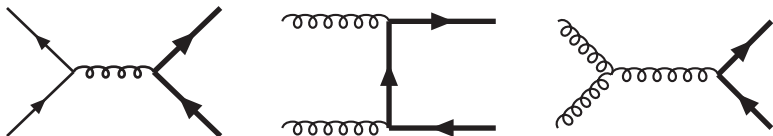
- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Top Quark Decay

- top quark's predominant decay channel  $t \rightarrow W^+ b$
- $\Gamma(t \rightarrow Wb) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \sim 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$   
 $m_t \gg m_b, m_W, |V_{tb}| \sim 1$
- $\Gamma(t \rightarrow sW) \sim |V_{ts}|^2 \Rightarrow B(t \rightarrow sW) \sim 0.2\%$  since  $|V_{ts}|$  small
- $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma_t} \sim 1$
- Lifetime  $\tau = \frac{1}{\Gamma} \sim \frac{1}{G_F m_t^3} \Rightarrow \tau \sim 10^{-25} \text{ s}$   
 $c\tau \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm}$ , hadron size  $\sim 1 \text{ fm} \Rightarrow$  **No Hadronization**
- subsequent decay of W:
  - leptonic:  $W \rightarrow l\nu$
  - hadronic:  $W \rightarrow qq'$
- $\Gamma(t \rightarrow be\nu) = \Gamma(t \rightarrow Wb) \frac{\Gamma(W \rightarrow e\nu)}{\Gamma_W}$  in the narrow-width approx.
- unconventional decays can alter  $B(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wb) + \Gamma(t \rightarrow X)}$

# Heavy Quark Production (LO)

Production by quark annihilation and gluon fusion at leading order:



Differential Cross Section:

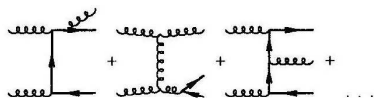
$$\frac{d\sigma}{dy_3 dy_4 d^2 p_T} = \frac{1}{64\pi^2 m_T^4 (1 + \cosh(\Delta y))^2} \times \sum_{ij} x_1 f_i(x_1, \mu^2) x_2 f_j(x_2, \mu^2) \overline{|\mathcal{M}_{ij}|^2}$$

In the limit  $\Delta y \gg 1$

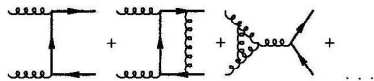
$$\overline{|\mathcal{M}_{qq}|^2} \sim \text{const}, \quad \overline{|\mathcal{M}_{gg}|^2} \sim e^{\Delta y}$$

# Heavy Quark Production (NLO)

- Perturbation Expansion in the coupling constant  $\alpha_s$
- LO  $\sim \alpha_s^2$
- NLO  $\sim \alpha_s^3$  include virtual and real corrections
- reduction of unphysical  $\mu$ -dependence  $\mu^2 \frac{d}{d\mu^2} \sigma = \mathcal{O}(\alpha_s^4)$



Real emission diagrams



Virtual emission diagrams

Examples of higher-order corrections to heavy quark production.

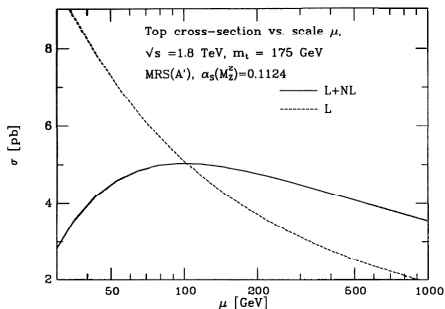


Fig. 10.10. Scale dependence of the top quark cross section in leading and next-to-leading order of perturbation theory.

[figures: QCD & Collider Physics, Ellis et al.]

# Experimental Analysis

## PART I

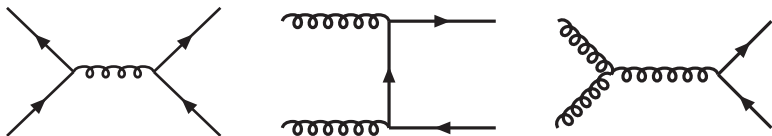
Cross Section Measurement for the Production of  $t\bar{t}$ -Pairs

[ATLAS-CONF-2012-031]

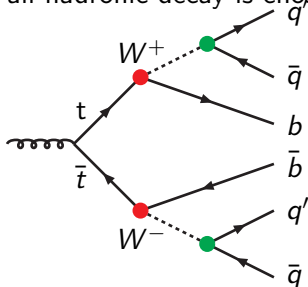


# Introduction

Production by quark annihilation and gluon fusion:



In the detector only decay products can be measured. For this analysis the all-hadronic decay is chosen:



● BR  $\sim$  100%

$\Rightarrow$  total BR  $\sim$  46%

● BR  $\sim$  68%

but **high background!**

# Event Characteristics

physics observables are **jets**:

- b-jets originating from top decay
- light quark jets originating from W decay
- no isolated lepton with high  $p_T$
- no missing energy  $E_{miss}$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20$  GeV
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20$  GeV
- $N_{jet} \geq 5$ ,  $p_T > 55$  GeV and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30$  GeV and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55$  GeV,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20$  GeV
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20$  GeV
- $N_{jet} \geq 5$ ,  $p_T > 55$  GeV and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30$  GeV and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55$  GeV,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|\text{JVF}| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$



# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|JVF| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

# Event Selection

- at least 1 primary vertex with  $\geq 5$  tracks
- no isolated lepton with  $p_T > 20 \text{ GeV}$ 
  - only hadronic decays
- no fake jets
  - reject unphysical jets from fake energy deposits in calorimeters
- no jets with  $|\text{JVF}| \leq 0.75$  and  $p_T > 20 \text{ GeV}$
- $N_{jet} \geq 5$ ,  $p_T > 55 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one more jet:  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.5$ 
  - all-hadronic decay needs at least six jets
- at least 2 b-tagged jets:  $p_T > 55 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - top/antitop decay into  $b/\bar{b}$  quark
- $S_t = E_T^{miss} / (0.5\sqrt{\text{GeV}} \times \sqrt{H_T}) < 6$ 
  - if significance of  $E_T^{miss}$  is high then veto
- angular distance between two b-jets:  $\Delta R(b, \bar{b}) > 1.2$ 
  - reject b-jets from gluon splitting
- angular distance between any two jets:  $\Delta R > 0.6$ 
  - increase trigger efficiency at high  $p_T$

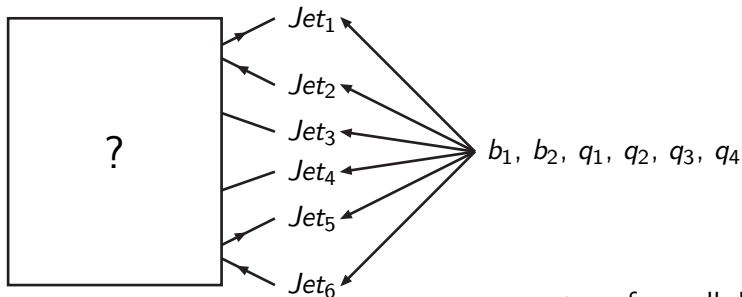
# Likelihood Method

- Idea:
  - determine jet energies  $E_{jet_i}$  by maximum likelihood method
  - use fitted  $E_{jet_i}$  to compute  $m_t$ -distribution
  - perform unbinned likelihood fit to the  $m_t$ -distribution to obtain  $\sigma_{t\bar{t}}$
- Likelihood (constrains on  $m_W$ ,  $\Gamma_W$  and  $m_{top}$ )

$$\begin{aligned}
 L_{kin} = & BW(m_{q_1 q_2} | m_W, \Gamma_W) \cdot BW(m_{q_3 q_4} | m_W, \Gamma_W) \cdot \\
 & BW(m_{q_1 q_2 b_1} | m_{top}^{reco}, \Gamma_{top}) \cdot BW(m_{q_3 q_4 b_2} | m_{top}^{reco}, \Gamma_{top}) \cdot \\
 & W(\hat{E}_{jet_1} | E_{b_1}) \cdot W(\hat{E}_{jet_2} | E_{b_2}) \cdot W(\hat{E}_{jet_3} | E_{q_1}) \cdot \\
 & W(\hat{E}_{jet_4} | E_{q_2}) \cdot W(\hat{E}_{jet_5} | E_{q_3}) \cdot W(\hat{E}_{jet_6} | E_{q_4}) \cdot \\
 & \prod_{i=1}^2 \left\{ \begin{array}{ll} \epsilon; & b_i \text{ b-tagged} \\ 1 - \epsilon; & b_i \text{ not b-tagged} \end{array} \right\} \cdot \\
 & \prod_{i=1}^4 \left\{ \begin{array}{ll} \frac{1}{R}; & q_i \text{ b-tagged} \\ 1 - \frac{1}{R}; & q_i \text{ not b-tagged} \end{array} \right\}
 \end{aligned}$$

## Likelihood Method II

A priori association of jets with quarks is impossible

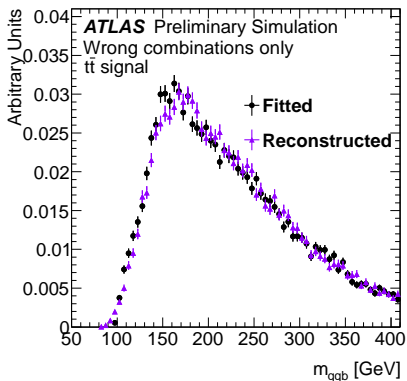
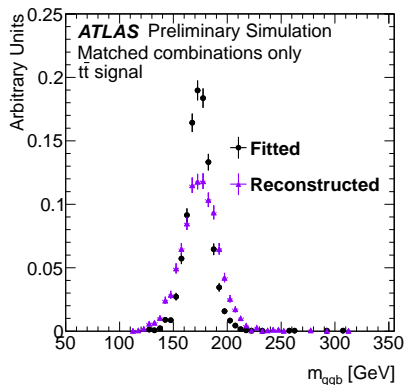


$\Rightarrow$  perform all distinguishable permutations for each event and select the one with the highest likelihood  
 = event probability  $P_{event}$

$$P_{event} = \frac{\text{Min}_{Perm}(-\ln L_{kin})}{\sum_{\{jet_i\} \in \mathcal{P}(1, \dots, 6)} (-\ln L_{kin}(\{jet_i\}))}$$

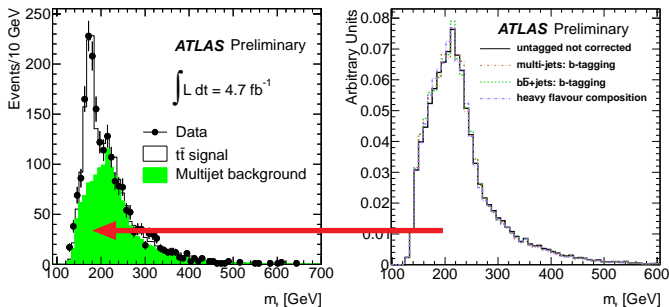
# Top Quark Mass Distribution

By the likelihood method one obtains the most likely jet combination which can be used to calculate  $m_t = m_{qqb}$



# Background Modelling

- Construct  $m_t$ -distribution from MC sample with eventselection as before but **without b-jet requirement**  
→ resulting distribution contains  $t\bar{t}$  all-hadronic contribution of estimated 4.6%
- take the shape of the distribution (+corrections) to be the probability density function of the background
- construct a Likelihood function for the background and fit the the data



# Cross Section Measurement

1 perform an additional event selection

- reconstructed  $m_t > 125 \text{ GeV}$
- $6 \leq N_{jet} \leq 10$
- $P_{event} > 0.8$
- $\text{Min}_{Perm} \chi^2 < 30$

$$\chi^2 = \frac{(m_{j_1, j_2} - m_w)^2}{\sigma_W^2} + \frac{(m_{j_1, j_2, b_1} - m_t)^2}{\sigma_t^2} + \frac{(m_{j_3, j_4} - m_W)^2}{\sigma_W^2} + \frac{(m_{j_3, j_4, b_2} - m_t)^2}{\sigma_t^2}$$

⇒ correct permutations rise up to 36%

- 2 after Likelihood fit a signal fraction of  $f_{t\bar{t}}^{sig} = (31.4 \pm 2.3)\%$  is obtained from final data sample with  $N_{t\bar{t}} = 2118$  events
- 3  $t\bar{t}$  selection efficiency derived from MC sim.:  $\epsilon = (0.086 \pm 0.003)\%$
- 4 Luminosity  $\mathcal{L} = (4.7 \pm 0.2) \text{ fb}^{-1}$
- 5  $\sigma_{t\bar{t}} = \frac{f_{t\bar{t}}^{sig} N_{t\bar{t}}}{\epsilon \mathcal{L}} = (168 \pm 12) \text{ pb}$



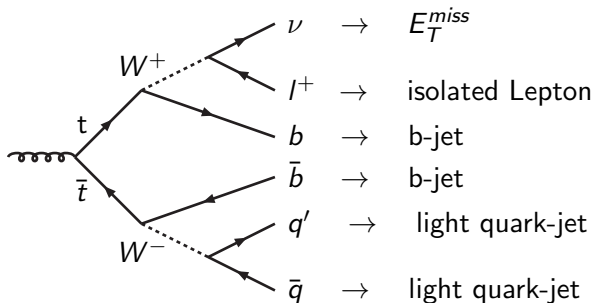
# PART II

## Measurement of $m_t$ from $t\bar{t}$ -Production

[ATLAS arXiv:1203.5755v1]

# Event Topology

Example for the event topology for  $e$ +jets and  $\mu$ +jets decay of the  $t\bar{t}$ -Pair



# Event Selection

- single  $e$  ( $E_T > 25 \text{ GeV}$ ) or single  $\mu$  ( $p_T > 20 \text{ GeV}$ )
- $\mu$ +jets:  $E_T^{\text{miss}} > 20 \text{ GeV}$  and  $E_T^{\text{miss}} + m_W^T > 60 \text{ GeV}$
- $e$ +jets:  $E_T^{\text{miss}} > 35 \text{ GeV}$  and  $m_W^T > 25 \text{ GeV}$
- $N_{\text{jets}} \geq 4$  with  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one jet is a b-jet

# Event Selection

- single  $e$  ( $E_T > 25 \text{ GeV}$ ) or single  $\mu$  ( $p_T > 20 \text{ GeV}$ )
- $\mu$ +jets:  $E_T^{\text{miss}} > 20 \text{ GeV}$  and  $E_T^{\text{miss}} + m_W^T > 60 \text{ GeV}$
- $e$ +jets:  $E_T^{\text{miss}} > 35 \text{ GeV}$  and  $m_W^T > 25 \text{ GeV}$
- $N_{\text{jets}} \geq 4$  with  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one jet is a b-jet

# Event Selection

- single  $e$  ( $E_T > 25 \text{ GeV}$ ) or single  $\mu$  ( $p_T > 20 \text{ GeV}$ )
- $\mu$ +jets:  $E_T^{miss} > 20 \text{ GeV}$  and  $E_T^{miss} + m_W^T > 60 \text{ GeV}$
- $e$ +jets:  $E_T^{miss} > 35 \text{ GeV}$  and  $m_W^T > 25 \text{ GeV}$
- $N_{jets} \geq 4$  with  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one jet is a b-jet

# Event Selection

- single  $e$  ( $E_T > 25 \text{ GeV}$ ) or single  $\mu$  ( $p_T > 20 \text{ GeV}$ )
- $\mu$ +jets:  $E_T^{miss} > 20 \text{ GeV}$  and  $E_T^{miss} + m_W^T > 60 \text{ GeV}$
- $e$ +jets:  $E_T^{miss} > 35 \text{ GeV}$  and  $m_W^T > 25 \text{ GeV}$
- $N_{jets} \geq 4$  with  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one jet is a b-jet

# Event Selection

- single  $e$  ( $E_T > 25 \text{ GeV}$ ) or single  $\mu$  ( $p_T > 20 \text{ GeV}$ )
- $\mu$ +jets:  $E_T^{miss} > 20 \text{ GeV}$  and  $E_T^{miss} + m_W^T > 60 \text{ GeV}$
- $e$ +jets:  $E_T^{miss} > 35 \text{ GeV}$  and  $m_W^T > 25 \text{ GeV}$
- $N_{jets} \geq 4$  with  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$
- at least one jet is a b-jet

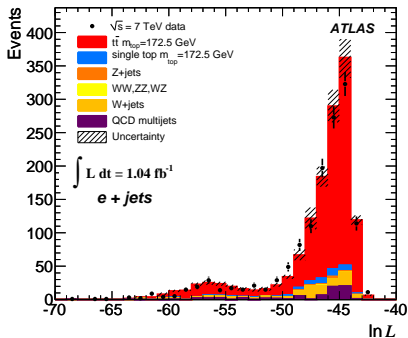
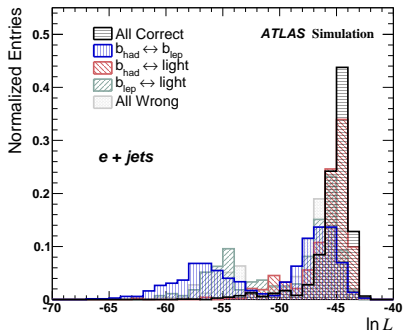
# 1d-Analysis



# Likelihood

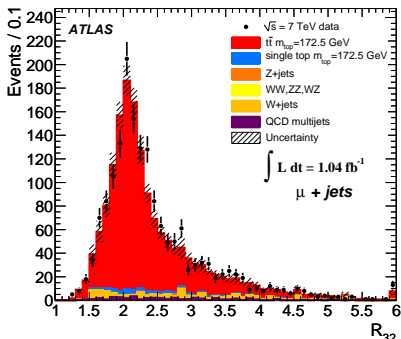
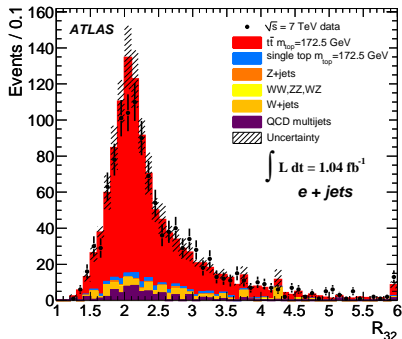
Use likelihood to determine the most likely jet permutation:

$$L = (\text{Breit-Wigner}) \cdot (\text{Transfer Functions}) \cdot (\text{bTag})$$



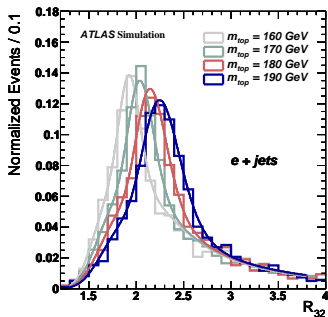
# Observable

- The used observable is  $R_{32} \equiv \frac{m_{top}^{reco}}{m_W^{reco}}$
- $R_{32} = \frac{m_{top}^{reco,like}}{m_W^{reco,like}} \Rightarrow R_{32}$  is dependent on JES (Jet Energy Scale)
- take most likely jet permutation from the likelihood fit and use these jets to reconstruct the masses  $m_{top}^{reco}$  and  $m_W^{reco}$   
 $\Rightarrow R_{32}$  less dependent on JES



# $R_{32}$ -Templates

- use MC data for different  $m_t = 160, 170, 172.5, 175, 180, 190$  to create  $R_{32}$  distributions  $\rightarrow$  templates
- parametrize the templates by functions  $f(R_{32}, \vec{x}_m)$

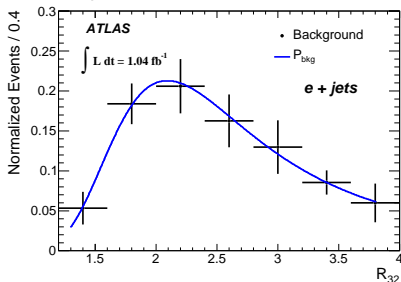


$$\chi^2 = \sum_{m=m_1}^{m_6} \sum_{i=1}^{N_{bins}} \frac{\left[ T((R_{32})_i, m) - f((R_{32})_i, \vec{x}_m) \right]^2}{\sigma_m^2}$$

- perform simultaneous  $\chi^2$ -fit to all templates to obtain a continuous function of  $m_{top}$  and  $R_{32}$  which can be interpreted as signal probability function  $P_{sig}(R_{32} | m_{top})$

# Background and Likelihood Fit for $m_{top}$

- get  $m_{top}$ -independent background  $R_{32}$  distribution from MC



- interpret this as  $P_{bkg}(R_{32})$
- perform likelihood fit with  $n_{bkg}$  and  $m_{top}$  as fit parameters

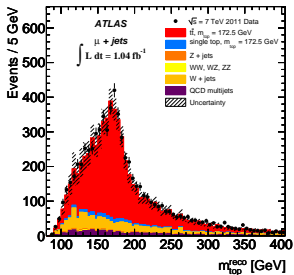
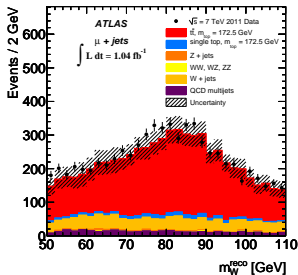
$$\mathcal{L}(R_{32}|m_{top}) = \mathcal{L}_{shape}(R_{32}|m_{top}) \times \mathcal{L}_{bkg}(R_{32}), \quad \mathcal{L}_{shape}(R_{32}|m_{top}) = \prod_{i=1}^{N_{bins}} \frac{\lambda_i^{N_i}}{N_i!} \cdot e^{-\lambda_i}$$

$$\mathcal{L}_{bkg} = \exp \left\{ - \frac{(n_{bkg} - n_{bkg}^{pred})^2}{2\sigma_{n_{bkg}^{pred}}^2} \right\}, \quad \lambda_i = (N - n_{bkg}) \cdot P_{sig}(R_{32}|m_{top})_i + n_{bkg} \cdot P_{bkg}(R_{32})_i$$

## 2d Analysis

# $m_t$ and $m_W$ Distributions

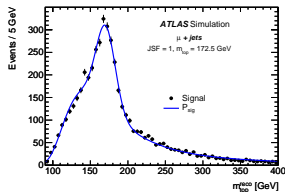
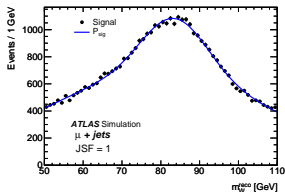
- in this analysis  $m_{top}$  and JSF (Jet Energy Scale Factor) are fitted
- t-quark candidate from hadronic decay is reconstructed from the combination of any b-jet and two light jets which has the highest  $p_T$
- for the light jet pair the partron scale factors  $\alpha_i$  are determined by a  $\chi^2$  fit: 
$$\chi^2 = \sum_{i=1}^2 \left[ \frac{E_{jet_i}(1-\alpha_i)}{\sigma(E_{jet_i})} \right]^2 + \left[ \frac{M_{jet, jet}(\alpha_1, \alpha_2) - m_W}{\Gamma_W} \right]^2$$
- $m_W^{reco}$  is reconstructed without partron rescaling
- $m_{top}^{reco}$  is reconstructed using  $\alpha_i$



# Constructing Templates

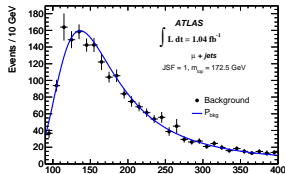
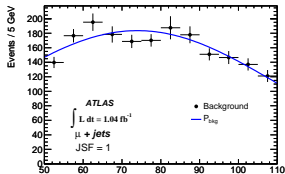
- templates for  $m_{top}^{reco}$ -distribution for values of  $m_t \in [160, 190]$  GeV and JSF  $\in [0.9, 1.1]$
- templates for  $m_W^{reco}$ -distribution for values of JSF  $\in [0.9, 1.1]$
- parametrize the templates for signal and background and interpret them as probability density functions.

$P_W^{sig}$



$P_{top}^{sig}$

$P_W^{bkg}$



$P_{top}^{bkg}$

# Likelihood Fit for JSF and $m_{top}$

Unbinned likelihood fit to the data for all events,  $i = 1, \dots, N$

$$\begin{aligned} \mathcal{L}_{shape}(m_W^{reco}, m_{top}^{reco} | m_{top}, JSF, n_{bkg}) \\ = \prod_{i=1}^N P_{top}(m_{top}^{reco} | m_{top}, JSF, n_{bkg})_i \times P_W(m_W^{reco} | JSF, n_{bkg})_i \end{aligned}$$

with

$$P_{top,i} = (N - n_{bkg}) \cdot P_{top}^{sig}(m_{top}^{reco} | m_{top}, JSF)_i + n_{bkg} \cdot P_{top}^{bkg}(m_{top}^{reco} | m_{top}, JSF)_i$$

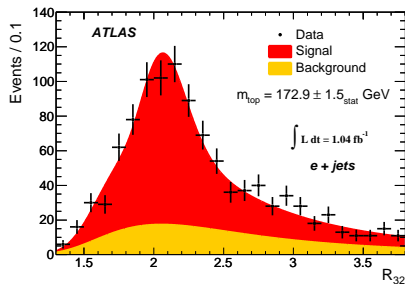
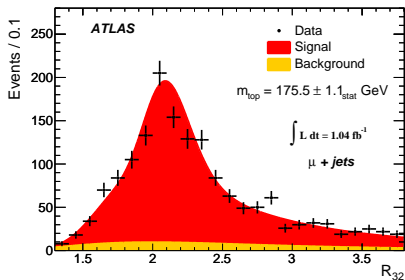
$$P_{W,i} = (N - n_{bkg}) \cdot P_W^{sig}(m_W^{reco} | JSF)_i + n_{bkg} \cdot P_W^{bkg}(m_W^{reco} | JSF)_i$$



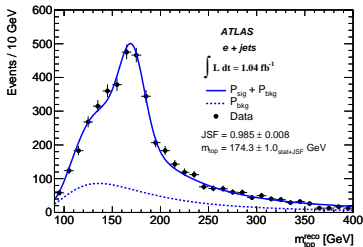
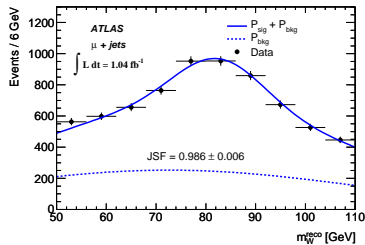
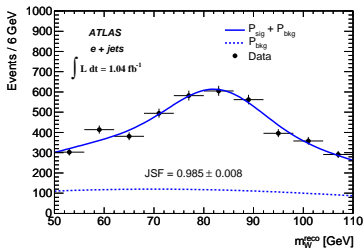
# Numerical Results

- 1d
  - e+jets :  $m_{top} = 172.9 \pm 1.5_{stat} \pm 2.5_{syst} \text{ GeV}$
  - $\mu$ +jets :  $m_{top} = 175.5 \pm 1.1_{stat} \pm 2.6_{syst} \text{ GeV}$
- 2d
  - e+jets :  $m_{top} = 174.3 \pm 0.8_{stat} \pm 2.3_{syst} \text{ GeV}$
  - $\mu$ +jets :  $m_{top} = 175.0 \pm 0.7_{stat} \pm 2.6_{syst} \text{ GeV}$
- e- and  $\mu$ -channel combined
  - 1d:  $m_{top} = 174.4 \pm 0.9_{stat} \pm 2.5_{syst} \text{ GeV}$
  - 2d:  $m_{top} = 174.5 \pm 0.6_{stat} \pm 2.3_{syst} \text{ GeV}$

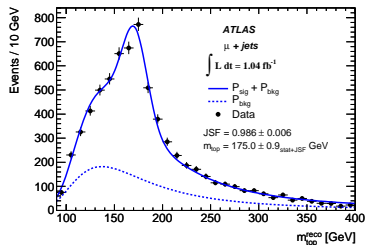
# Graphical Results: 1d-Analysis

 $e + \text{jets}$  $\mu + \text{jets}$

# Graphical Results: 2d-Analysis



e+jets



$\mu$ +jets

# Topquark Mass Comparison

