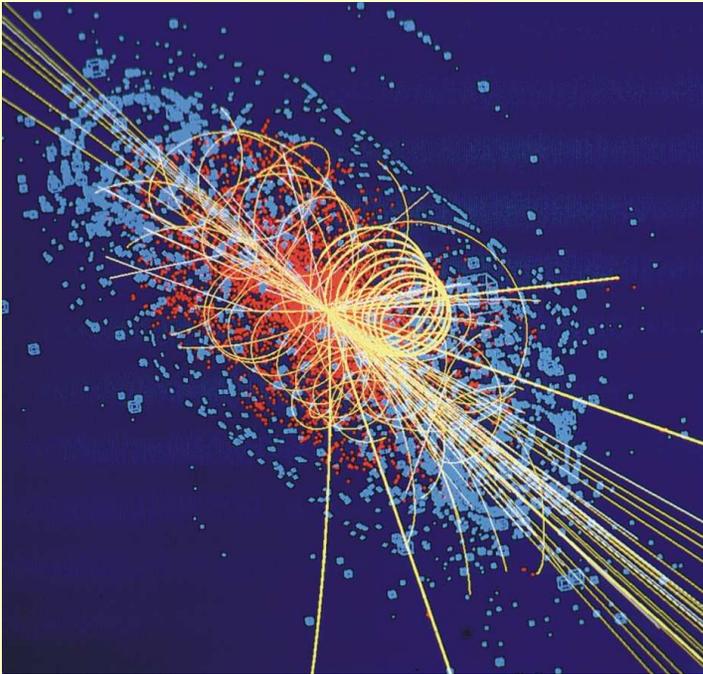


Physics at Hadron Colliders

Part 2



Standard Model Physics

Jet calibration

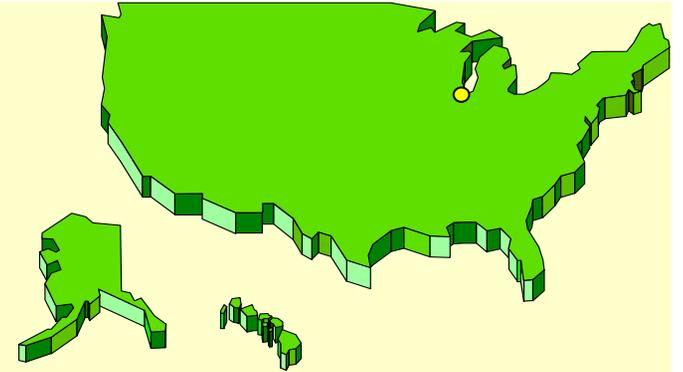
Test of Quantum Chromodynamics

- Jet production
- W/Z production
- Production of Top quarks

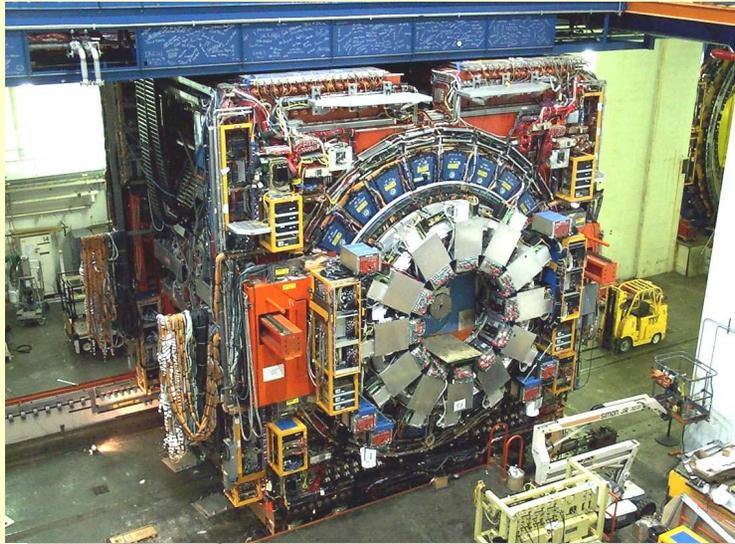
Electroweak measurements

- W mass
- Top-quark mass and other properties
- Single top production

Back to the Tevatron



The CDF experiment



**12 countries, 59 institutions
706 physicists**

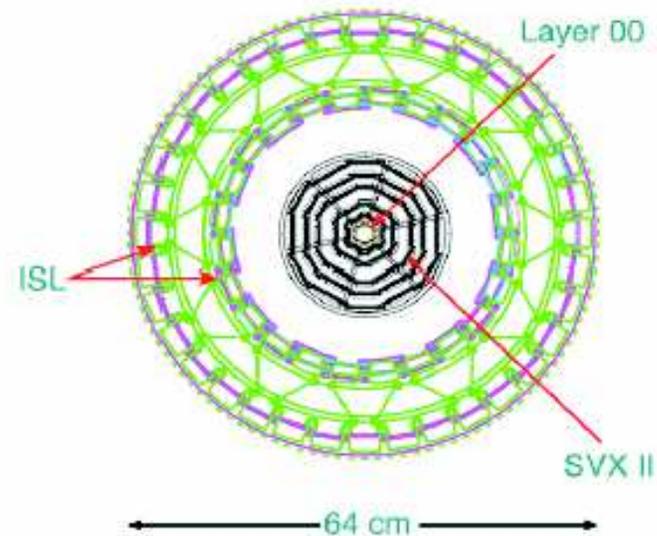
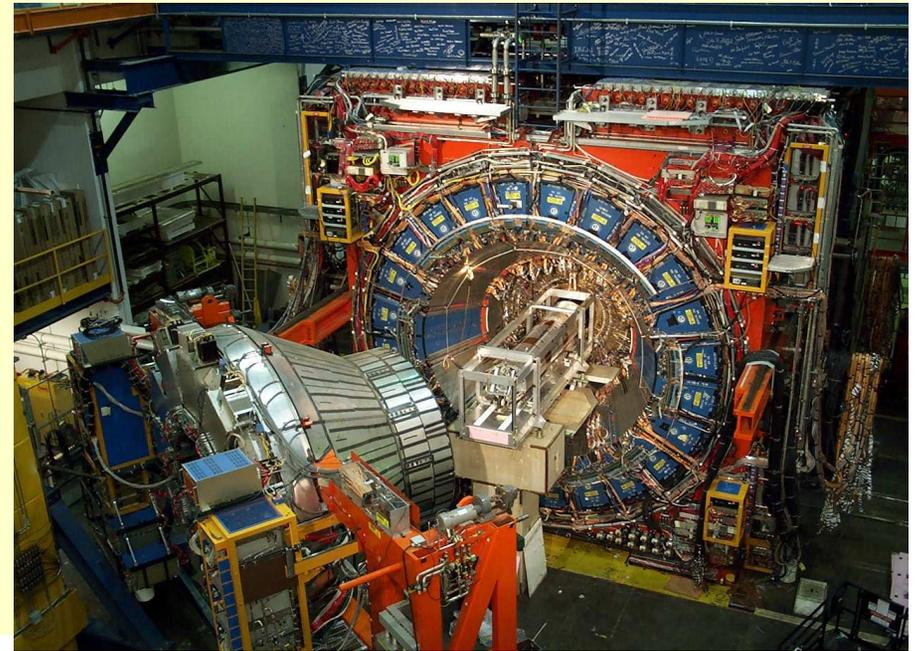
The DØ collaboration



**19 countries, 83 institutions
664 physicists**

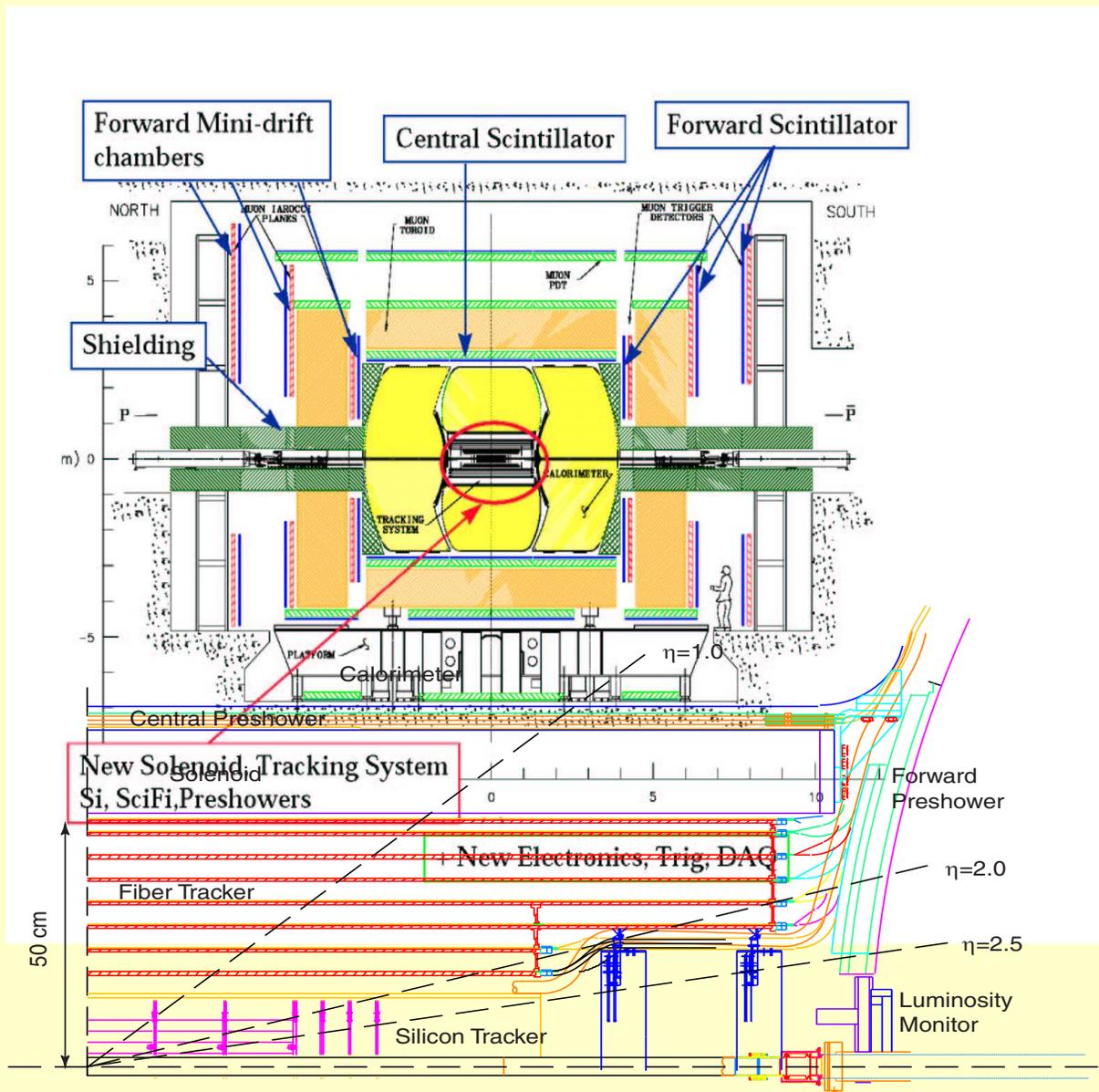
The CDF detector in Run II

- Core detector operates since 1985:
 - Central Calorimeters
 - Central muon chambers
- Major upgrades for Run II:
 - Drift chamber (central tracker)
 - **Silicon tracking detector:**
SVX, ISL, Layer 00
 - 8 layers
 - 700k readout channels
 - 6 m²
 - material: 15% X₀
 - Forward calorimeters
 - Forward muon system
 - Time-of-flight system
 - Trigger and DAQ
 - Front-end electronics





The DØ Run II Detector



Retained from Run I
LAr calorimeter
Central muon detector
Muon toroid

New for Run II

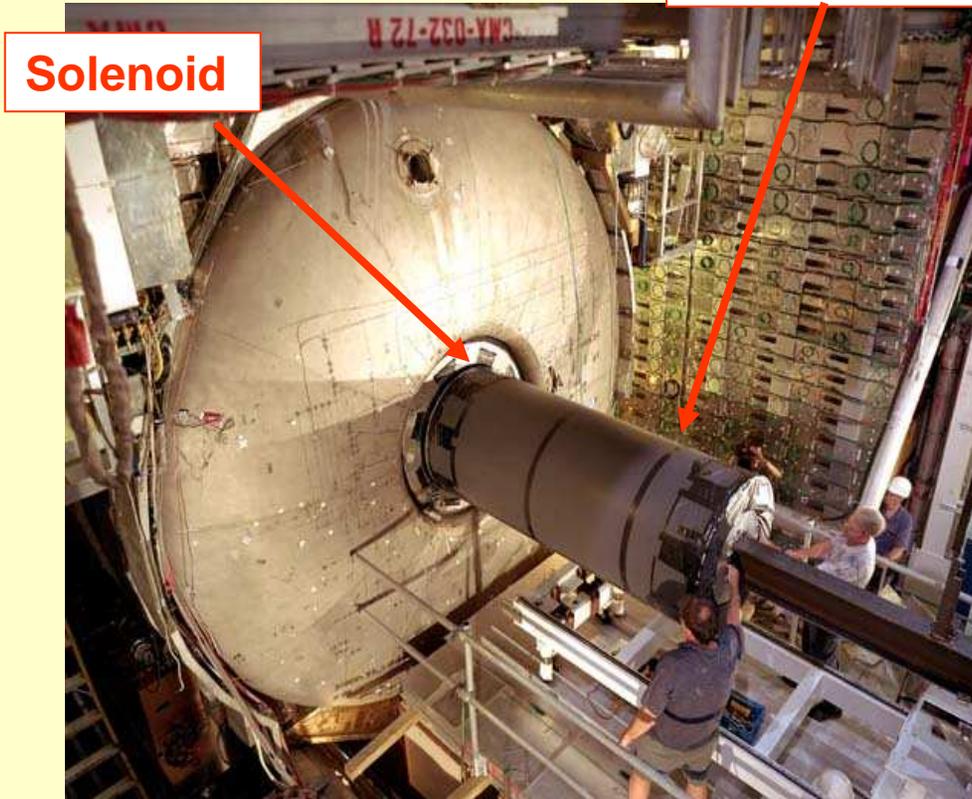
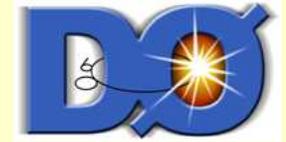
Inner detector
(tracking)
Magnetic field added

Preshower detectors
Forward muon detector

Front-end electronics
Trigger and DAQ

In addition: Inner B-layer
(similar to CDF)

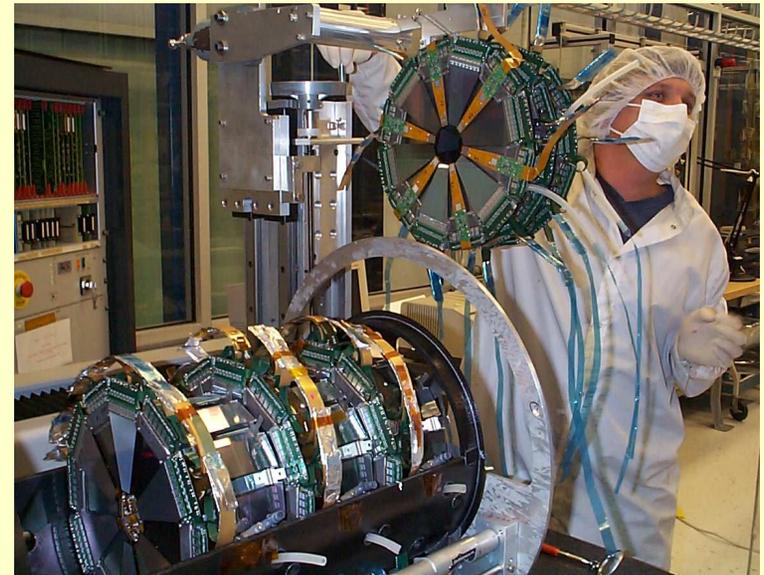
DØ Detector



Solenoid

Fiber Tracker

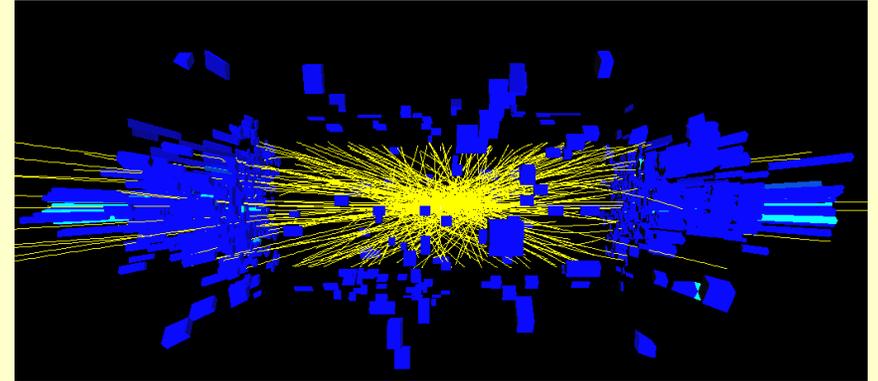
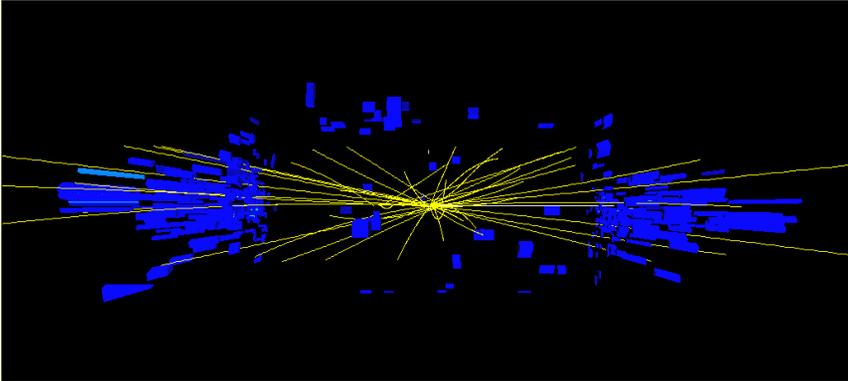
Silicon Detector



Challenges with high luminosity

Min. bias pileup at the Tevatron, at $0.6 \cdot 10^{32} \text{ cm}^2\text{s}^{-1}$

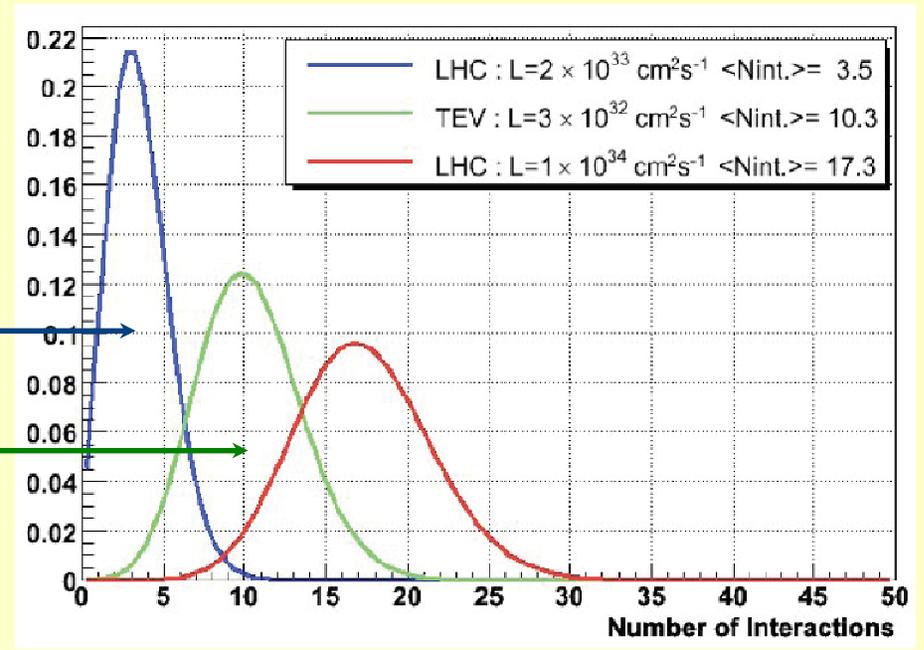
... and at $2.4 \cdot 10^{32} \text{ cm}^2\text{s}^{-1}$



Average number of interactions:

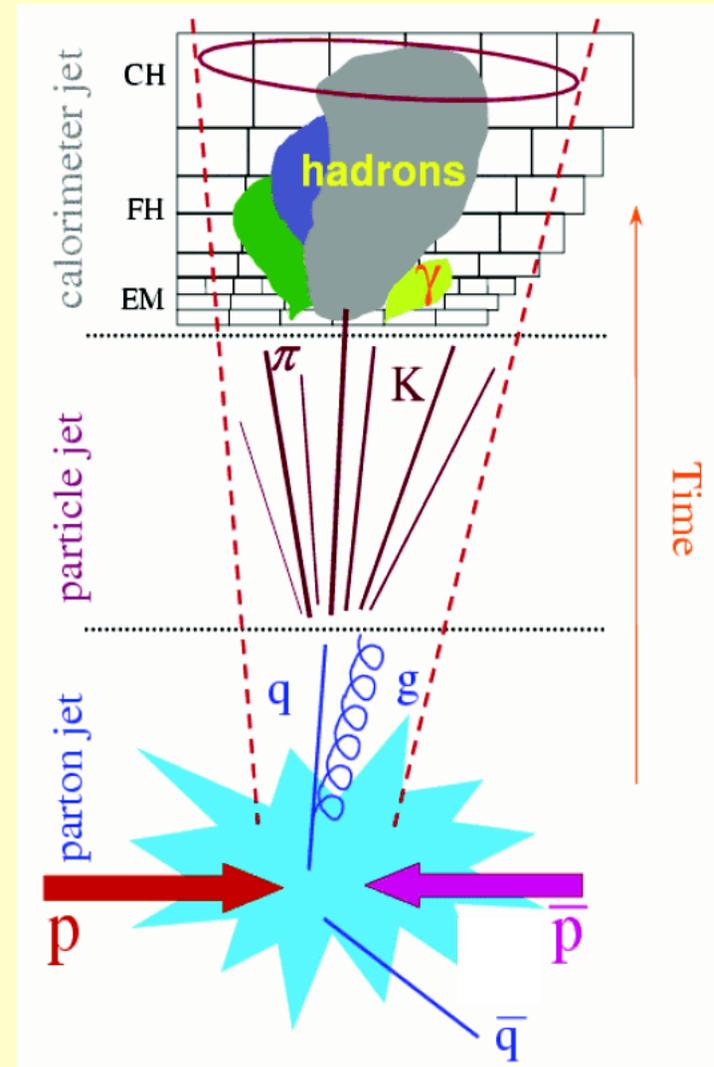
LHC: initial “low” luminosity run
($L=2 \cdot 10^{33} \text{ cm}^2\text{s}^{-1}$): $\langle N \rangle = 3.5$

TeV: ($L=3 \cdot 10^{32} \text{ cm}^2\text{s}^{-1}$): $\langle N \rangle = 10$



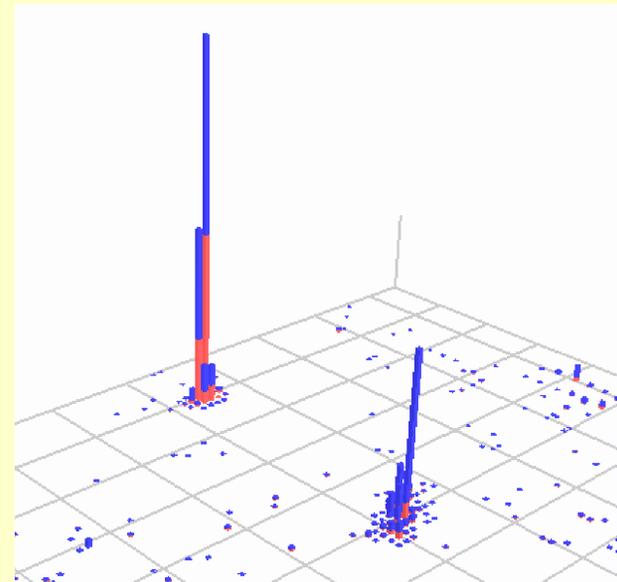
Jet reconstruction and energy measurement

- A jet is NOT a well defined object
(fragmentation, gluon radiation, detector response)
- The detector response is different for particles interacting electromagnetically (e, γ) and for hadrons
→ for comparisons with theory, one needs to correct back the calorimeter energies to the „particle level“ (particle jet)
Common ground between theory and experiment
- One needs an algorithm to define a jet and to measure its energy
conflicting requirements between experiment and theory (exp. simple, e.g. cone algorithm, vs. theoretically sound (no infrared divergencies))
- Energy corrections for losses of fragmentation products outside jet definition and underlying event or pileup energy inside



Main corrections:

- In general, calorimeters show different response to electrons/photons and hadrons
- Subtraction of offset energy not originating from the hard scattering (inside the same collision or pile-up contributions, use minimum bias data to extract this)
- Correction for jet energy out of cone (corrected with jet data + Monte Carlo simulations)



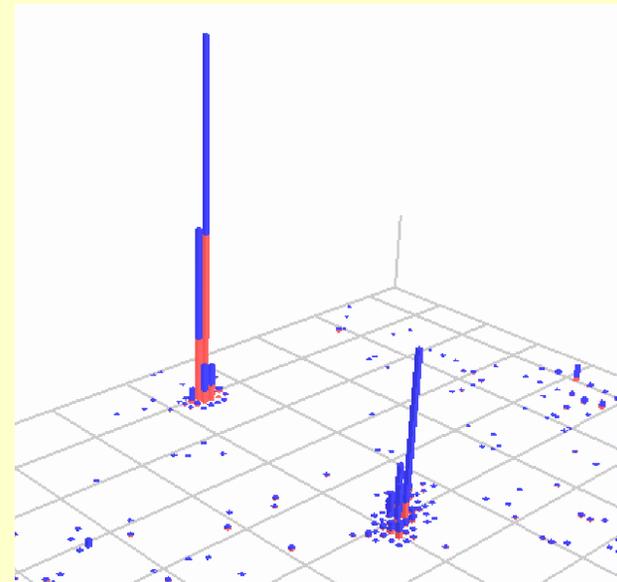
Main performance parameters of the different hadronic calorimeter components of the ATLAS and CMS detectors, as measured in test beams using charged pions in both stand-alone and combined mode with the ECAL

	ATLAS					
	Barrel LAr/Tile		End-cap LAr		CMS	
	Tile	Combined	HEC	Combined	Had. barrel	Combined
Electron/hadron ratio	1.36	1.37	1.49			
Stochastic term	$45\%/\sqrt{E}$	$55\%/\sqrt{E}$	$75\%/\sqrt{E}$	$85\%/\sqrt{E}$	$100\%/\sqrt{E}$	$70\%/\sqrt{E}$
Constant term	1.3%	2.3%	5.8%	< 1%		8.0%
Noise	Small	3.2 GeV		1.2 GeV	Small	1 GeV

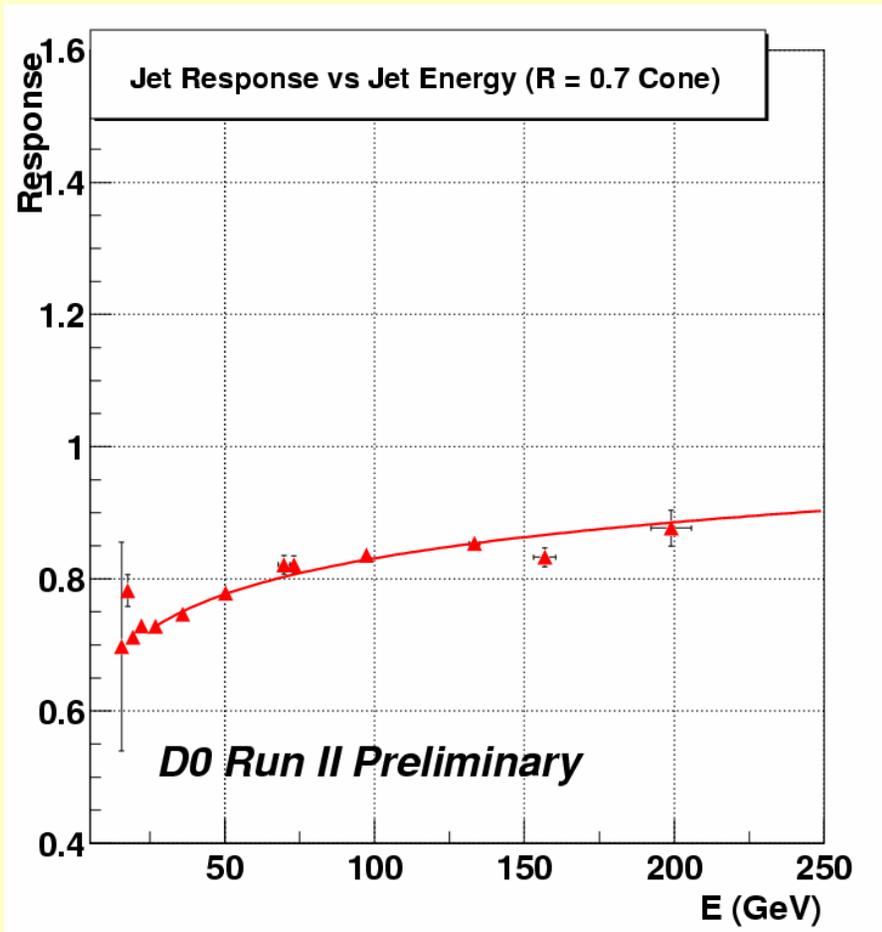
The measured electron/hadron ratios are given separately for the hadronic stand-alone and combined calorimeters when available, and for the contributions (added quadratically except for the stand-alone ATLAS tile calorimeter) to the pion energy resolution from the stochastic term, the local constant term, and the noise are also shown, when available from published data.

Main corrections:

- In general, calorimeters show different response to electrons/photons and hadrons
- Subtraction of offset energy not originating from the hard scattering (inside the same collision or pile-up contributions, use minimum bias data to extract this)
- Correction for jet energy out of cone (corrected with jet data + Monte Carlo simulations)

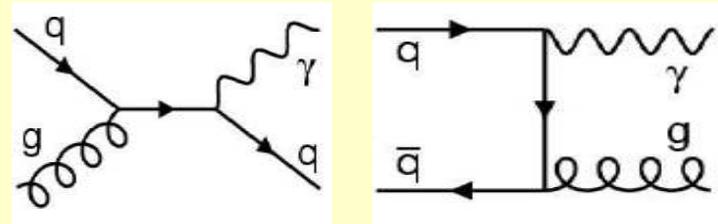


Jet Energy Scale



Jet response correction in DØ:

- Measure response of particles making up the jet
- Use photon + jet data - calibrate jets against the better calibrated photon energy



- Achieved jet energy scale uncertainty:

DØ: $\Delta E / E \sim 1-2\%$
(excellent result, a huge effort)

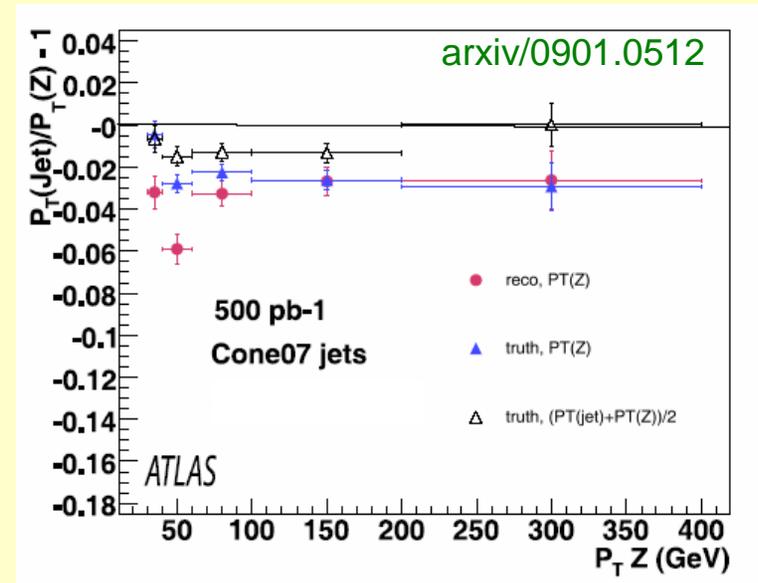
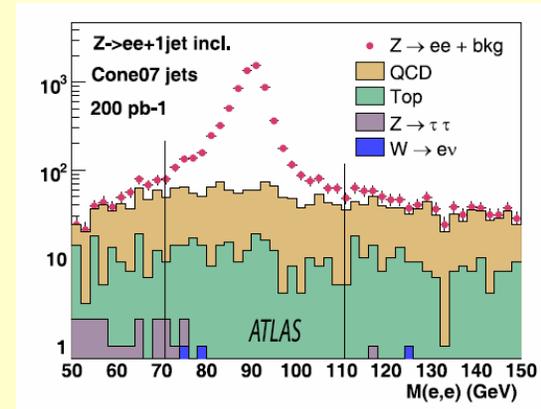
Jet energy scale at the LHC

- A good jet-energy scale determination is essential for many QCD measurements (arguments similar to Tevatron, but kinematic range (jet p_T) is larger, ~ 20 GeV – ~ 3 TeV)
- Propagate knowledge of the EM scale to the hadronic scale, but several processes are needed to cover the large p_T range

Measurement process	Jet p_T range
Z + jet balance	$20 < p_T < 100 - 200$ GeV
γ + jet balance	$50 < p_T < 500$ GeV (trigger, QCD background)
Multijet balance	500 GeV $< p_T$

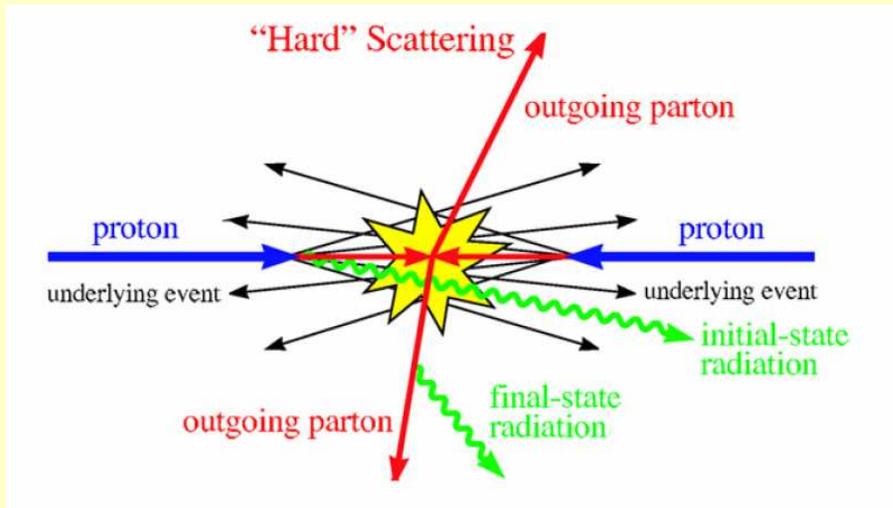
Reasonable goal: 5-10% in first runs (1 fb^{-1})
1- 2% long term

Example: Z + jet balance



Stat. precision (500 pb^{-1}): 0.8%
Systematics: 5-10% at low p_T , 1% at high p_T

Study of minimum bias events



... and of the underlying event

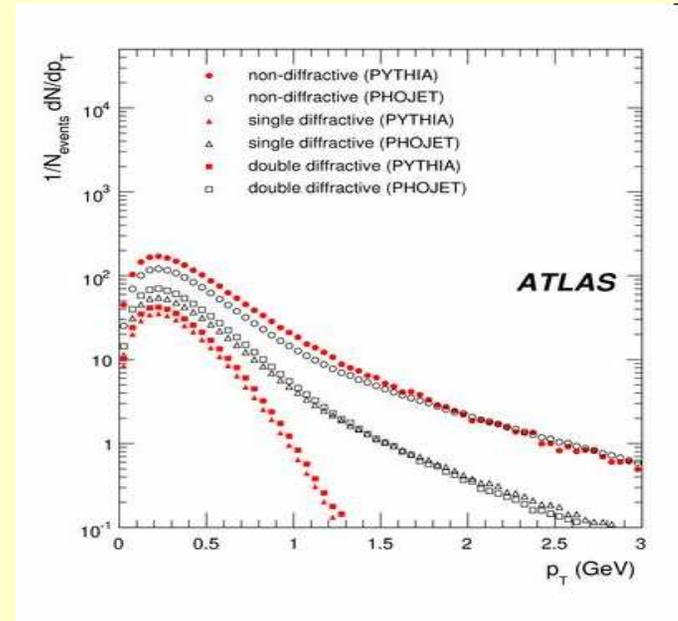
Understanding and modelling of the underlying event and min. bias events is important for:

- Simulation of pileup effects at the LHC
- Understanding of lepton and jet isolation
- Event selections with jet vetos (often low p_T (~ 20 GeV) jet vetos used in searches, e.g. $H \rightarrow WW \rightarrow \ell\nu \ell\nu$)
- Calibration of the jet energy scale

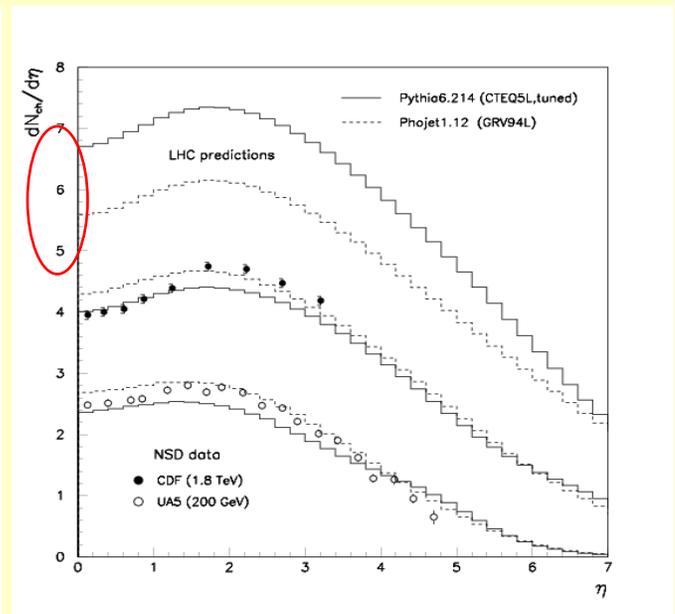
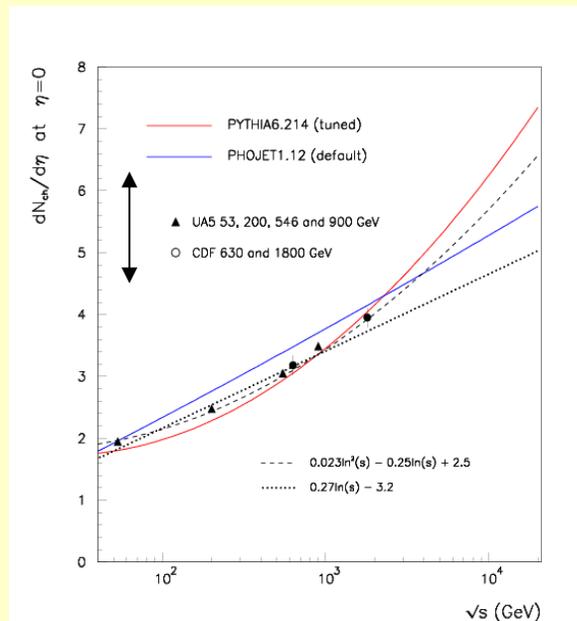
Measurement of properties of minimum bias events

First measurement at the LHC

- Measure charged particle distributions: **rapidity distribution and p_T -spectrum**
- Multiplicity distributions and $\langle p_T \rangle$
- Large uncertainties on model predictions



$\langle p_T \rangle$ ($\eta = 0$): 550 – 640 MeV (15%)

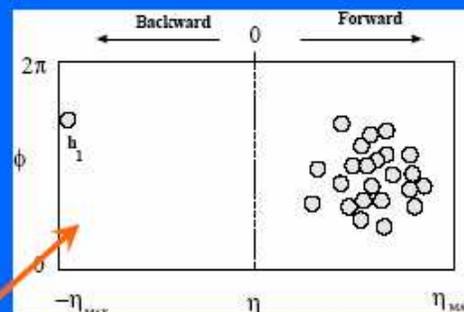


$dN_{ch}/d\eta$ ($\eta=0$): 5-7 (~ 33%)

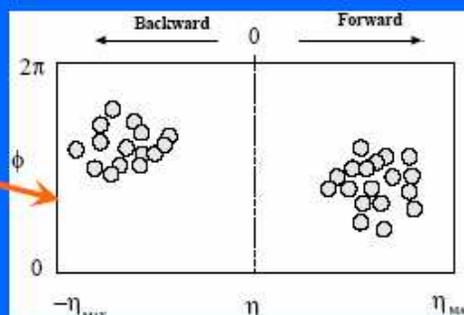
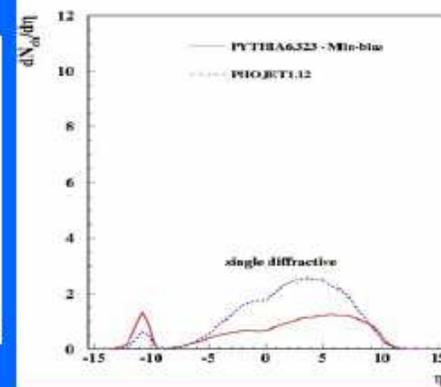
Soft pp collisions

pp collisions at $\sqrt{s} = 14\text{TeV}$	PYTHIA6.323	PHOJET1.12
σ_{tot}	101.5 mb	119.1 mb
σ_{elas}	22.2 mb	34.5mb
$2^*\sigma_{\text{SD}}$	14.4mb	11.0mb
σ_{DD}	10.3mb	4.1mb
σ_{ND}	54.7mb	69.5mb

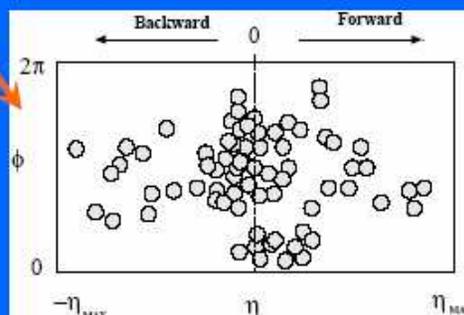
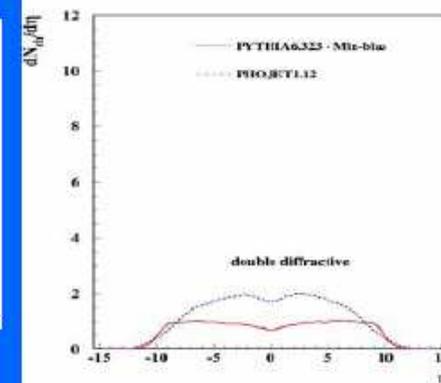
Minimum bias
Made up of
combination of
non-diffractive
and diffractive



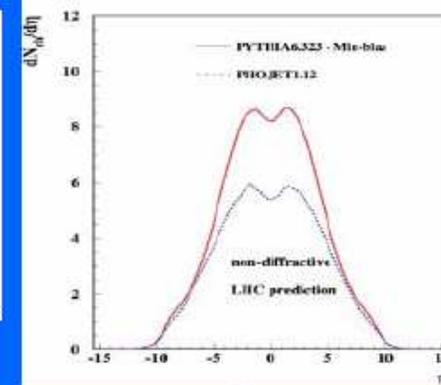
Single diffractive SD



Double diffractive DD



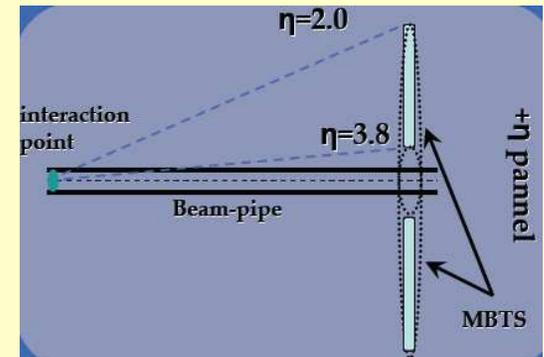
Non-diffractive ND



Early QCD measurements with ATLAS, DIS09

“Minimum bias events”

- Minimum bias is an experimental definition, defined by experimental trigger selection and analysis
- Relation to Physics:



$$\sigma_{\text{measured}} = f_{\text{sd}} \sigma_{\text{sd}} + f_{\text{dd}} \sigma_{\text{dd}} + f_{\text{nd-inelelastic}} \sigma_{\text{nd-inelelastic}}$$

where f_i are the efficiencies for different physics processes determined by the trigger

NB: need to understand what is measured to allow comparison to previous results, often presented for non-single diffractive events



Present experimental preparations / studies

Measurements of minimum bias physics require special **triggers** and **reconstruction**:

Trigger:

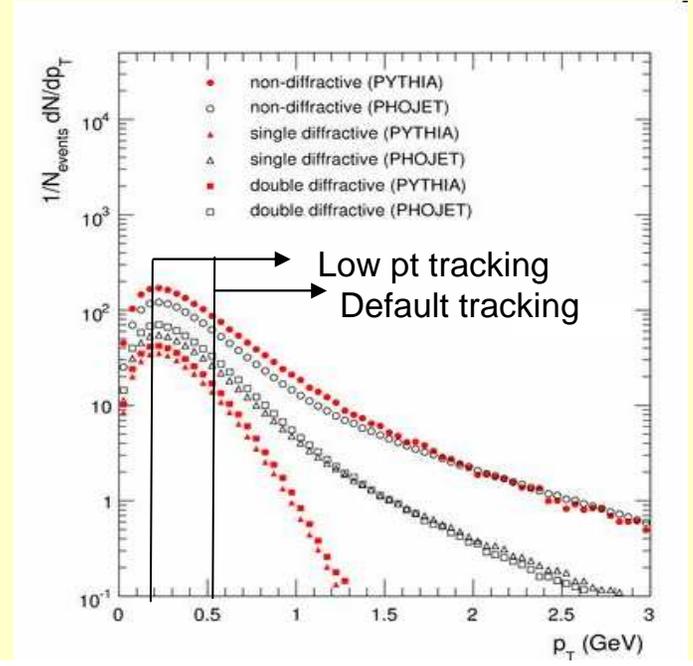
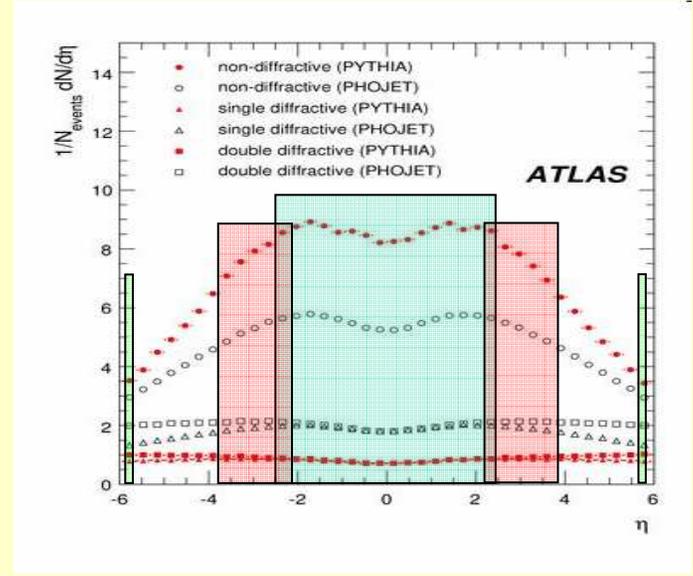
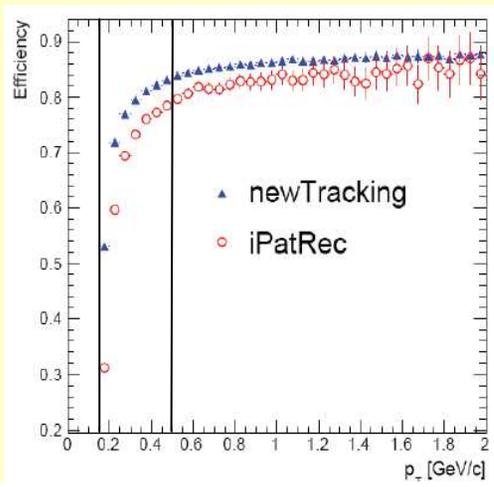
For **early running** up to $\sim 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$, number of events/crossing $\ll 1$

- Inner detector space points and tracks $|\eta| < 2.5$
- Trigger scintillators (MBTS) $2.1 < |\eta| < 3.8$
- LUCID detector $5.6 < |\eta| < 5.9$
- Zero degree calorimeter (ZDC) $|\eta| > 8.3$

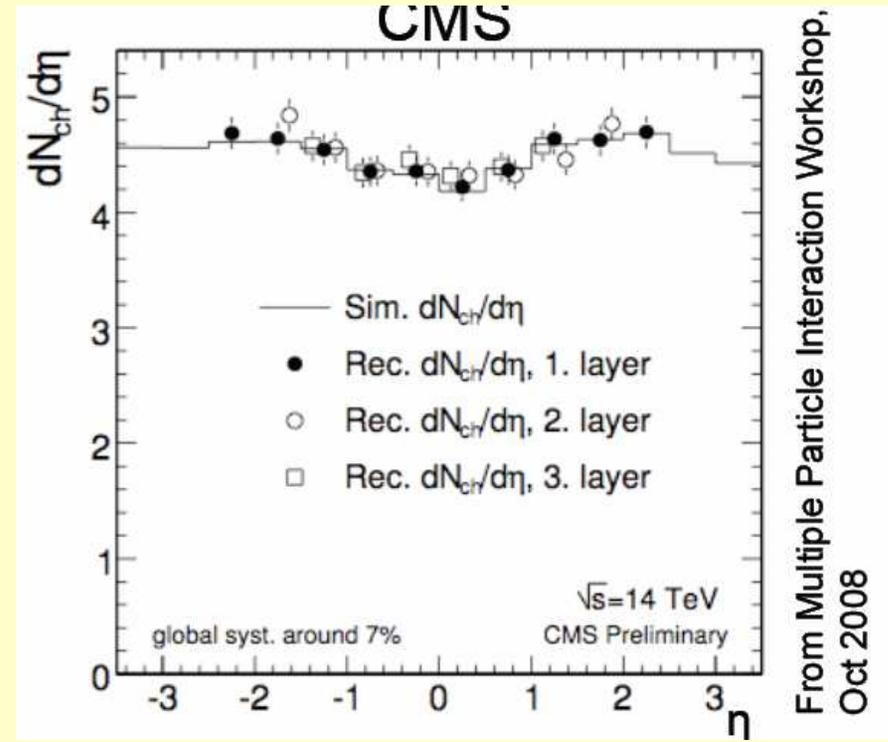
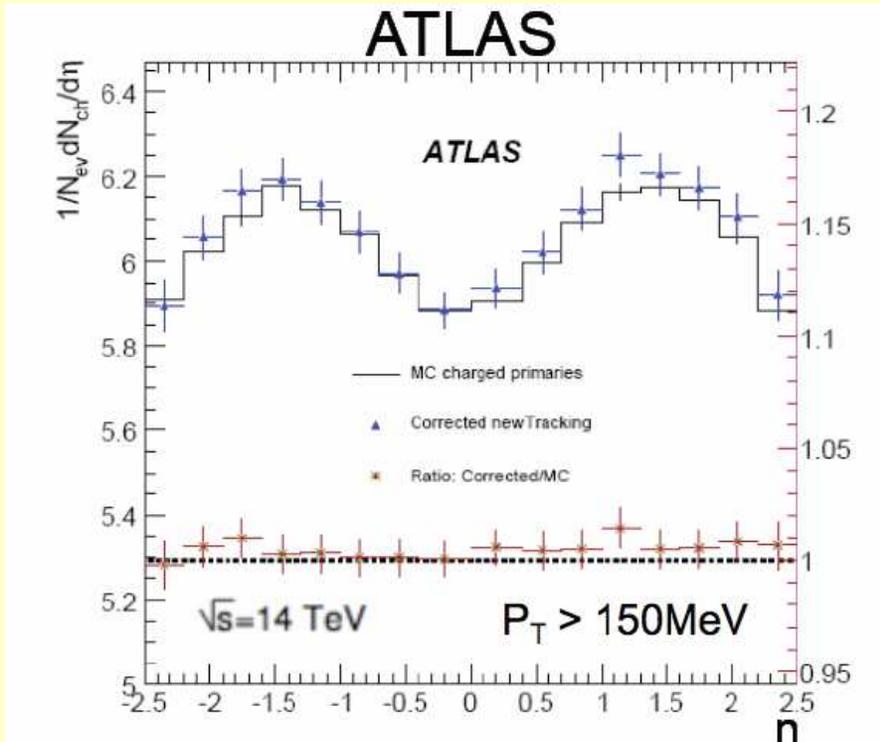
Later, for $L = 10^{33} - 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$: use random trigger

Reconstruction:

Track reconstruction down to very low p_T is required



An ATLAS / CMS comparison



From Multiple Particle Interaction Workshop,
Oct 2008

Uses a **tracking-based** method

Dominant uncertainties from the Inner Detector misalignment and diffractive cross sections

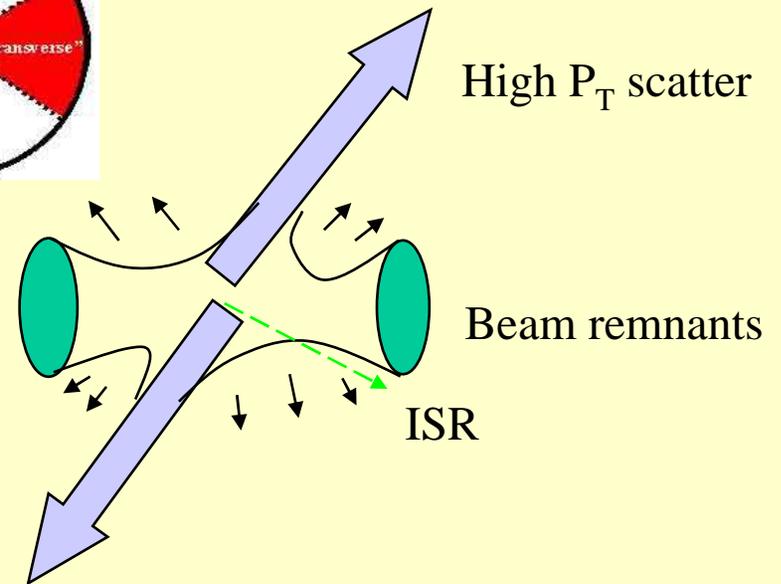
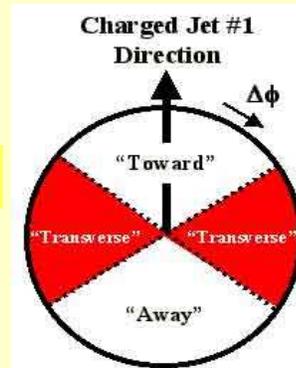
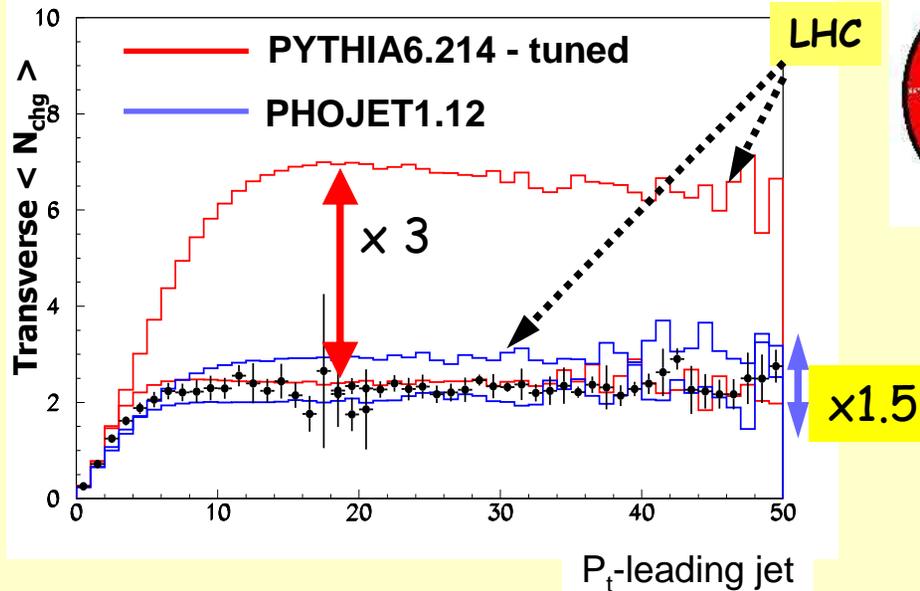
Goal: total systematic uncertainty $\sim 8\%$

Uses a **hit-counting** method

Dominant uncertainties from reconstruction (hit numbers to charged particle conversion functions)

The underlying event

Average charged particle density in transverse region



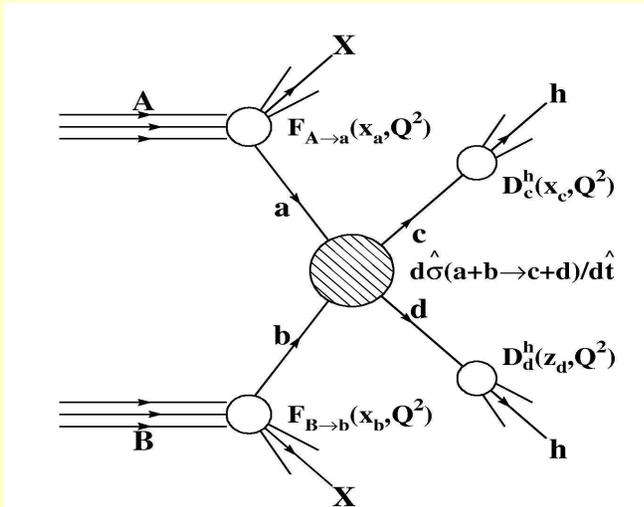
Extrapolation of the underlying event to LHC energies is unknown;

underlying event depends on:

- Multiple interactions
- Radiation
- PDFs
- String formation

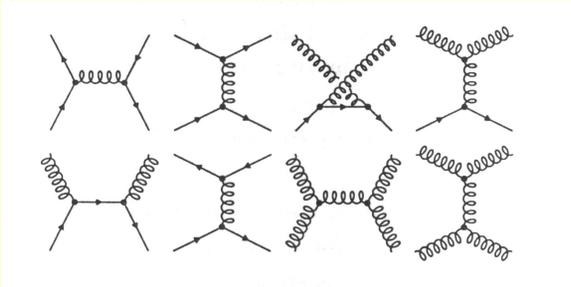
- A lot of Monte Carlo tuning is needed;
- Early measurements at the LHC (low p_{T} jets, but also in W/Z production) will considerably extend our knowledge

QCD processes at hadron colliders

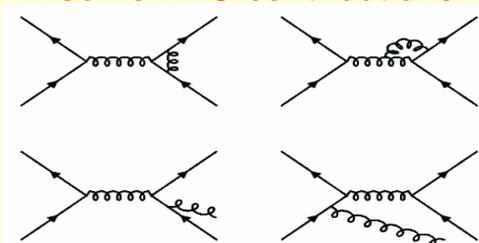


- Hard scattering processes are dominated by QCD jet production
- Originating from qq , qg and gg scattering
- Cross sections can be calculated in QCD (perturbation theory)

Leading order



...some NLO contributions



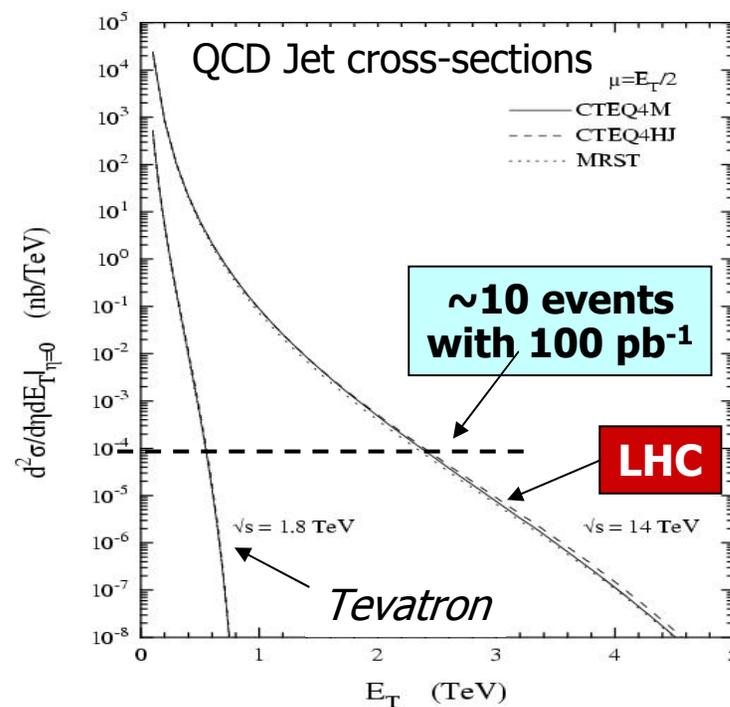
Comparison between experimental data and theoretical predictions constitutes an important test of the theory.

Deviations?

- Problem in the experiment ?
- Problem in the theory (QCD) ?
- New Physics, e.g. quark substructure ?

Jets from QCD production: Tevatron vs LHC

- Rapidly probe perturbative QCD in a new energy regime (at a scale above the Tevatron, large cross sections)
- **Experimental challenge:** understanding of the detector
 - main focus on **jet energy scale**
 - resolution
- **Theory challenge:**
 - improved calculations... (renormalization and factorization scale uncertainties)
 - pdf uncertainties



In addition to QCD test: Sensitivity to New Physics

- Contact interactions:

Despite the relatively large jet energy scale uncertainties (5-10%) expected with **early data**, the LHC has large sensitivity to contact interactions parametrized by a scale parameter Λ

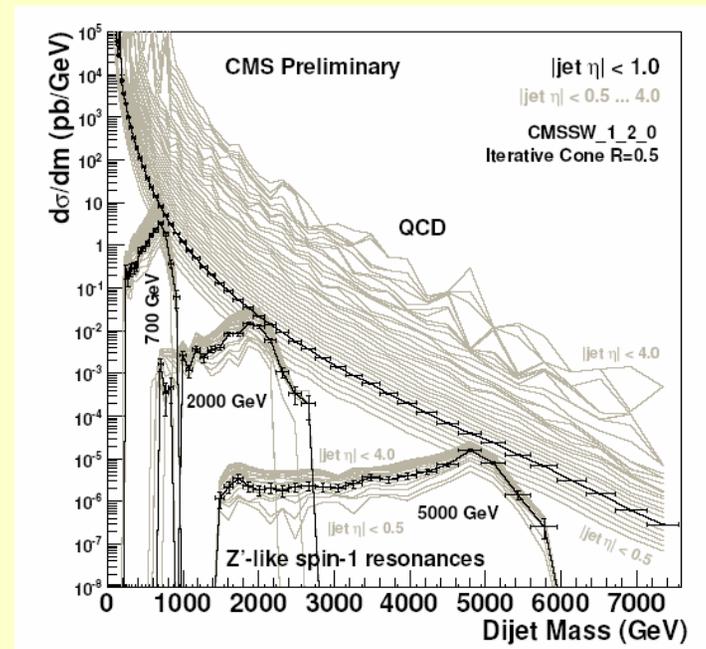
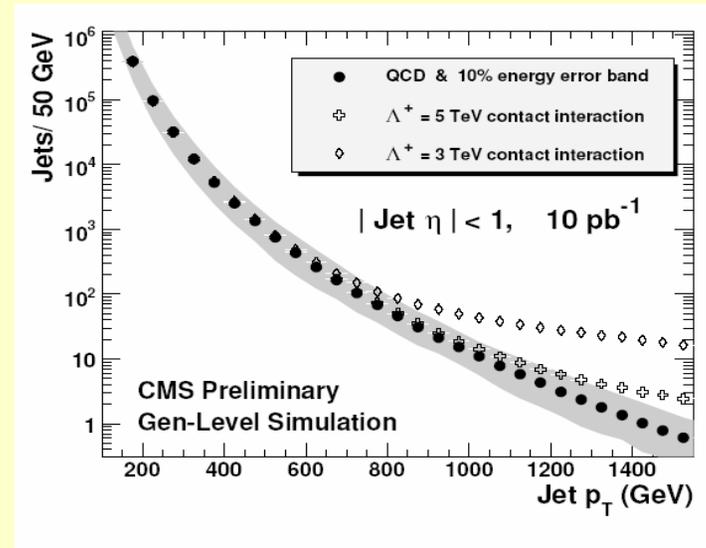
Search for deviations from QCD in the high p_T region

- Heavy resonances decaying into jets

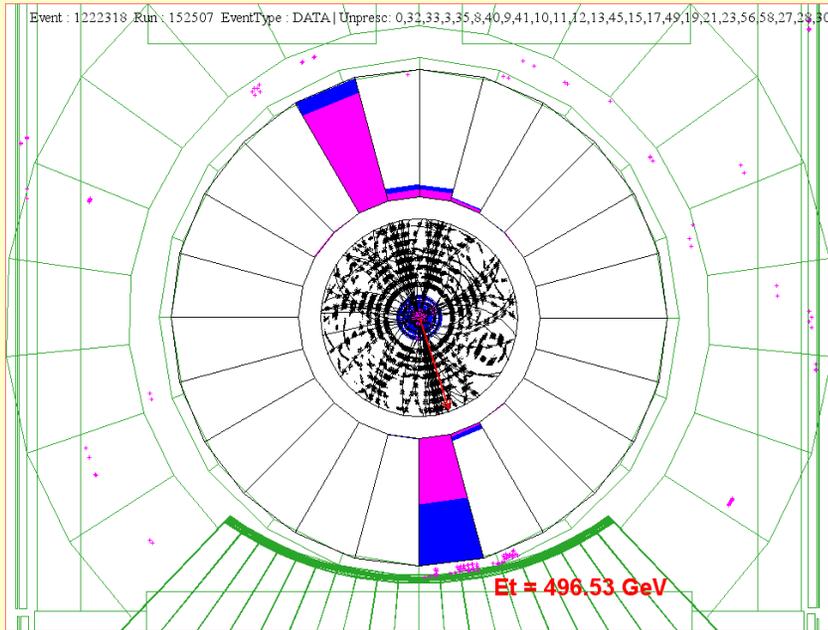
e.g. $Z' \rightarrow qq$

Search for resonant structures in dijet invariant mass spectrum

→ estimates tomorrow



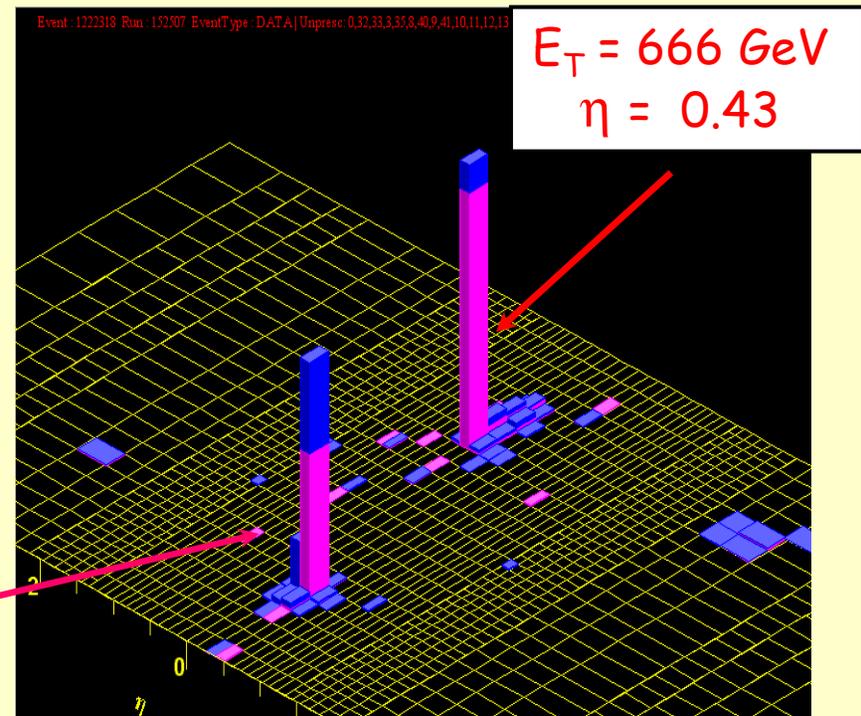
A two jet event at the Tevatron (CDF)



CDF (ϕ -r view)

$E_T = 633 \text{ GeV}$
 $\eta = -0.19$

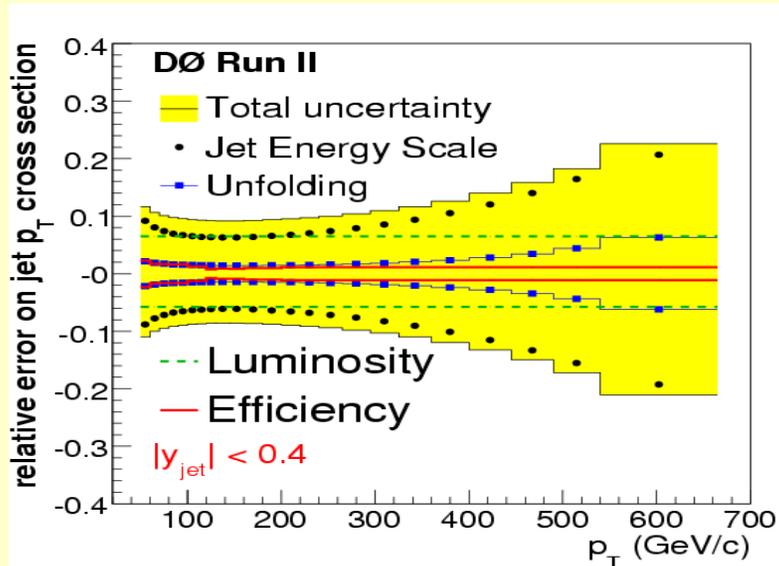
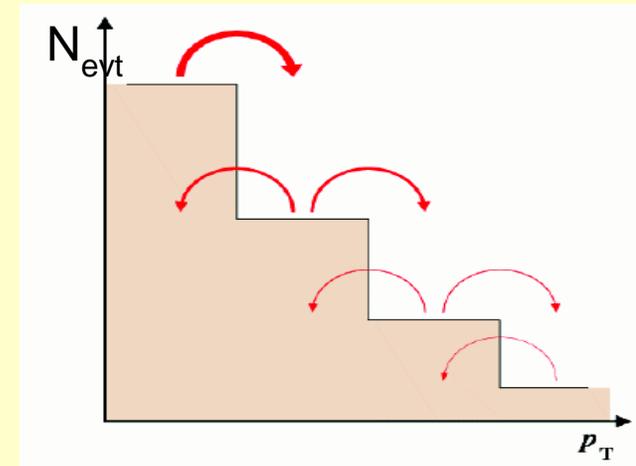
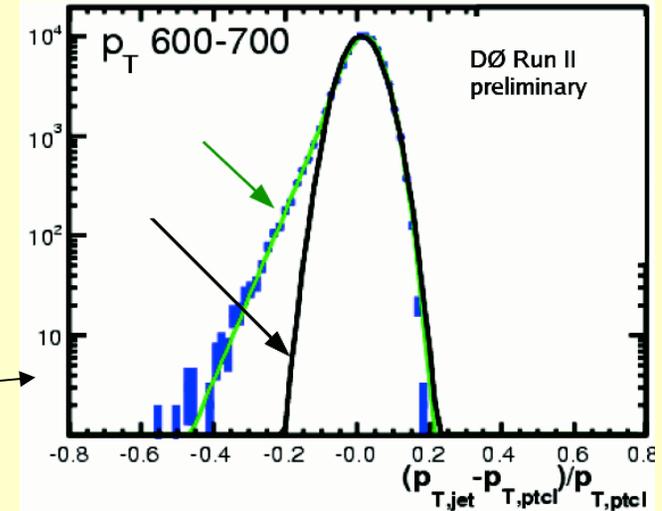
Dijet Mass = $1364 \text{ GeV}/c^2$



Jet measurements

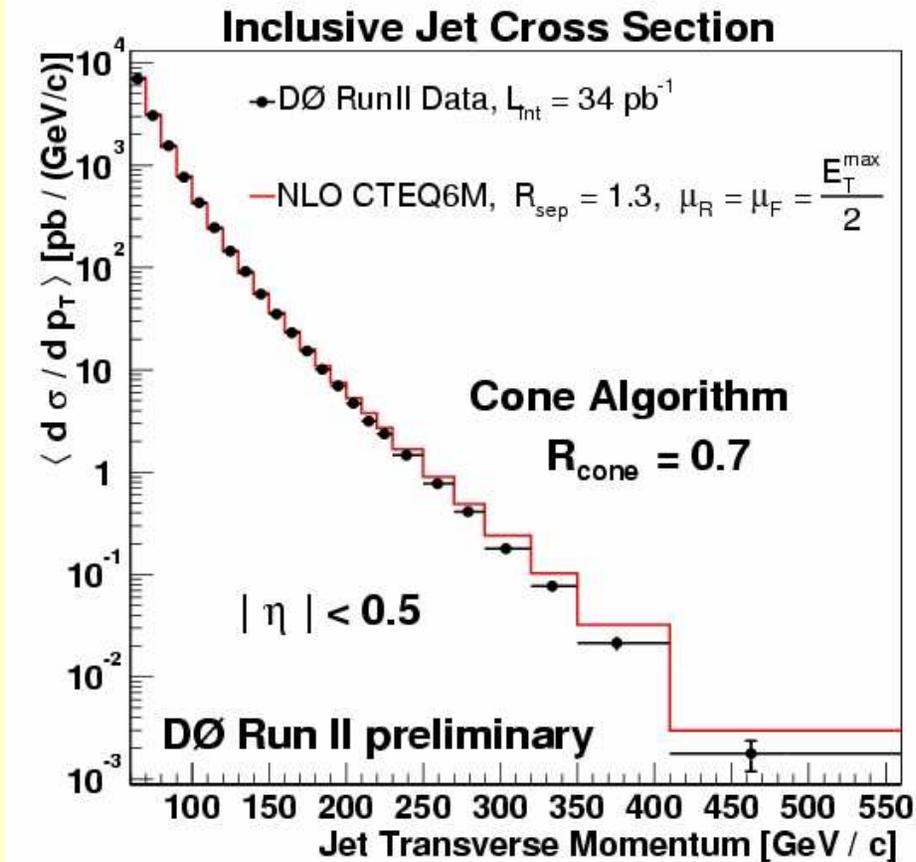
$$d^2\sigma / dp_T d\eta = N / (\epsilon \cdot L \cdot \Delta p_T \cdot \Delta\eta)$$

- In principle a simple counting experiment
- However, steeply falling p_T spectra are sensitive to jet energy scale uncertainties and resolution effects (migration between bins) → corrections (unfolding) to be applied
- Sensitivity to jet energy scale uncertainty:
 DØ: 1% energy scale error
 → 10% cross section uncert. at $|\eta| < 0.4$



Major exp. errors:
 energy scale, luminosity (6%),...

Test of QCD Jet production



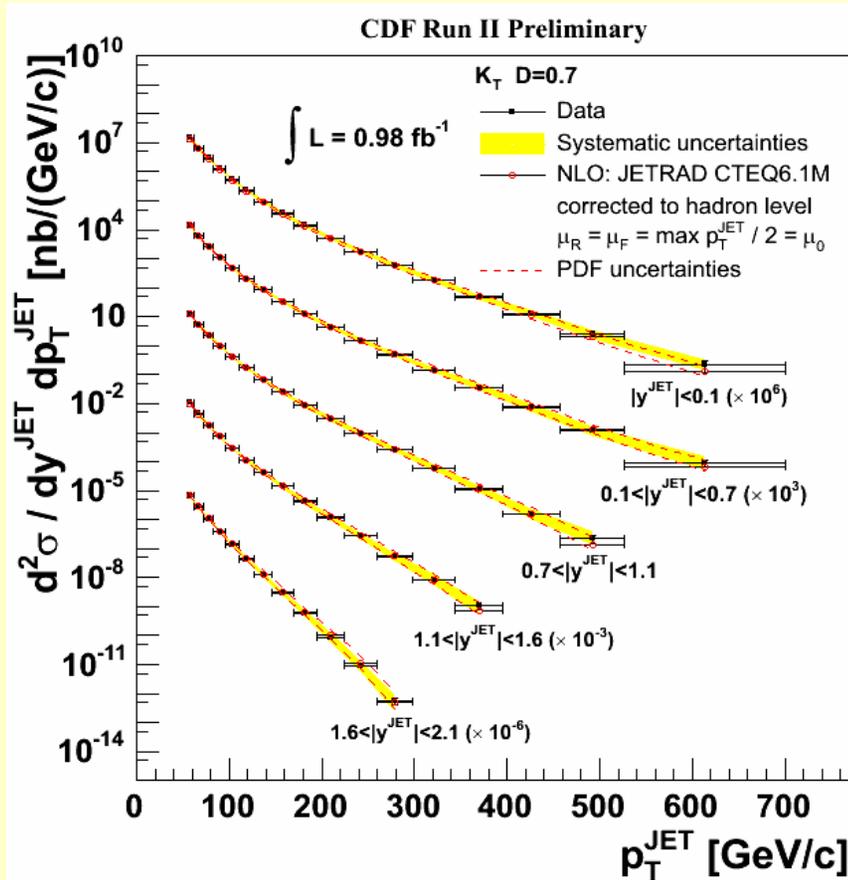
An “**early**” result from the DØ experiment (34 pb^{-1})

Inclusive Jet spectrum as a function of Jet- P_T

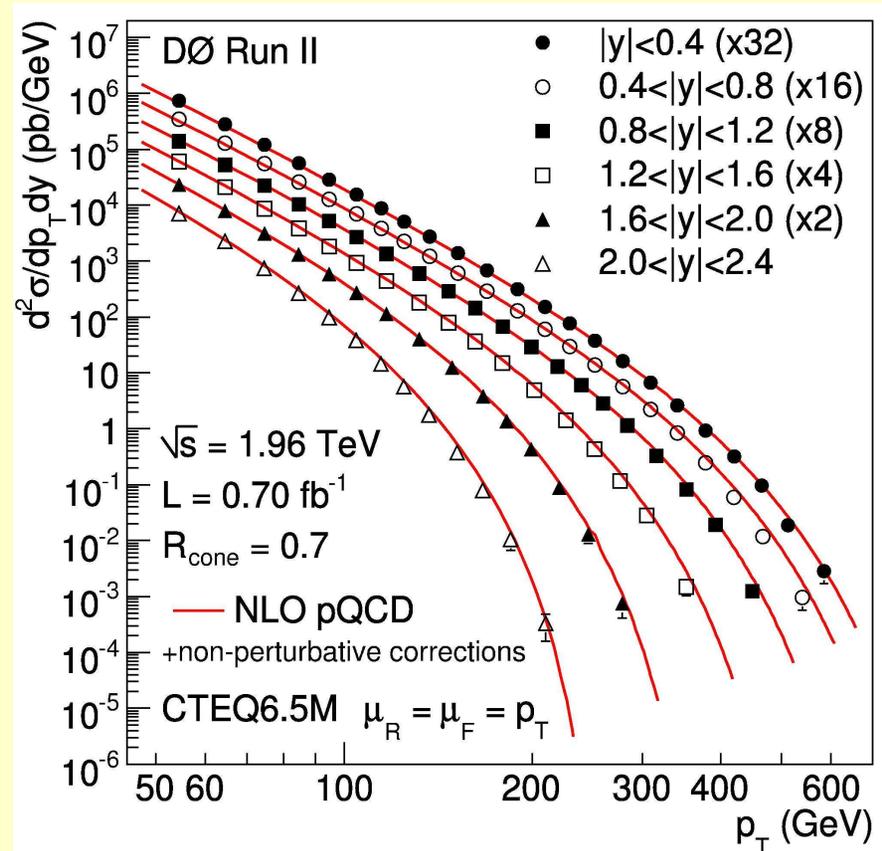
very good agreement with NLO pQCD calculations over many orders of magnitude !

within the large theoretical and experimental uncertainties

Double differential distributions in p_T and η



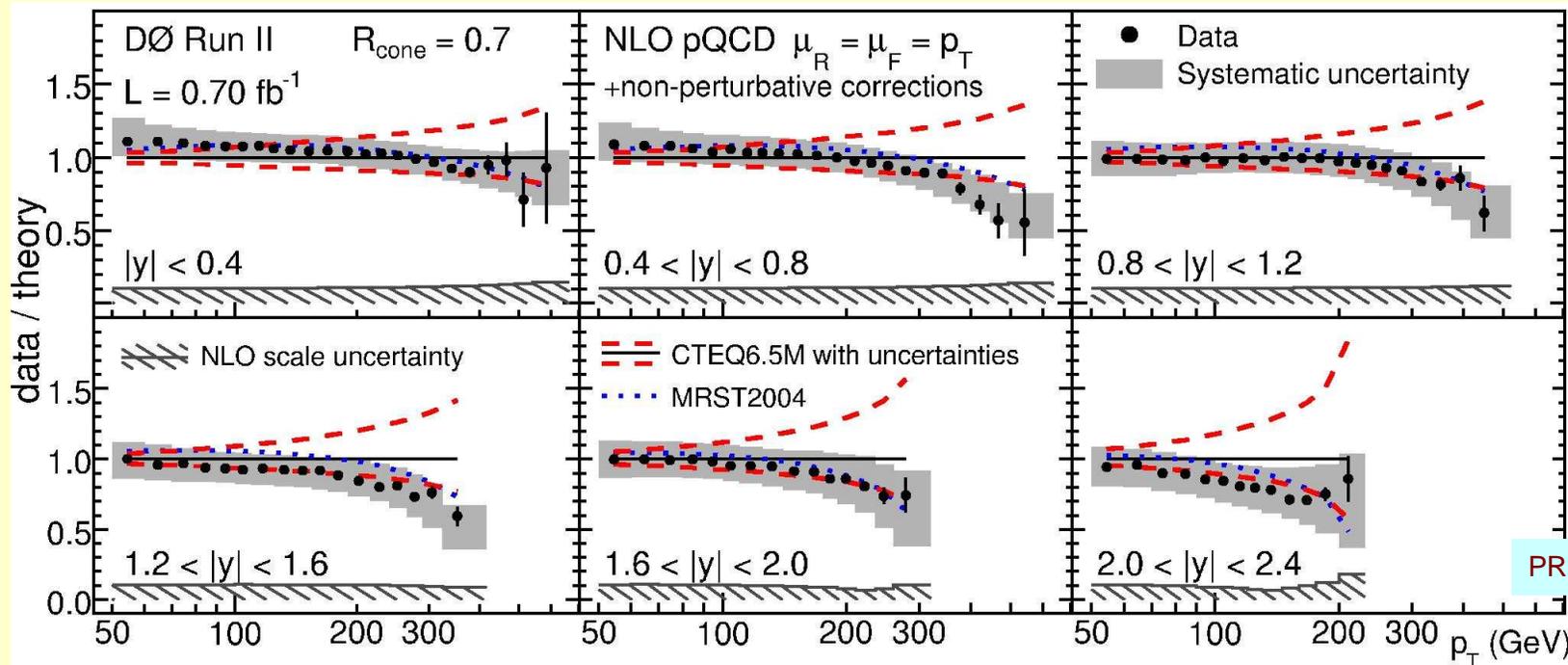
PRD 78 052006 ('08)



PRL 101 062001 ('08)

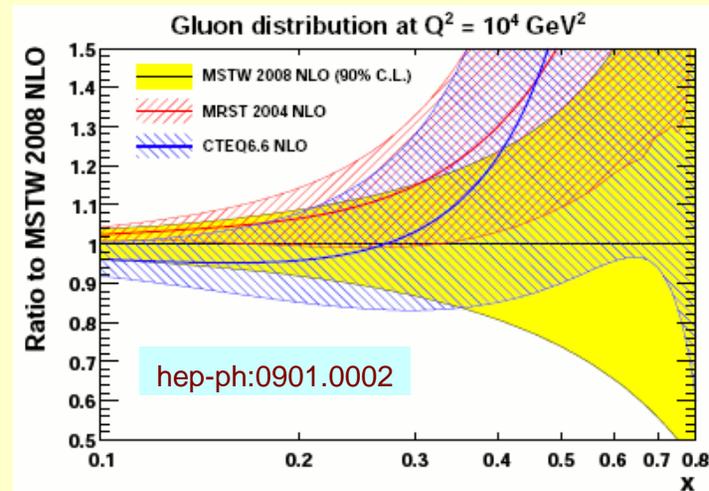
- Measurement in 5-6 different rapidity bins, over 9 orders of magnitude, up to $p_T \sim 650 \text{ GeV}$
- Data corresponding to $\sim 1 \text{ fb}^{-1}$ (CDF) and 0.7 fb^{-1} (DØ)

Comparison between data and theory



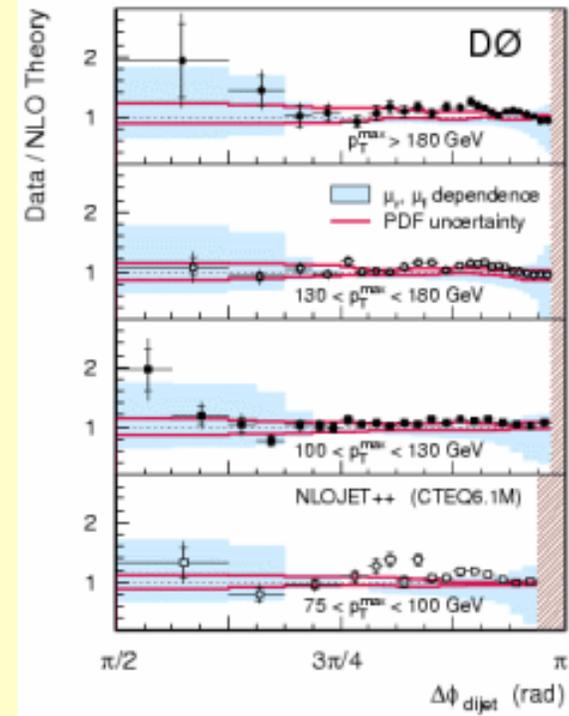
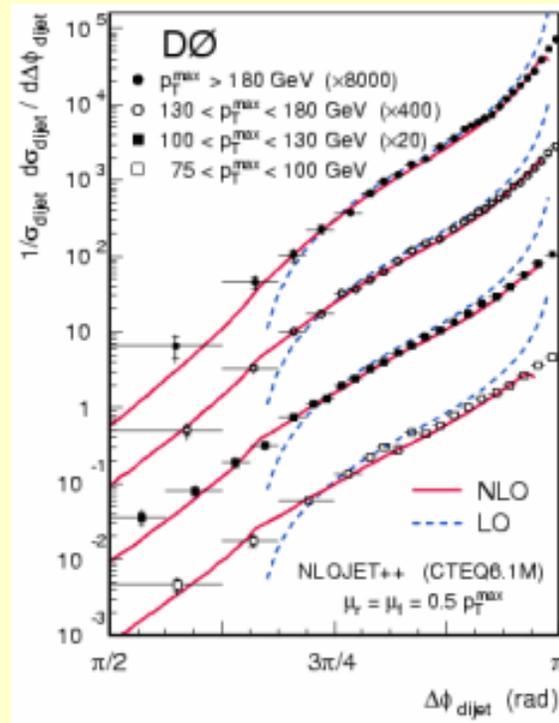
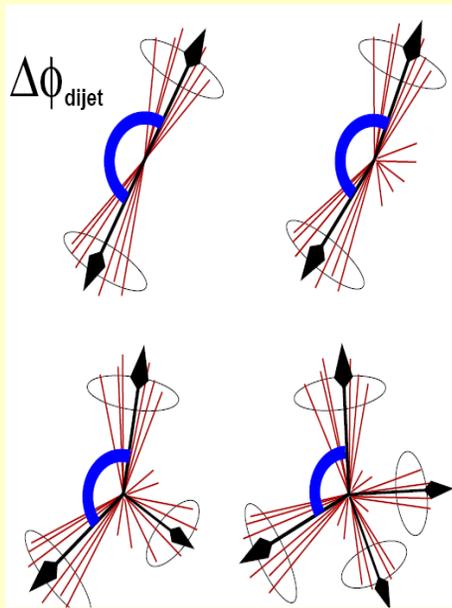
PRL 101 062001 ('08)

- CDF and DØ agree within uncertainties
- Experimental uncertainties are smaller than the pdf uncertainties (in particular large for large x , gluon distribution)
- Wait for updated (2009) parametrizations (plans to include Tevatron data, to better constrain the high x -region)



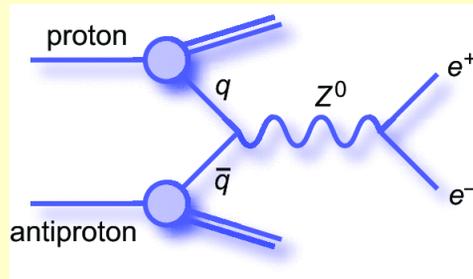
Di-jet angular distributions

- *reduced sensitivity to Jet energy scale*
- *sensitivity to higher order QCD corrections preserved*

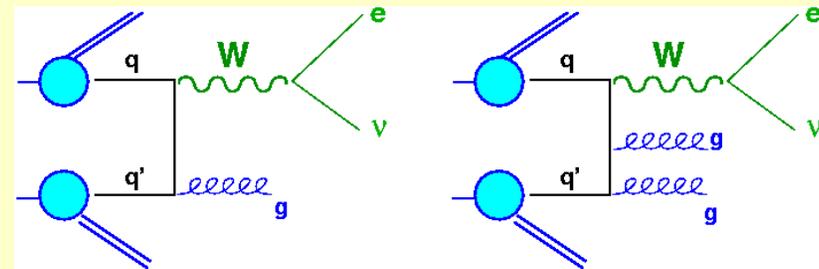


**Good agreement with
Next-to-leading order QCD-predictions**

QCD aspects in W/Z (+ jet) production



QCD at work

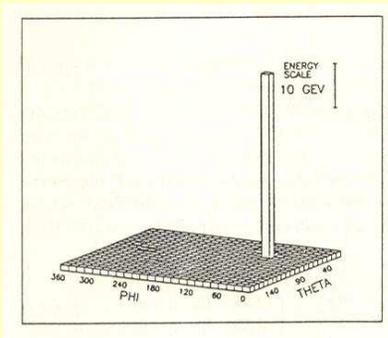


- Important test of NNLO Drell-Yan QCD prediction for the total cross section
- Test of perturbative QCD in high p_T region (jet multiplicities, p_T spectra,....)
- Tuning and „calibration“ of Monte Carlos for background predictions in searches at the LHC

How do W and Z events look like ?

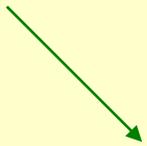
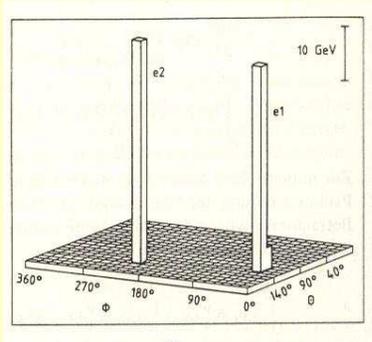
As explained, leptons, photons and missing transverse energy are key signatures at hadron colliders

→ Search for leptonic decays: $W \rightarrow \ell \nu$ (large $P_T(\ell)$, large P_T^{miss})
 $Z \rightarrow \ell \ell$

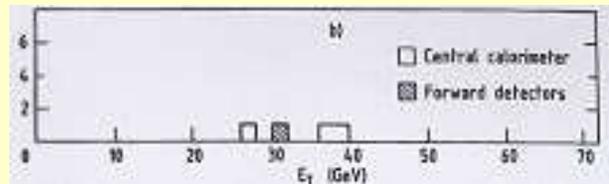


A bit of history: one of the first W events seen; UA2 experiment

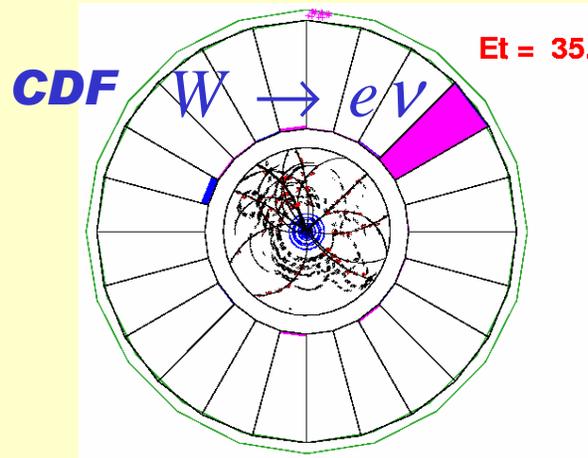
W/Z discovery by the UA1 and UA2 experiments at CERN (1983/84)



Transverse momentum of the electrons



Today's W / Z \rightarrow ev / ee signals



Trigger:

- Electron candidate > 20 GeV/c

Electrons

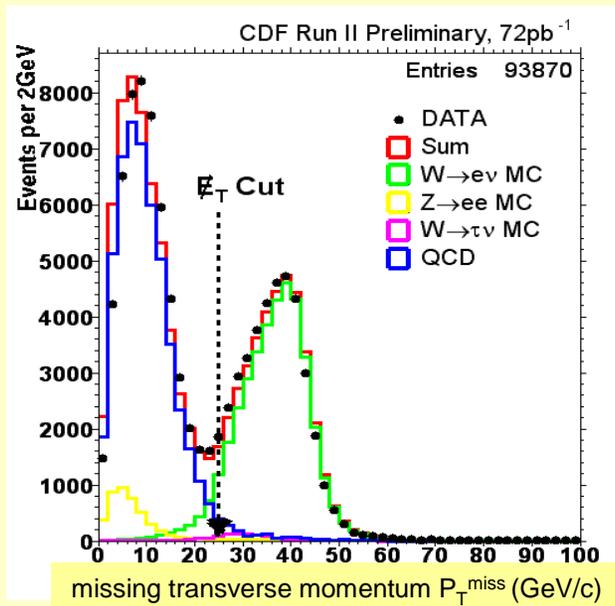
- Isolated el.magn. cluster in the calorimeter
- $P_T > 25$ GeV/c
- Shower shape consistent with expectation for electrons
- Matched with tracks

Z \rightarrow ee

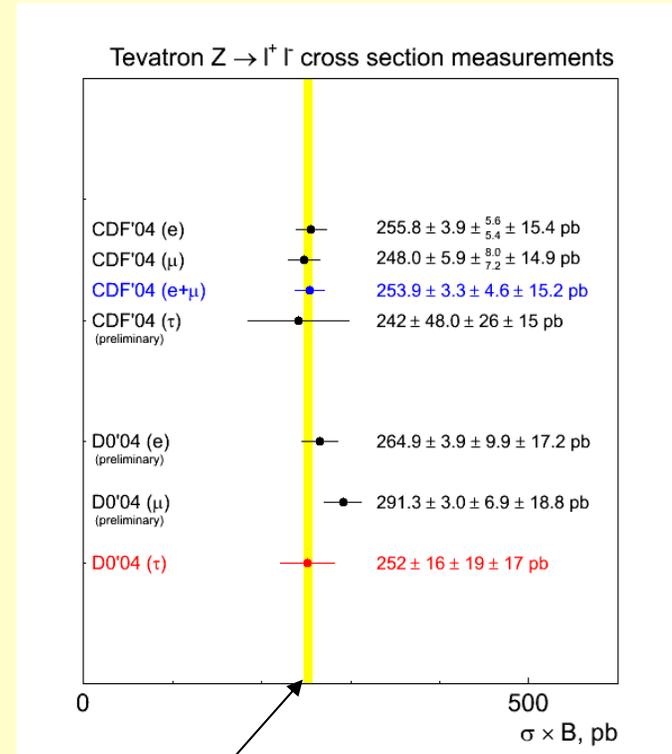
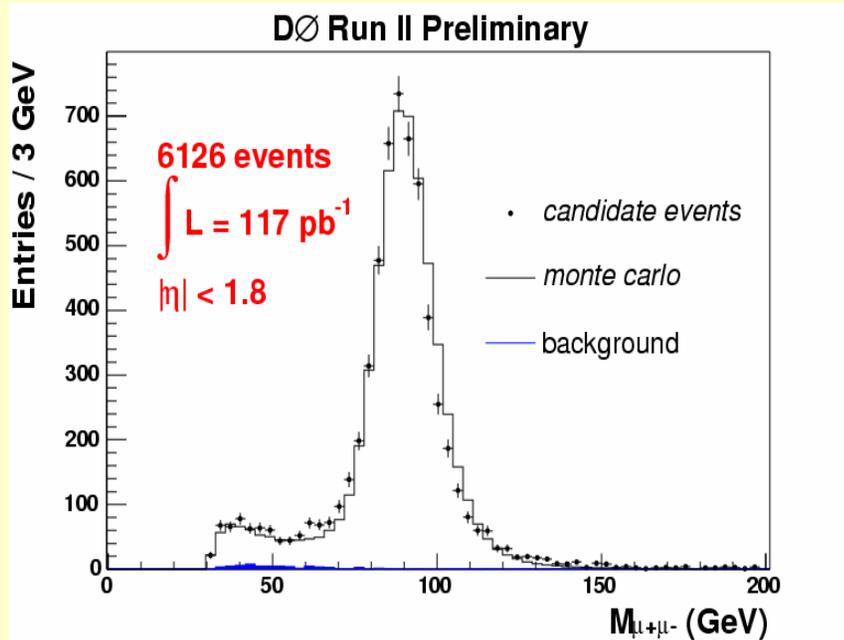
- $70 \text{ GeV}/c^2 < m_{ee} < 110 \text{ GeV}/c^2$

W \rightarrow ev

- Missing transverse momentum > 25 GeV/c



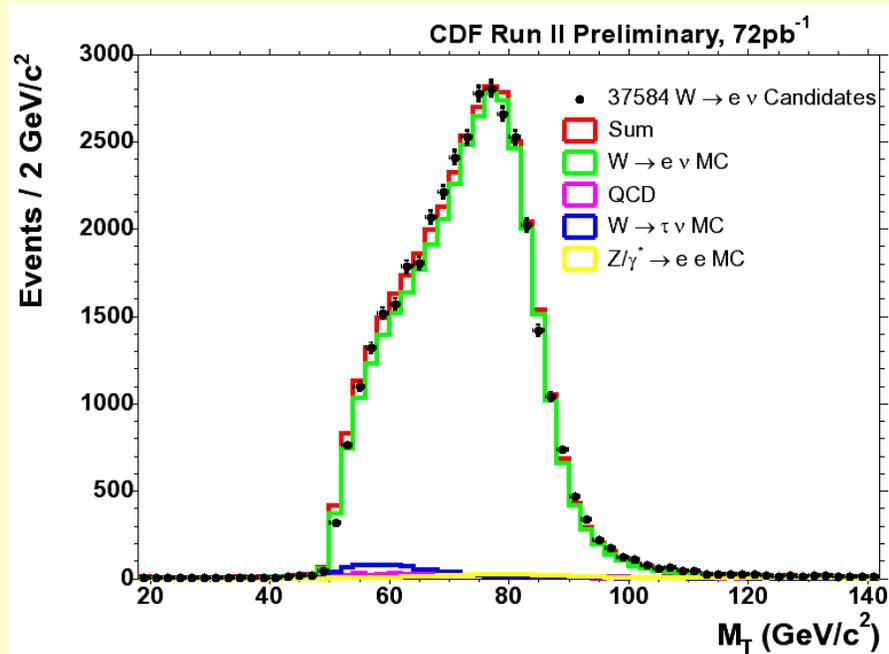
Z → ℓℓ cross sections



Good agreement with NNLO QCD calculations, QCD corrections are large: factor 1.3-1.4
 C.R.Hamberg et al, Nucl. Phys. B359 (1991) 343.

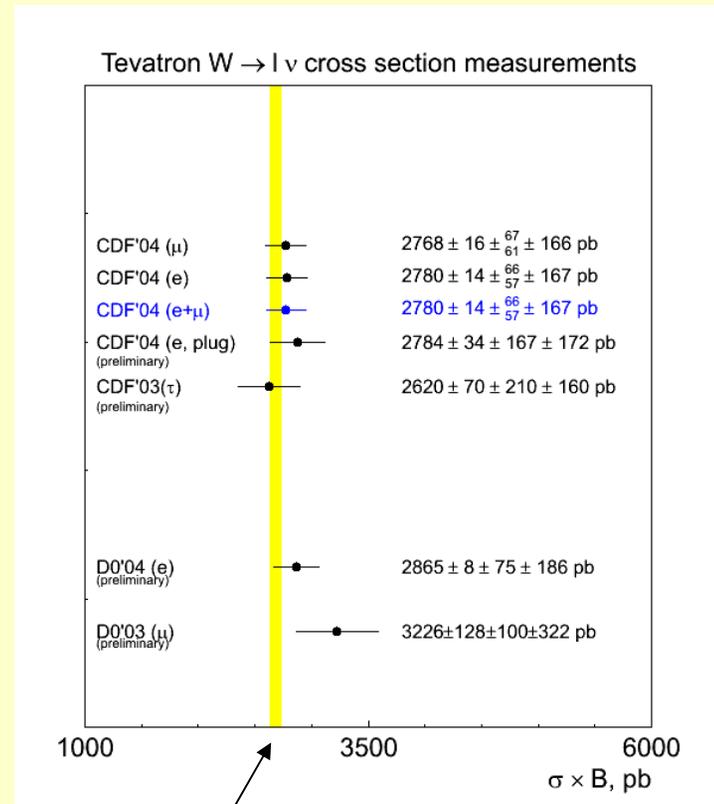
Precision is limited by systematic effects (uncertainties on luminosity, parton densities,...)

W → ℓν Cross Section



$$M_W^T = \sqrt{2 \cdot P_T^l \cdot P_T^\nu \cdot (1 - \cos \Delta\phi^{l,\nu})}$$

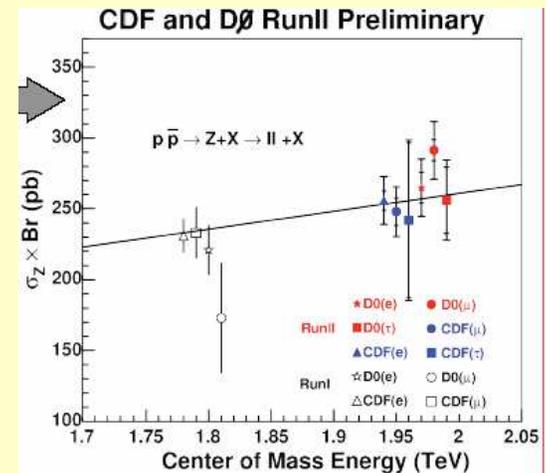
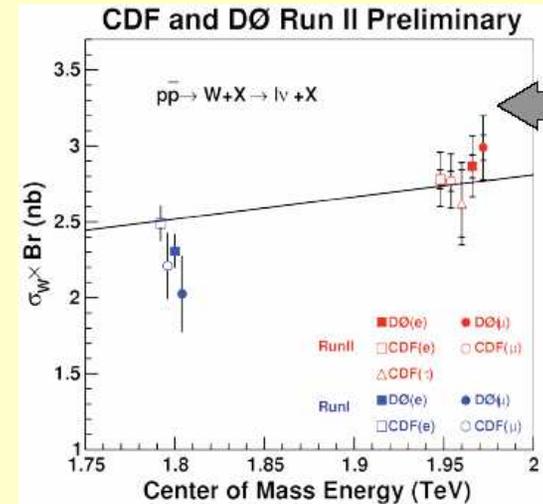
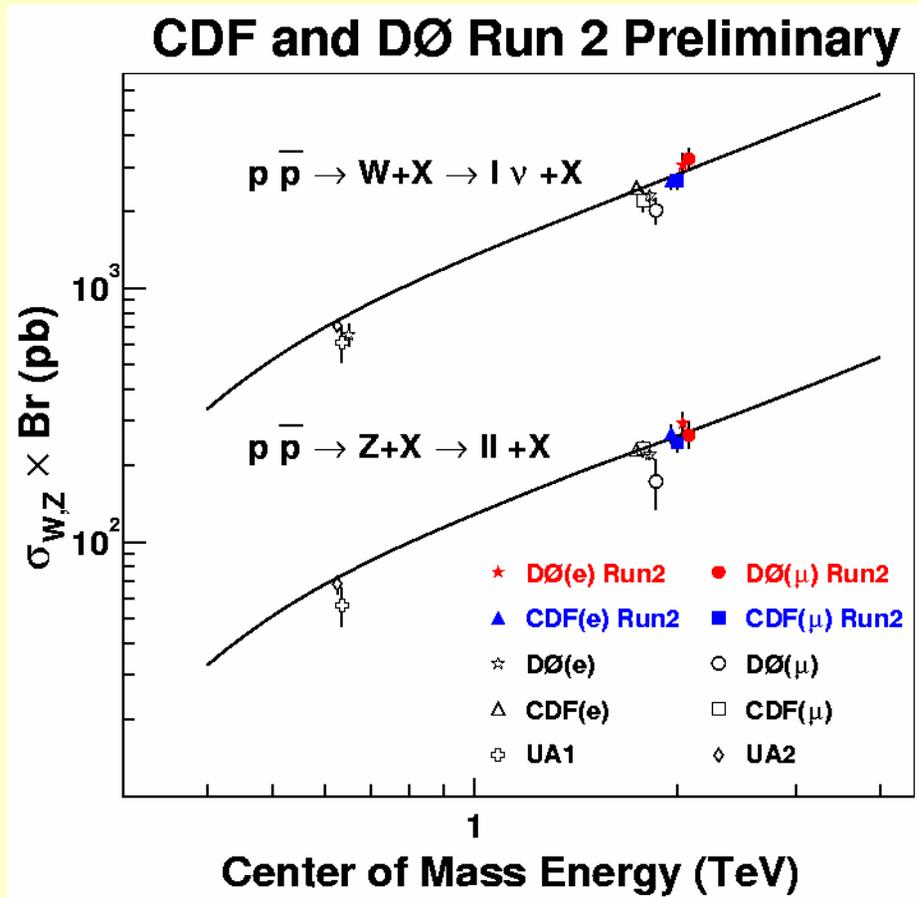
Note: the longitudinal component of the neutrino cannot be measured
 → only transverse mass can be reconstructed



Good agreement with NNLO QCD calculations
 C.R.Hamberg et al, Nucl. Phys. B359 (1991) 343.

Precision is limited by systematic effects (uncertainties on luminosity, parton densities,...)

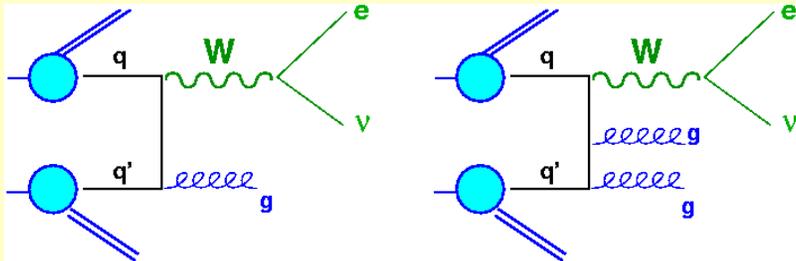
Comparison between measured W/Z cross sections and theoretical prediction (QCD)



C. R. Hamberg, W.L. van Neerven and T. Matsuura, Nucl. Phys. B359 (1991) 343

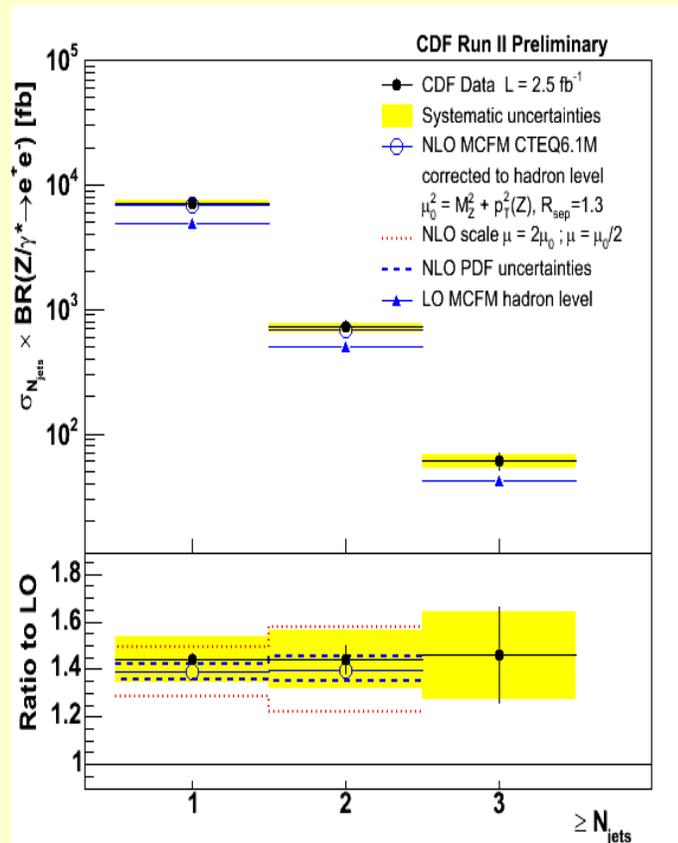


QCD Test in W/Z + jet production

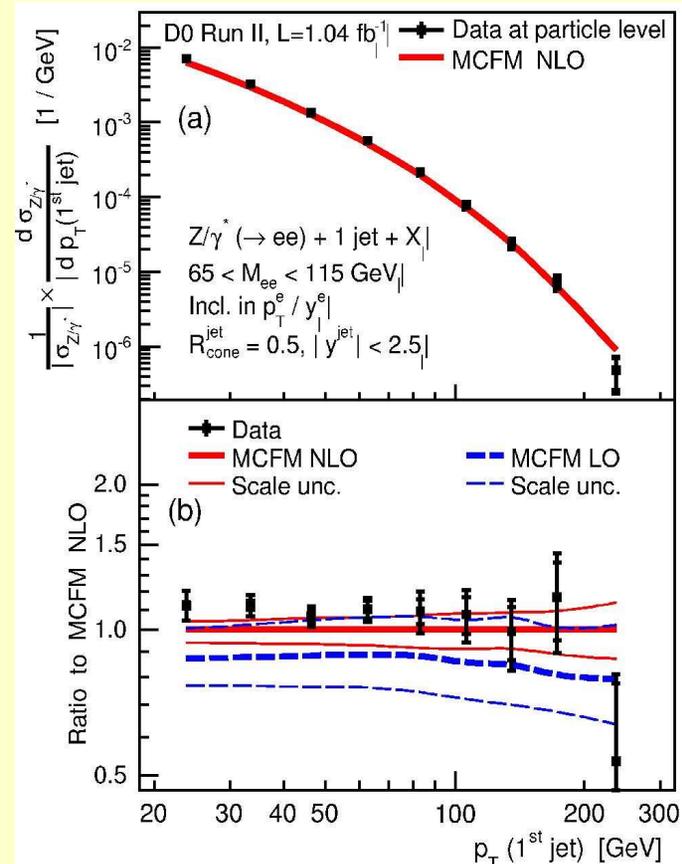


- LO predictions fail to describe the data;
- Jet multiplicities and p_T spectra in agreement with NLO predictions within errors;
- NLO central value $\sim 10\%$ low

Jet multiplicities in Z+jet production



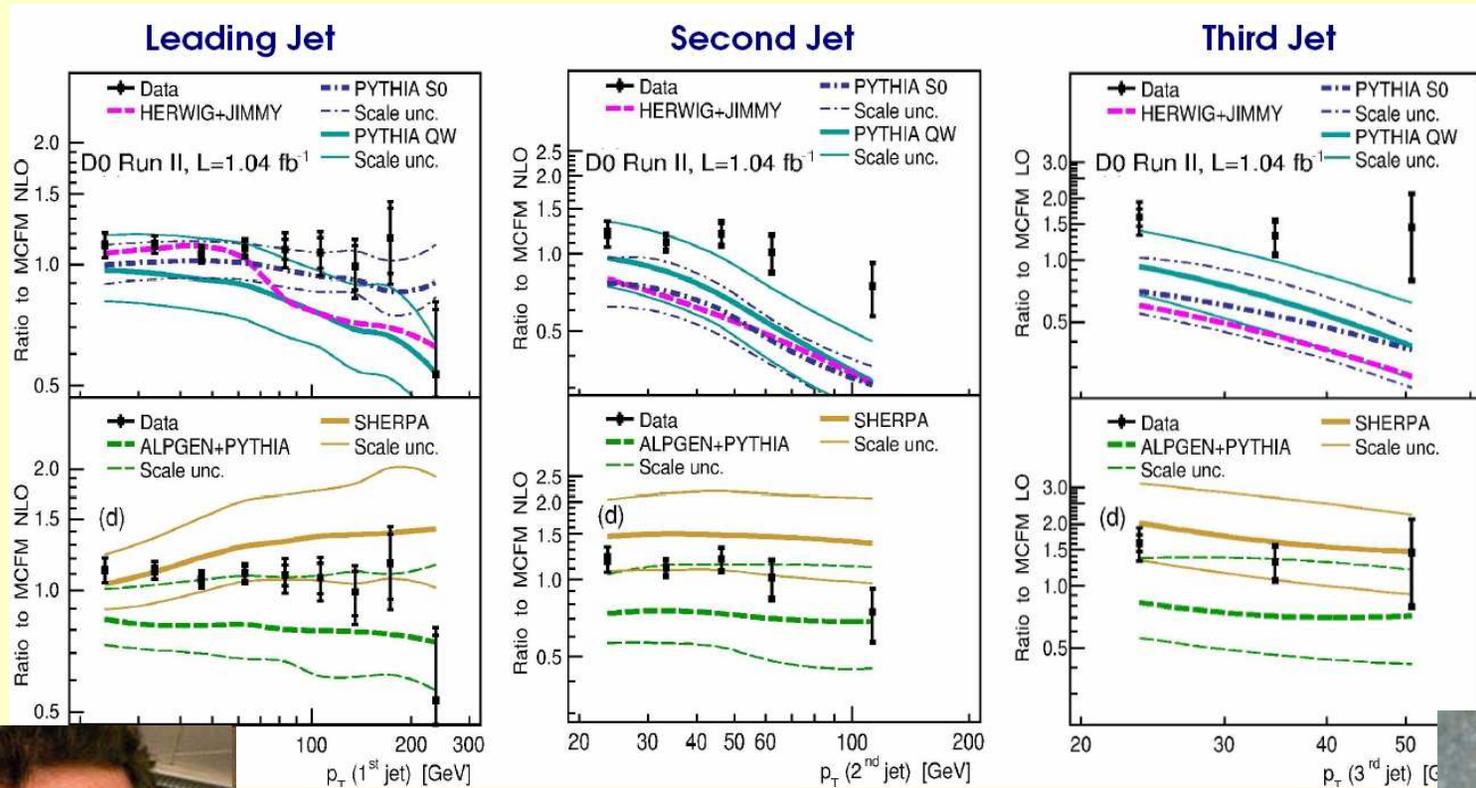
p_T spectrum of leading jet



comparison to different Monte Carlo predictions



- Comparison of p_T spectra of leading, second and third jet in Z+jet events to
 - PYTHIA and HERWIG (parton shower based Monte Carlos)
 - ALPGEN and SHERPA (explicit matrix elements (tree level) matched to parton showers)



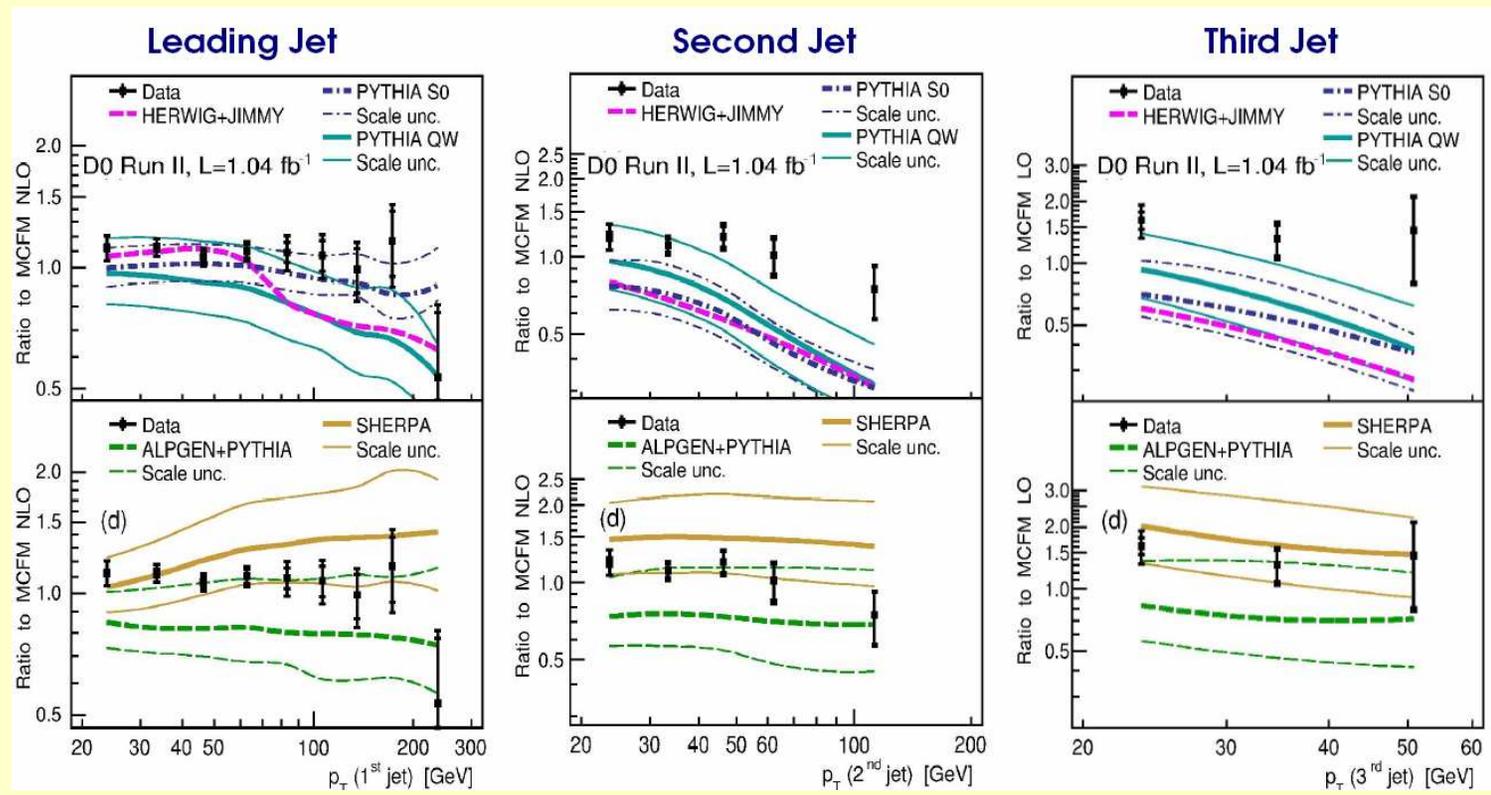
....they might have to try harder



comparison to different Monte Carlo predictions



- Comparison of p_T spectra of leading, second and third jet in Z+jet events to
 - PYTHIA and HERWIG (parton shower based Monte Carlos)
 - ALPGEN and SHERPA (explicit matrix elements (tree level) matched to parton showers)



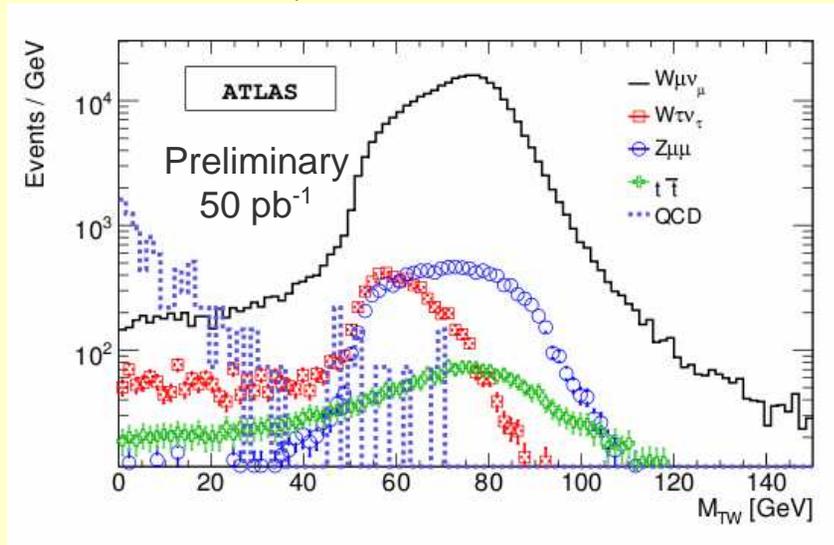
- **Conclusions:** (important for LHC)
 - Parton shower Monte Carlos fail to describe the higher jet p_T spectra;
 - Better agreement for ALPGEN and SHERPA, parameters can be tuned to describe them, but uncertainties -linked to the underlying tree level calculations- remain large;
 - It would be desirable to have NLO matched calculations

W and Z Cross sections at the LHC

Even with early data ($10\text{-}50 \text{ pb}^{-1}$),
high statistics of W and Z samples

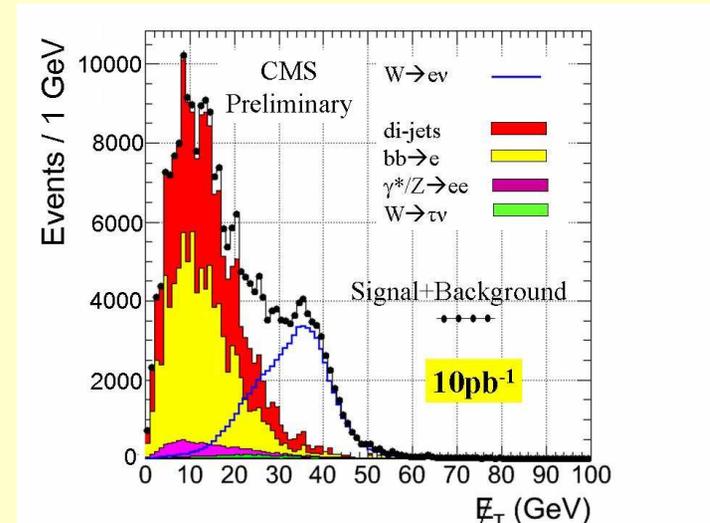
→ data-driven cross-section measurements

W → μ ν

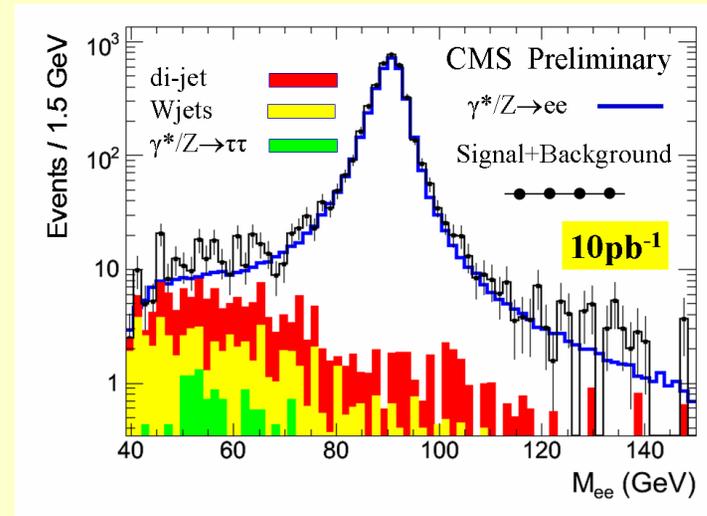


Limited by luminosity error: ~ 5-10% in first year,
Longer term goal ~ 2-3%
(process might be used later for luminosity measurement)

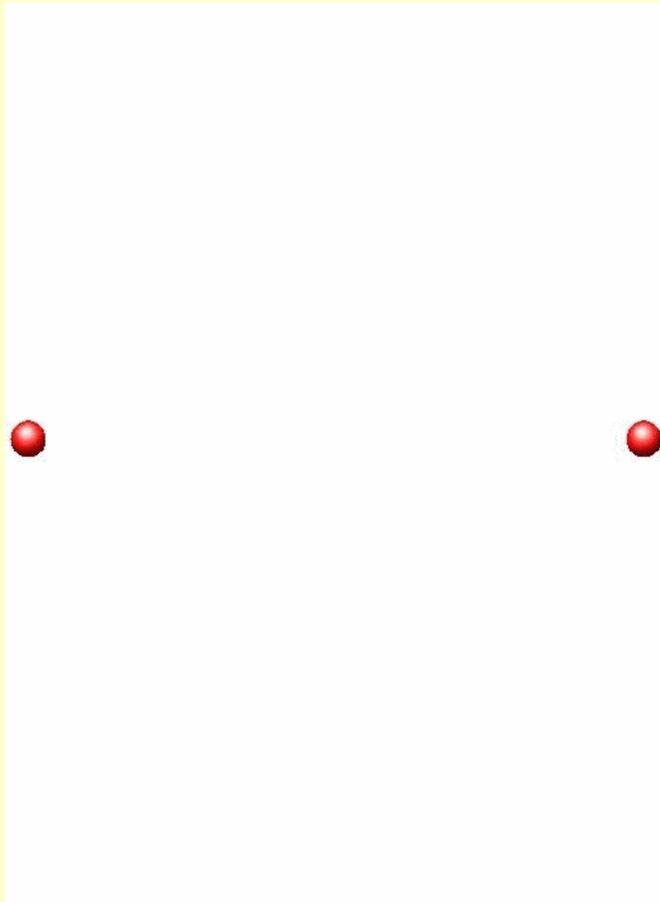
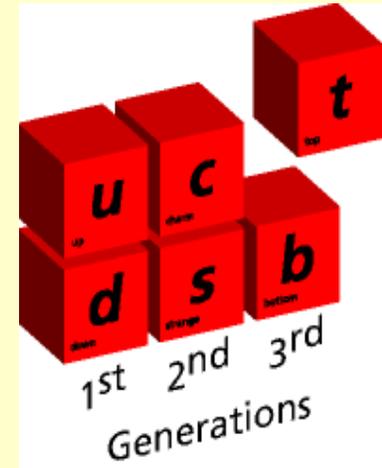
W → e ν



Z → ee

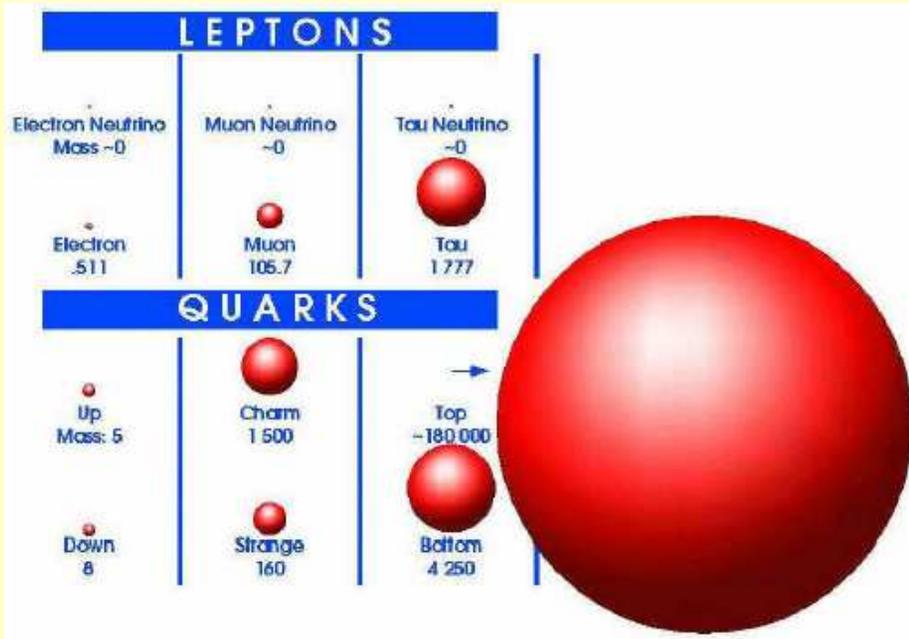


Top Quark Physics



- Discovered by CDF and DØ collaborations at the Tevatron in 1995
- Run I top physics results are consistent with the Standard Model
(Errors dominated by statistics)
- Run II top physics program will take full advantage of higher statistics
 - Better precision
 - Search for deviations from Standard Model expectations

Why is Top-Quark so important ?



The top quark may serve as a window to **New Physics** related to the electroweak symmetry breaking;

Why is its Yukawa coupling ~ 1 ??

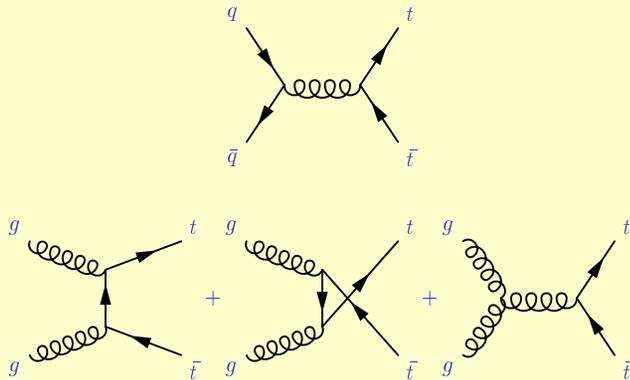
$$M_t = \frac{1}{\sqrt{2}} \lambda_t v$$

$$\Rightarrow \lambda_t = \frac{M_t}{173.9 \text{ GeV} / c^2}$$

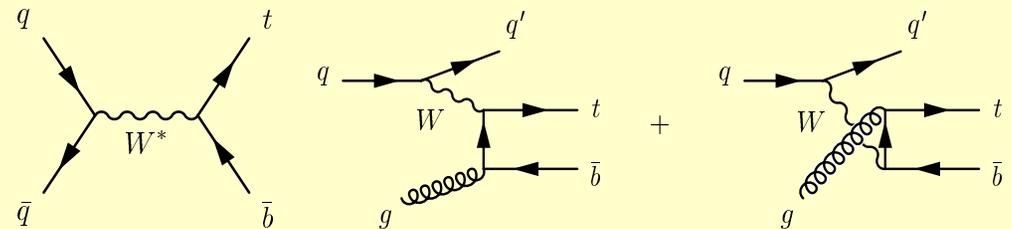
- We still know little about the properties of the top quark: mass, spin, charge, lifetime, decay properties (rare decays), gauge couplings, Yukawa coupling,...
- A unique quark: decays before it hadronizes, lifetime $\sim 10^{-24}$ s
no “toponium states”
remember: bb, bd, bs..... cc, cs..... Bound states (Mesons)

Top Quark Production

Pair production: qq and gg-fusion



Electroweak production of **single top-quarks** (Drell-Yan and Wg-fusion)



recently discovered by CDF and DØ at Fermilab

	Tevatron 1.96 TeV	LHC 14 TeV
qq	85%	5%
gg	15%	95%
σ (pb)	7 pb	830 pb

	Tevatron 1.96 TeV	LHC 14 TeV
σ (qq) (pb)	0.9	10
σ (gW) (pb)	2.4	250
σ (gb) (pb)	0.1	60

Top Quark Decays

BR ($t \rightarrow Wb$) $\sim 100\%$

Dilepton channel:

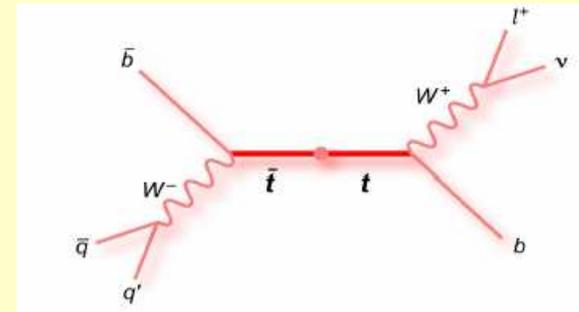
Both W 's decay via $W \rightarrow \ell\nu$ ($\ell=e$ or μ ; 4%)

Lepton + jet channel:

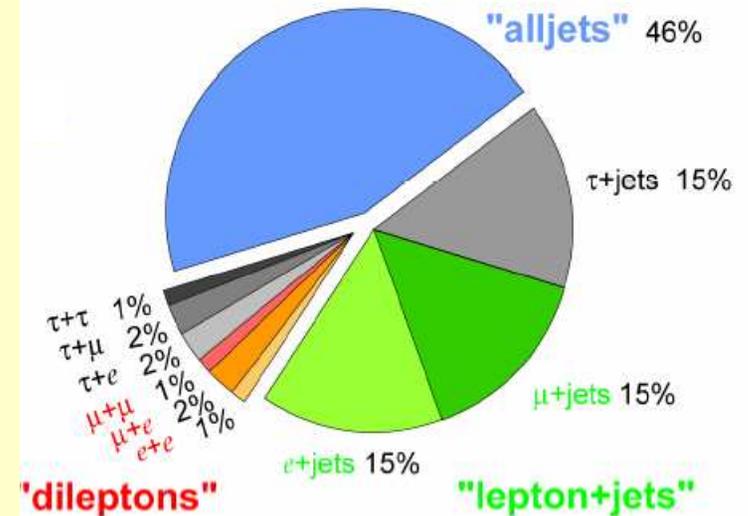
One W decays via $W \rightarrow \ell\nu$ ($\ell=e$ or μ ; 30%)

Full hadronic channel:

Both W 's decay via $W \rightarrow qq$ (46%)



Top Pair Branching Fractions

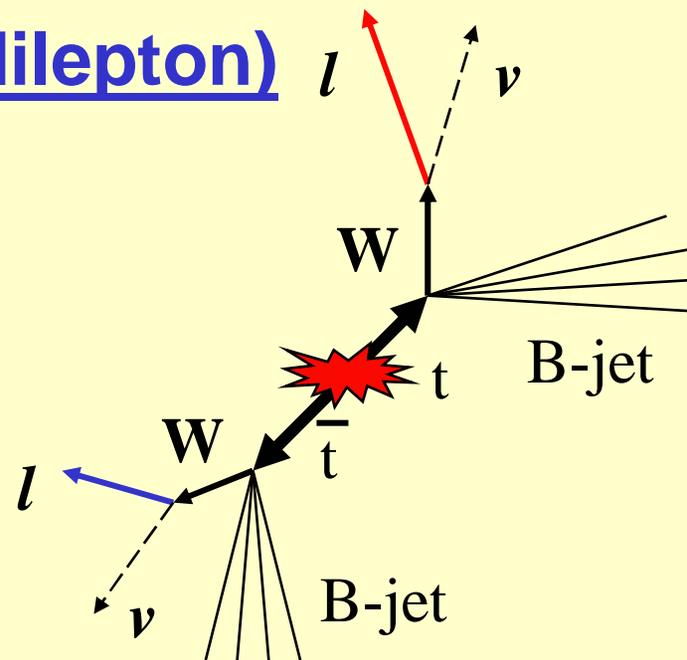


Important experimental signatures: - Lepton(s)

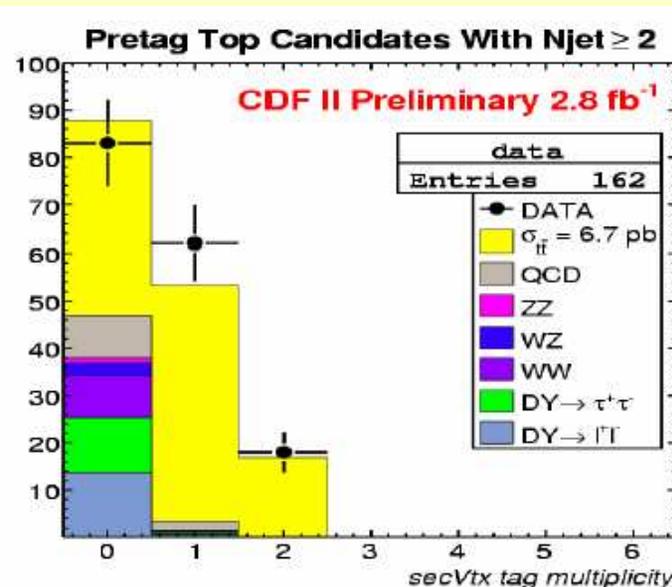
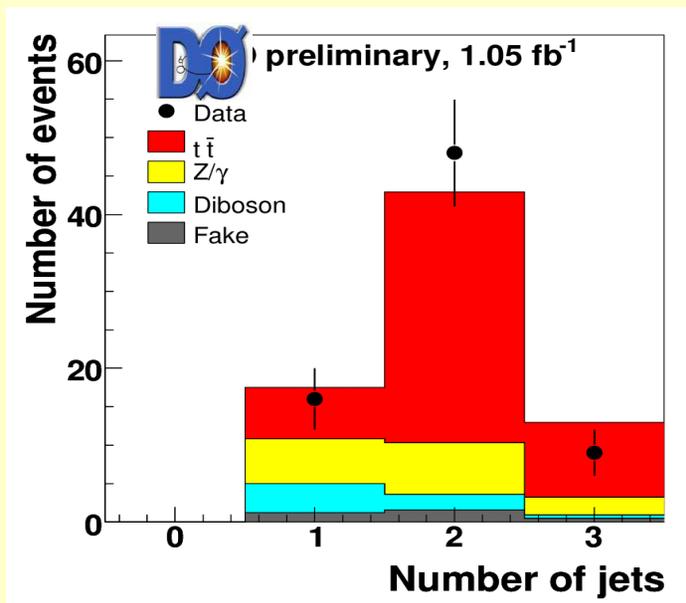
- Missing transverse momentum
- b-jet(s)

tt cross section (dilepton)

- Two high p_T leptons (opposite charge)
ee, e μ , $\mu\mu$
- Significant missing transverse momentum
- ≥ 1 jet (e μ), ≥ 2 jets (ee, $\mu\mu$)



ee, e μ and $\mu\mu$ combined

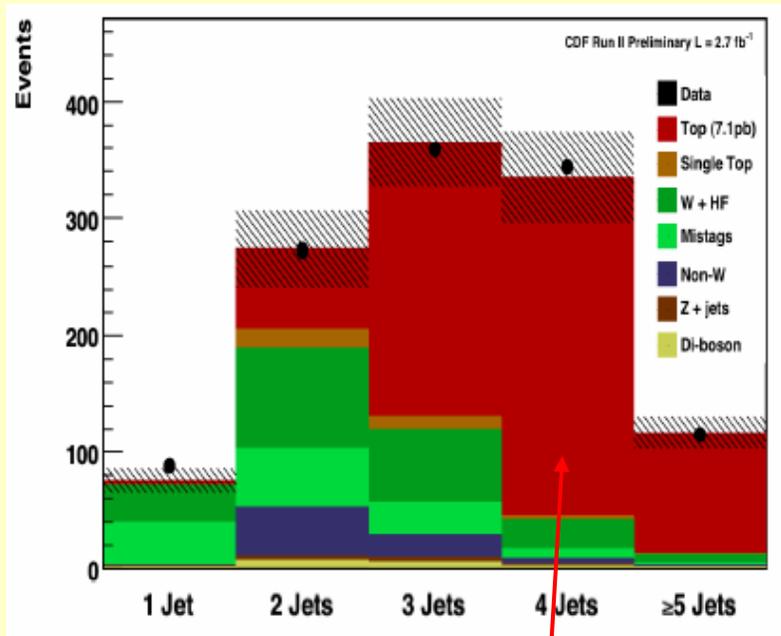


Top quark is needed to describe the b-jet multiplicity distribution in dilepton events

tt cross section (lepton + jets) (including b-tagging)

b-tag selection:

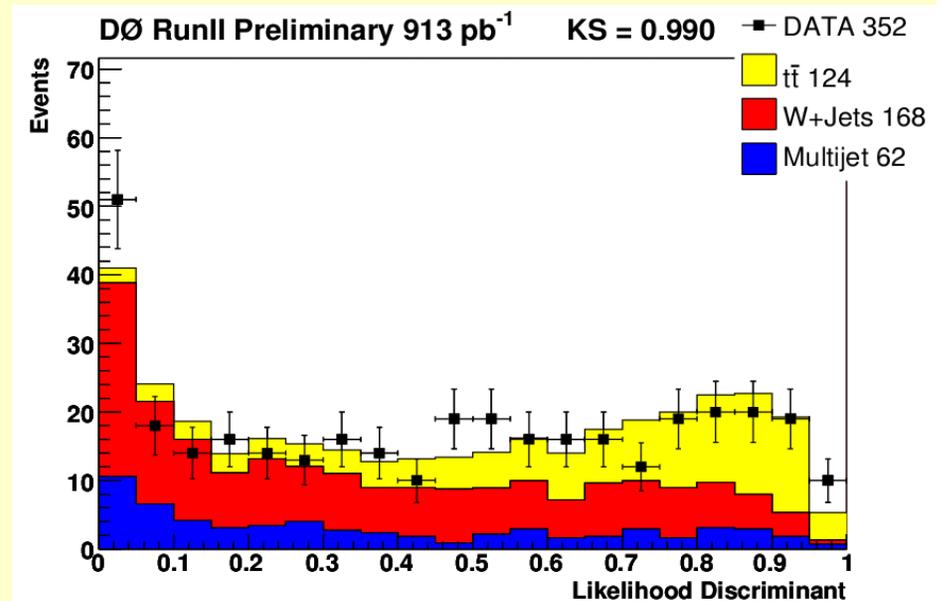
- One high P_T lepton (e, μ)
- Significant E_T^{miss}
- ≥ 1 b-tagged jet



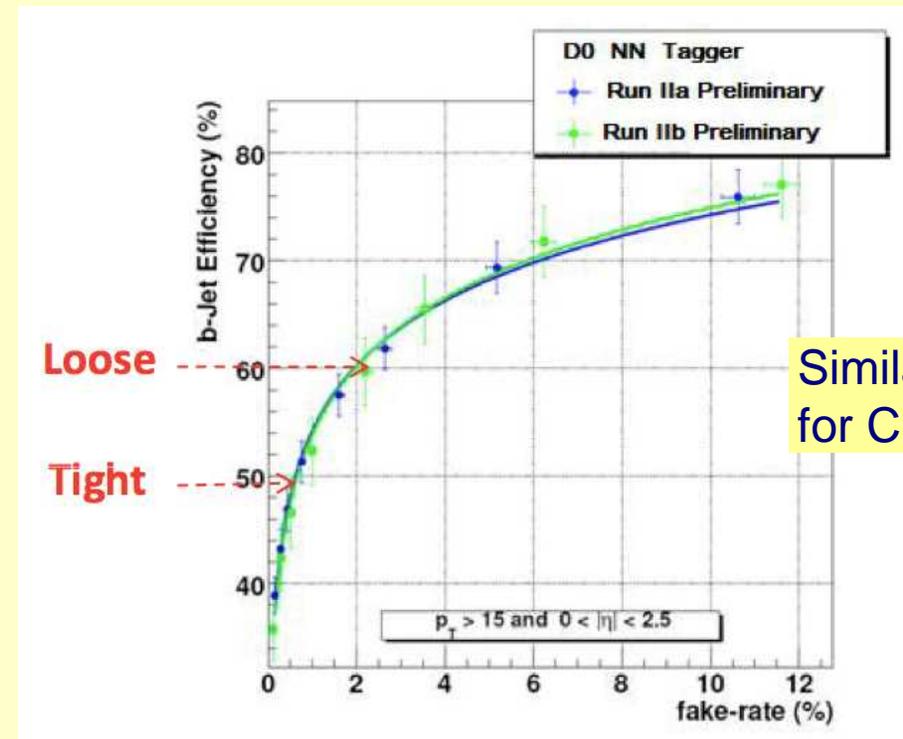
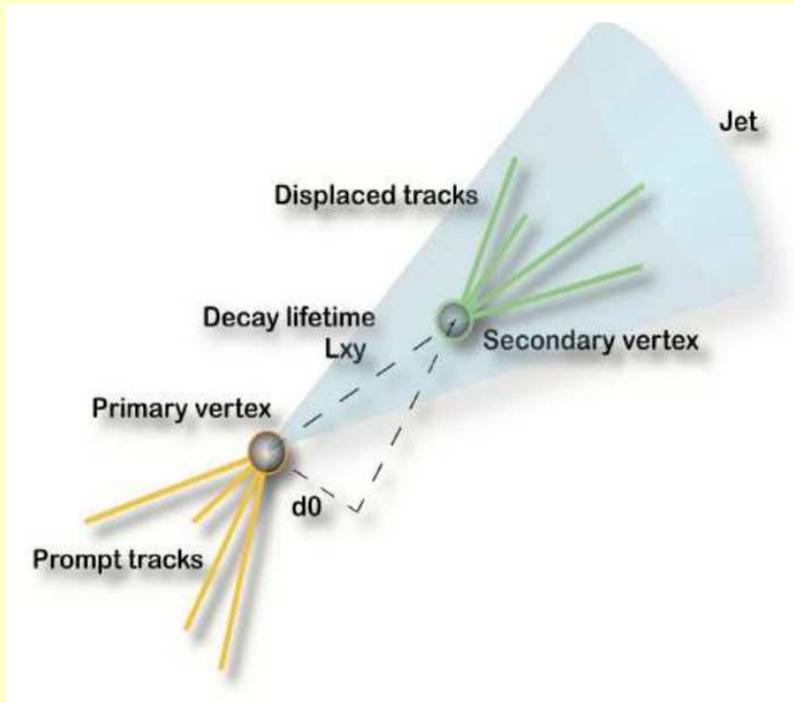
Clear excess above the W+ jet background in events with high jet multiplicity

Kinematic selection:

- One high P_T lepton (e, μ)
- Significant E_T^{miss}
- ≥ 4 jets
- **Likelihood discriminant (tt vs. W+jets)**

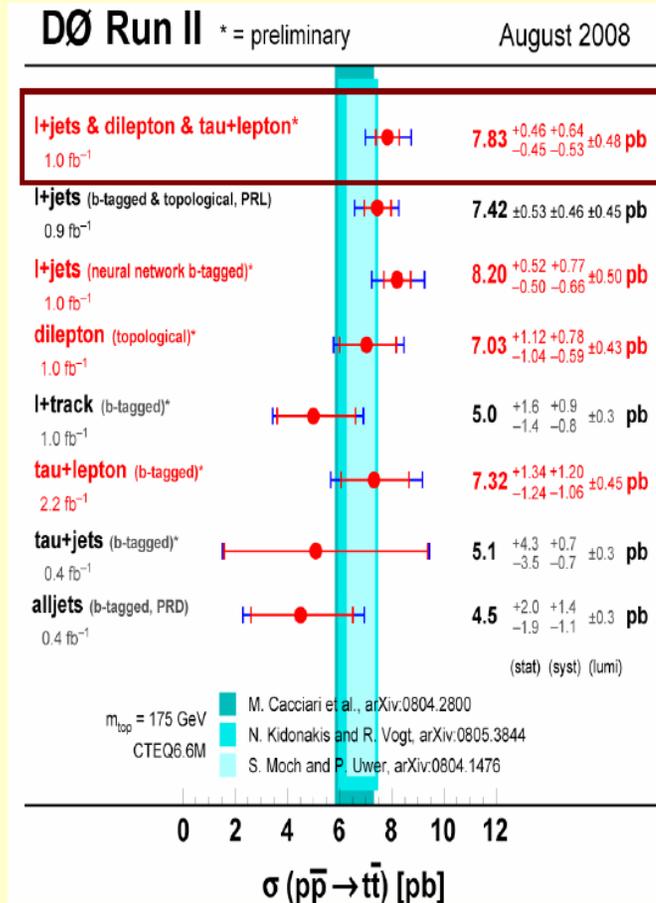
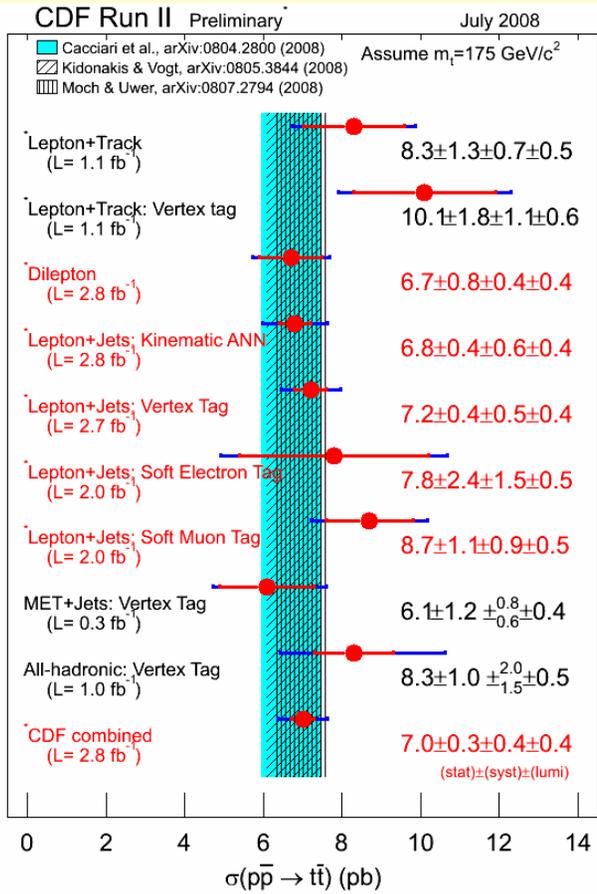


Tevatron b-tagging performance



Neural networks are used for optimal combination of tagging information

tt cross section summary (preliminary)



Summary of syst. uncertainties

b-tag analysis (2.7 fb⁻¹):

SYSTEMATIC	$\Delta \sigma$ pb	$\Delta \sigma / \sigma$ %
JET ENERGY SCALE	0.16	2.2
BOTTOM TAGGING	0.38	5.2
CHARM TAGGING	0.08	1.1
MIS-TAGS	0.15	2.1
HEAVY FLAVOR CORRECTION	0.23	3.2
LUMINOSITY	0.42	5.8
QCD FRACTION	0.02	0.2
PARTON SHOWER MODELING	0.13	1.8
INITIAL/FINAL STATE RADIATION	0.04	0.6
TRIGGER EFFICIENCY	0.05	0.6
PDF	0.06	1.0
TOTAL	0.67	9.3

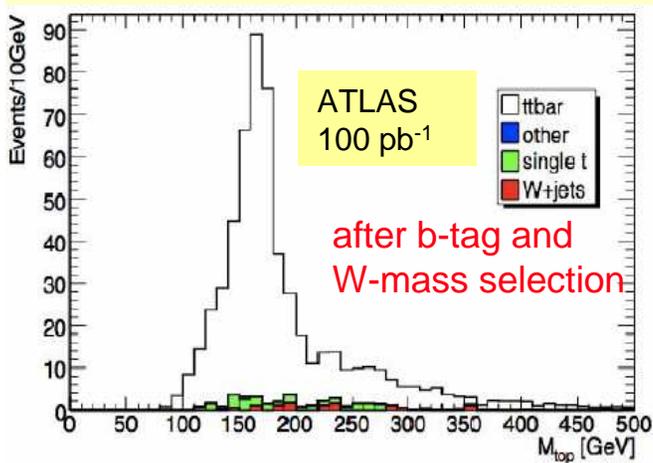
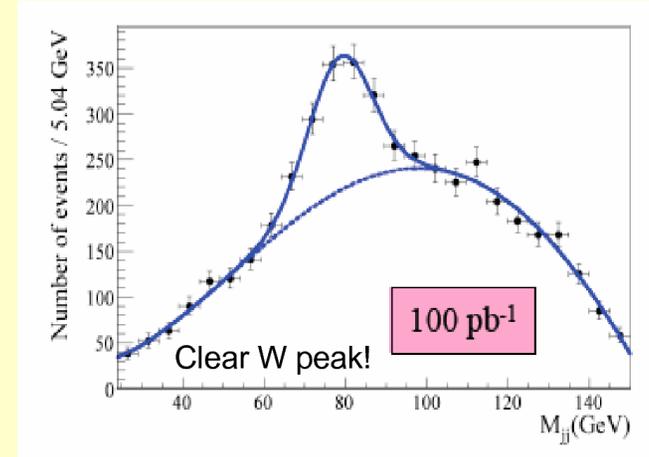
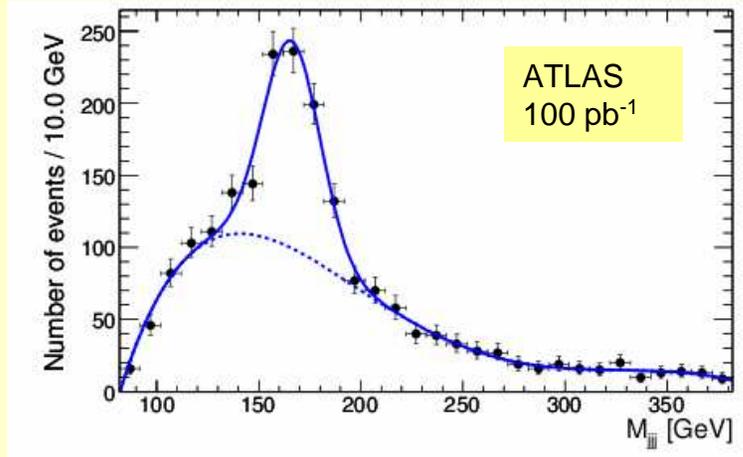
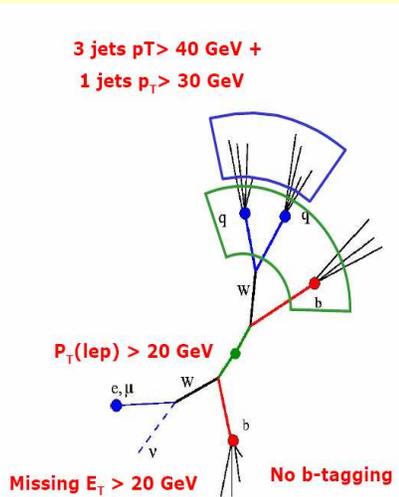
Good agreement:

- among various exp. measurements (two experiments)
- and with NLO + LL QCD prediction
- Systematic uncertainties at the 10% level (luminosity, b-tagging)

Top cross section in early LHC data

Large cross section: ~ 830 pb at $\sqrt{s} = 14$ TeV

Reconstructed mass distribution after a simple selection of $tt \rightarrow Wb Wb \rightarrow \ell\nu b qqb$ decays:



- Cross section measurement (test of perturbative QCD) with data corresponding to 100 pb^{-1} possible with an accuracy of $\pm 10\text{-}15\%$
- Errors are dominated by systematics (jet energy scale, Monte Carlo modelling (ISR, FSR),...)
- **Ultimate reach (100 fb^{-1}): $\pm 3\text{-}5\%$** (limited by uncertainty on the luminosity)

Electroweak parameters



- W mass
- Top Quark Mass & Properties
- Single top, V_{tb}

Precision measurements of m_W and m_{top}

Motivation:

W mass and top quark mass are **fundamental parameters** of the Standard Model;
 The standard theory provides well defined **relations between m_W , m_{top} and m_H**

Electromagnetic constant
 measured in atomic transitions,
 e^+e^- machines, etc.

$$m_W = \left(\frac{\pi \alpha_{EM}}{\sqrt{2} G_F} \right)^{1/2} \frac{1}{\sin \theta_W \sqrt{1 - \Delta r}}$$

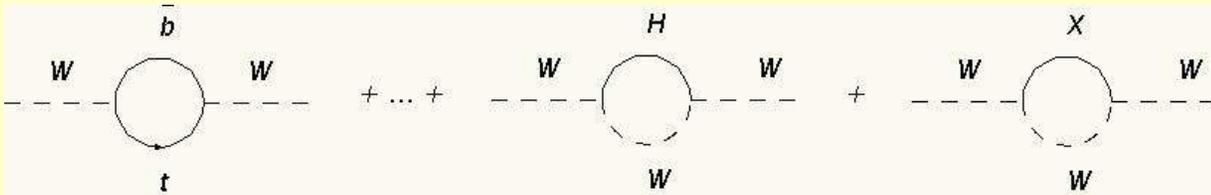
Fermi constant
 measured in muon
 decay

weak mixing angle
 measured at
 LEP/SLC

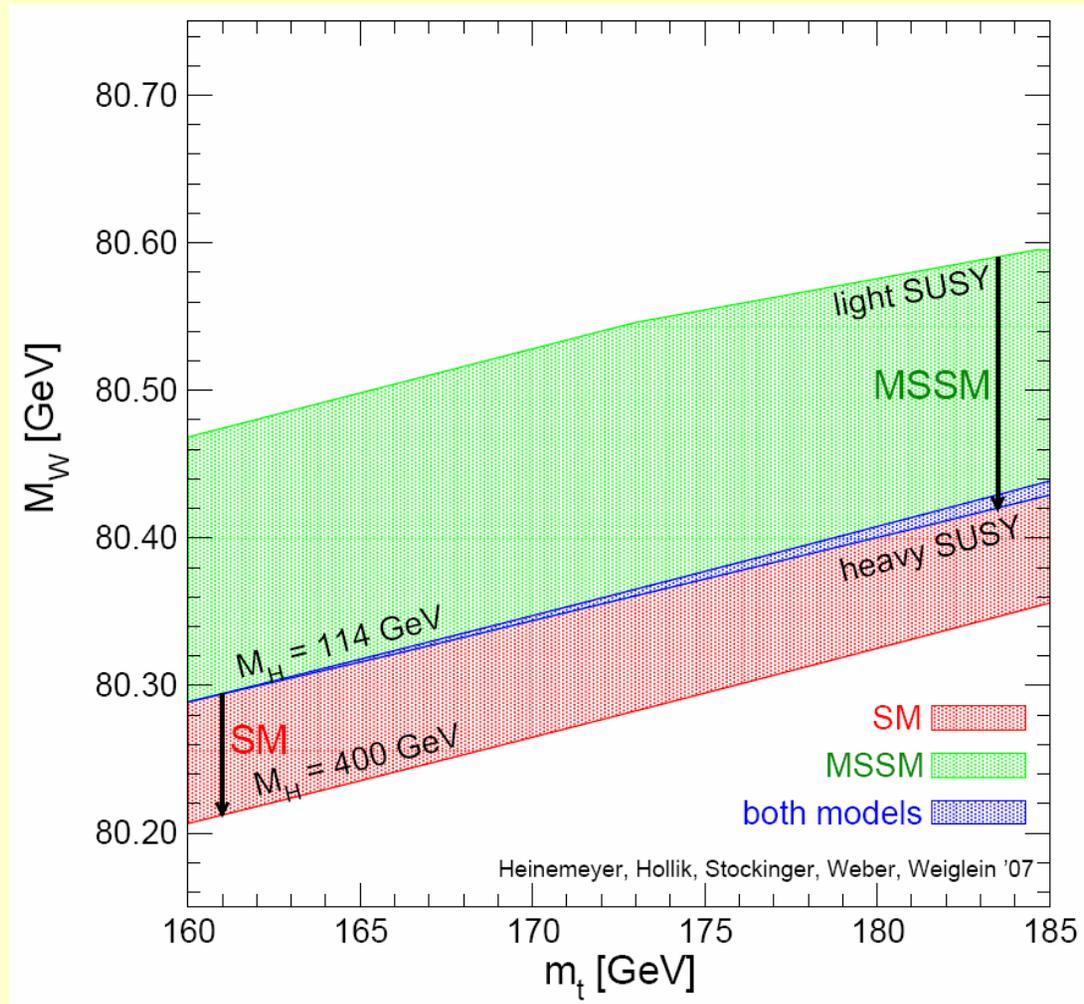
radiative corrections
 $\Delta r \sim f(m_{top}^2, \log m_H)$
 $\Delta r \approx 3\%$

$G_F, \alpha_{EM}, \sin \theta_W$
 are known with high precision

Precise measurements of the
 W mass and the top-quark
 mass constrain the Higgs-
 boson mass
 (and/or the theory,
 radiative corrections)



Relation between m_W , m_t , and m_H



The W-mass measurement

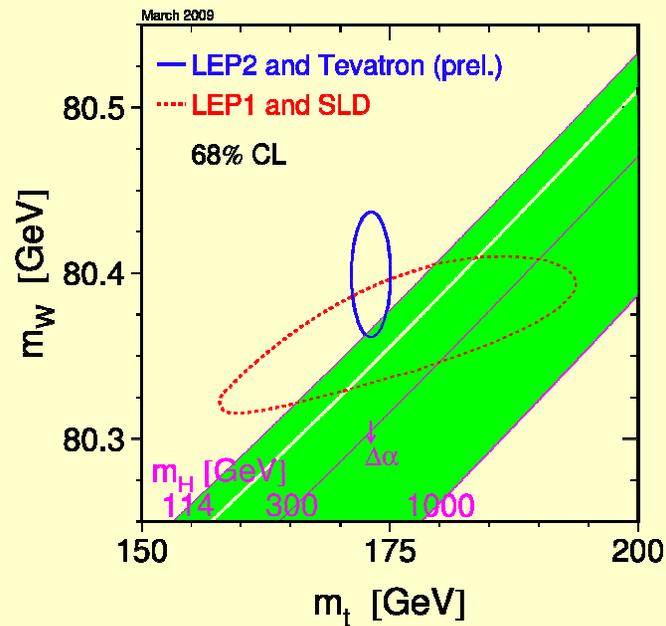
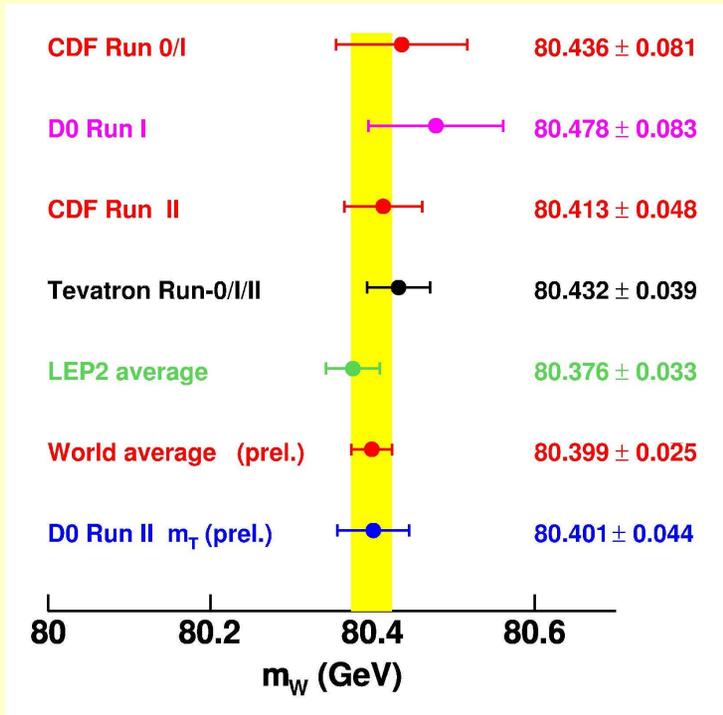
$$m_W = \left(\frac{\pi \alpha_{EM}}{\sqrt{2} G_F} \right)^{1/2} \frac{1}{\sin \theta_W \sqrt{1 - \Delta r}}$$

$3 \cdot 10^{-4}$

m_W (from LEP2 + Tevatron) = 80.399 ± 0.025 GeV

m_{top} (from Tevatron) = 173.1 ± 1.3 GeV

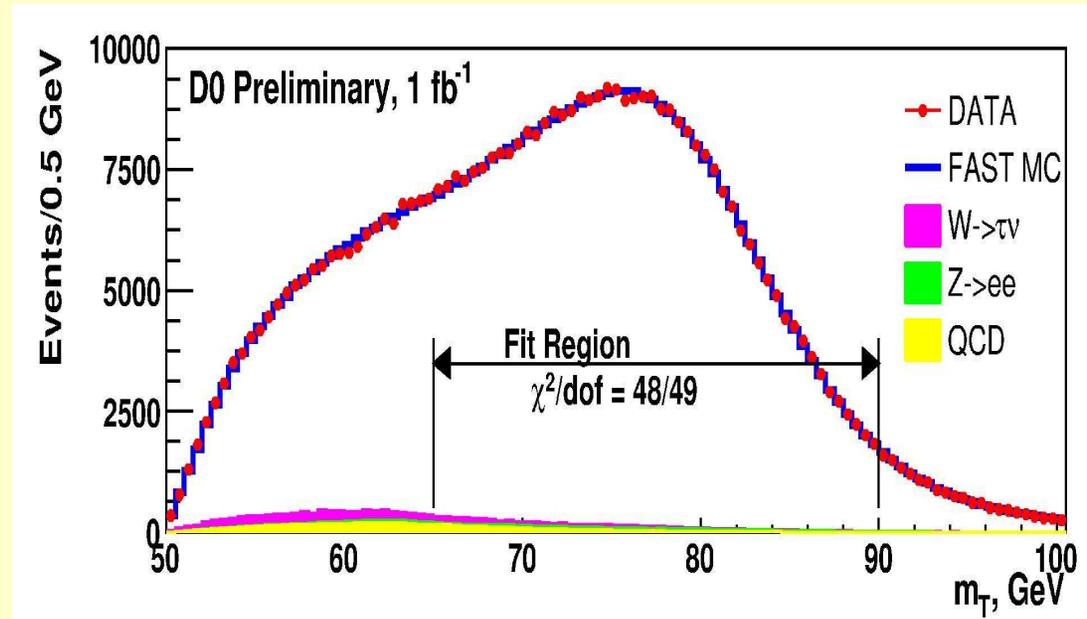
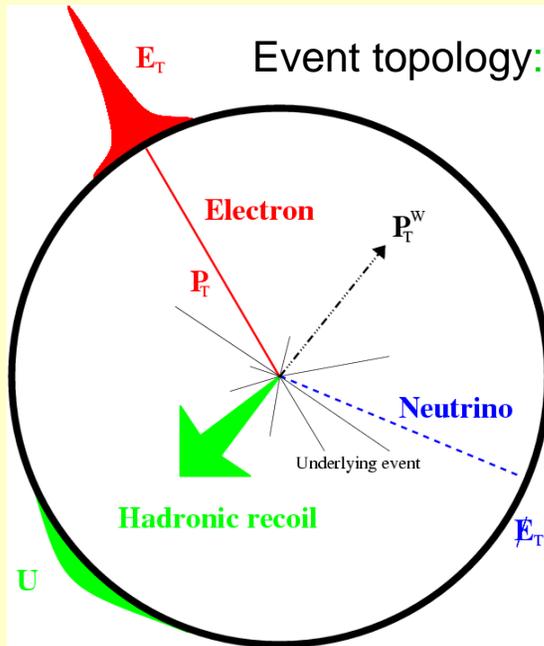
0.8%



A light Higgs boson is favoured by present measurements

Ultimate test of the Standard Model: comparison between the direct Higgs boson mass and predictions from radiative corrections....

Technique used for W mass measurement at hadron colliders:



Observables: $P_T(e)$, $P_T(\text{had})$

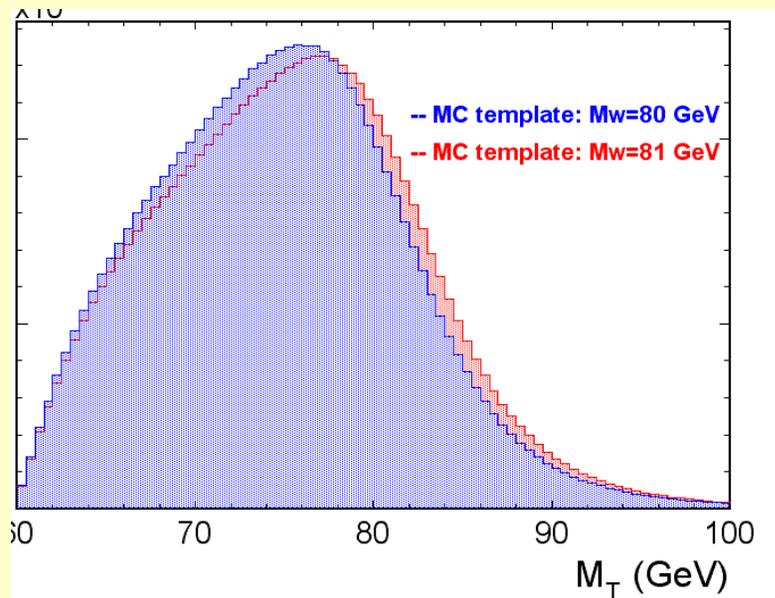
$$\Rightarrow P_T(\nu) = - (P_T(e) + P_T(\text{had}))$$

$$\Rightarrow M_W^T = \sqrt{2 \cdot P_T^l \cdot P_T^\nu \cdot (1 - \cos \Delta\phi^{l,\nu})}$$

long. component cannot be measured

In general the **transverse mass** M_T is used for the determination of the W-mass (smallest systematic uncertainty).

Shape of the transverse mass distribution is sensitive to m_W , the measured distribution is fitted with Monte Carlo predictions, where m_W is a parameter



Main uncertainties:

Ability of the Monte Carlo to reproduce real life:

- Detector performance
(energy resolution, energy scale,)
- Physics: production model
 $p_T(W)$, Γ_W ,
- Backgrounds

Systematic Uncertainties (Tevatron measurements)

CDF II : 200 pb⁻¹

m_T Uncertainty [MeV]	Electrons	Muons	Common
Lepton Scale	30	17	17
Lepton Resolution	9	3	0
Recoil Scale	9	9	9
Recoil Resolution	7	7	7
u_{ij} Efficiency	3	1	0
Lepton Removal	8	5	5
Backgrounds	8	9	0
$p_T(W)$	3	3	3
PDF	11	11	11
QED	11	12	11
Total Systematic	39	27	26
Statistical	48	54	0
Total	62	60	26

Total	48
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D0 Preliminary : 1 fb⁻¹

Source	$\sigma(m_W)$ MeV m_T
Experimental	
Electron Energy Scale	34
Electron Energy Resolution Model	2
Electron Energy Nonlinearity	4
W and Z Electron energy loss differences	4
Recoil Model	6
Electron Efficiencies	5
Backgrounds	2
Experimental Total	35
W production and decay model	
PDF	9
QED	7
Boson p_T	2
W model Total	12
Total	37

Statistical	23
Total	44

Dominant error: knowledge of the lepton energy scale of the detector !

What precision can be reached in Run II and at the LHC ?

Numbers for a
single decay
channel

$W \rightarrow e\nu$

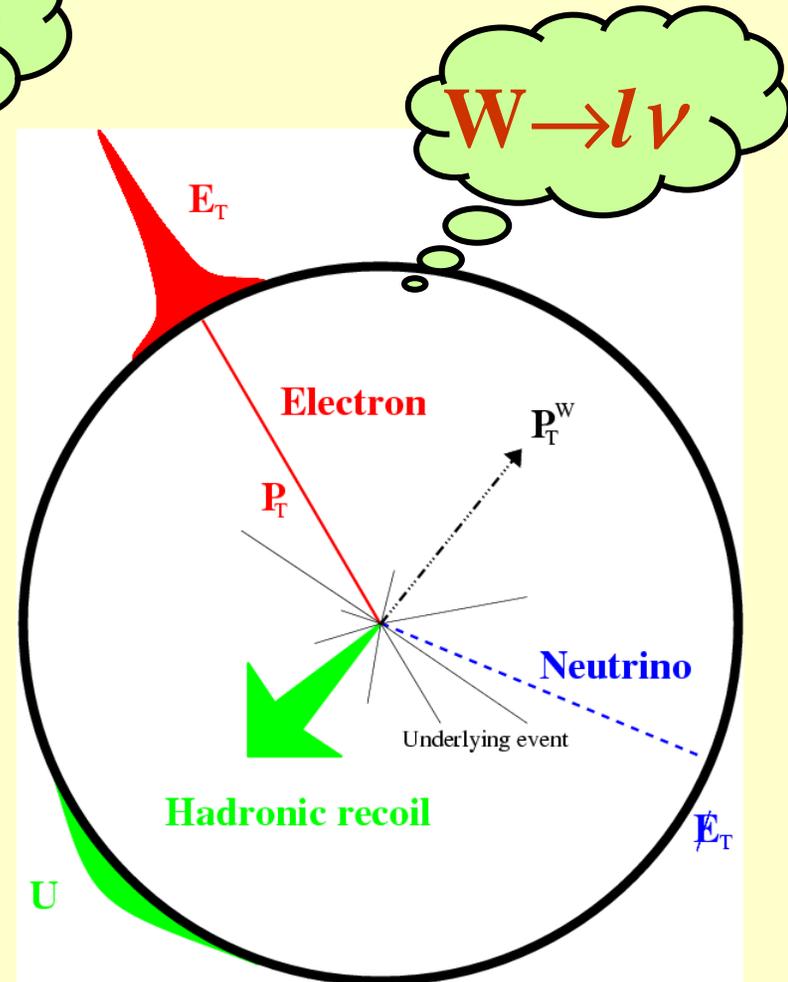
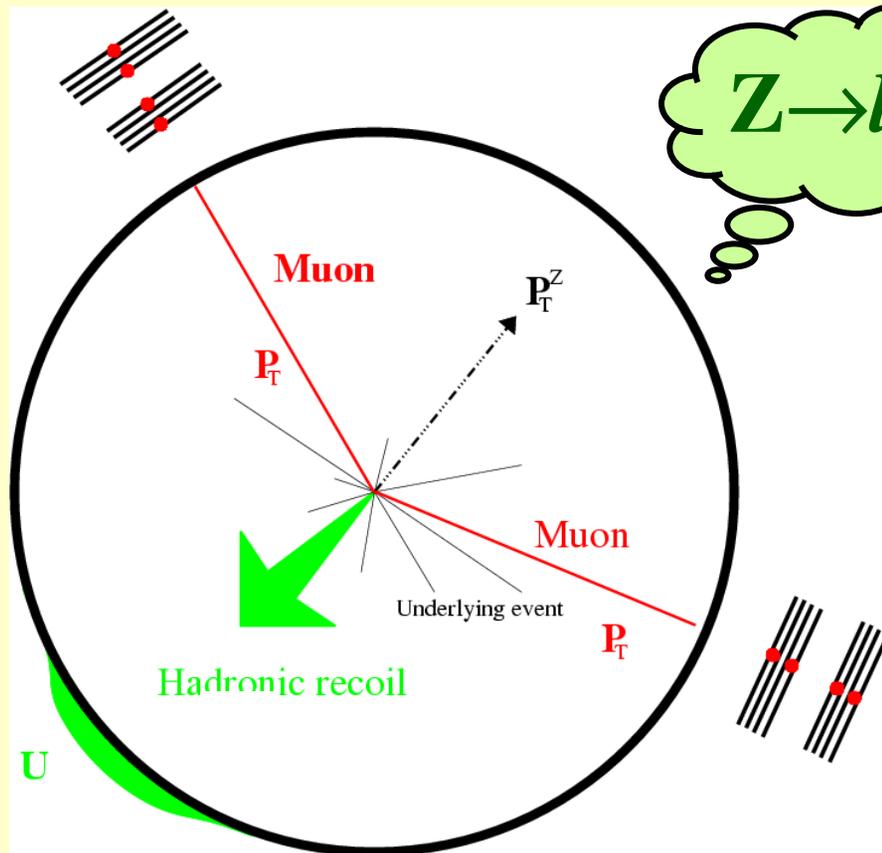
Int. Luminosity	CDF 0.2 fb ⁻¹	DØ 1 fb ⁻¹	LHC 10 fb ⁻¹
Stat. error	48 MeV	23 MeV	2 MeV
Energy scale, lepton res.	30 MeV	34 MeV	4 MeV
Monte Carlo model (P_T^W , structure functions, photon-radiation....)	16 MeV	12 MeV	7 MeV
Background	8 MeV	2 MeV	2 MeV
Tot. Syst. error	39 MeV	37 MeV	8 MeV
Total error	62 MeV	44 MeV	~10 MeV

- Tevatron numbers are based on real data analyses
- LHC numbers should be considered as „ambitious goal“
 - Many systematic uncertainties can be controlled in situ, using the large $Z \rightarrow \ell\ell$ sample ($p_T(W)$, recoil model, resolution)
 - Lepton energy scale of $\pm 0.02\%$ has to be achieved to reach the quoted numbers

Combining both experiments (ATLAS + CMS, 10 fb⁻¹), both lepton species and assuming a scale uncertainty of $\pm 0.02\%$ a total error in the order of

$\Rightarrow \Delta m_W \sim \pm 10 - 15 \text{ MeV}$ might be reached.

Signature of Z and W decays



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Tot. Syst. error	39 MeV	37 MeV	8 MeV
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 - Many systematic uncertainties can be controlled in situ, using the large $Z \rightarrow \ell\ell$ sample (PT(W), recoil model, resolution)
 - Lepton energy scale of $\pm 0.02\%$ has to be achieved to reach the quoted numbers

Combining both experiments (ATLAS + CMS, 10 fb⁻¹), both lepton species and assuming a scale uncertainty of $\pm 0.02\%$ a total error in the order of

$\Rightarrow \Delta m_W \sim \pm 10 - 15 \text{ MeV}$ might be reached.

Top mass measurements

- Top mass determination:
No simple mass reconstruction possible,
Monte Carlo models needed

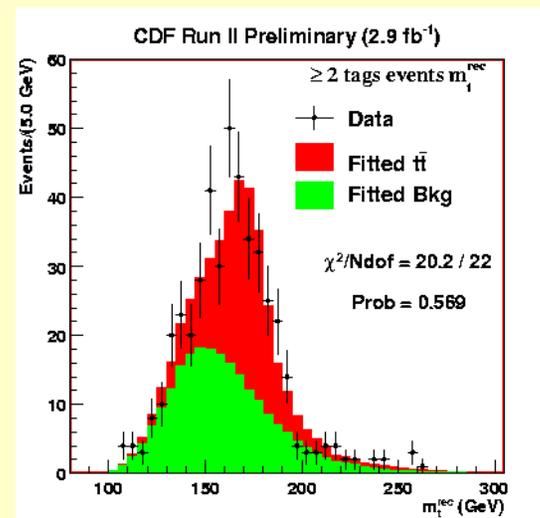
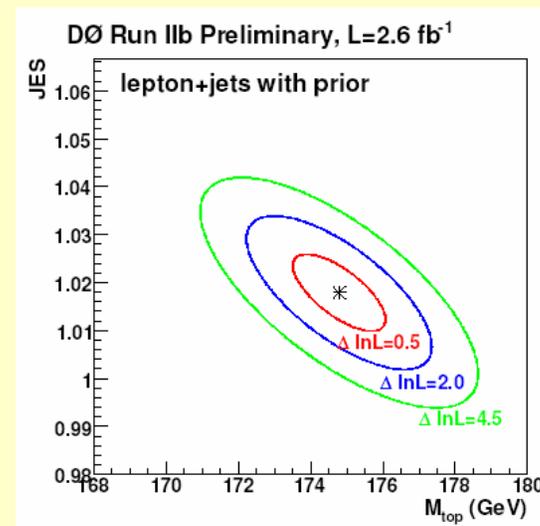
→ **template methods**,...
matrix element method...

Most precise single measurements:

$$m_{\text{top}} = 172.1 \pm 0.9 \text{ (stat)} \pm 1.3 \text{ (syst)} \text{ GeV}/c^2 \quad (\text{CDF})$$

$$m_{\text{top}} = 173.7 \pm 0.8 \text{ (stat)} \pm 1.6 \text{ (syst)} \text{ GeV}/c^2 \quad (\text{D}\emptyset)$$

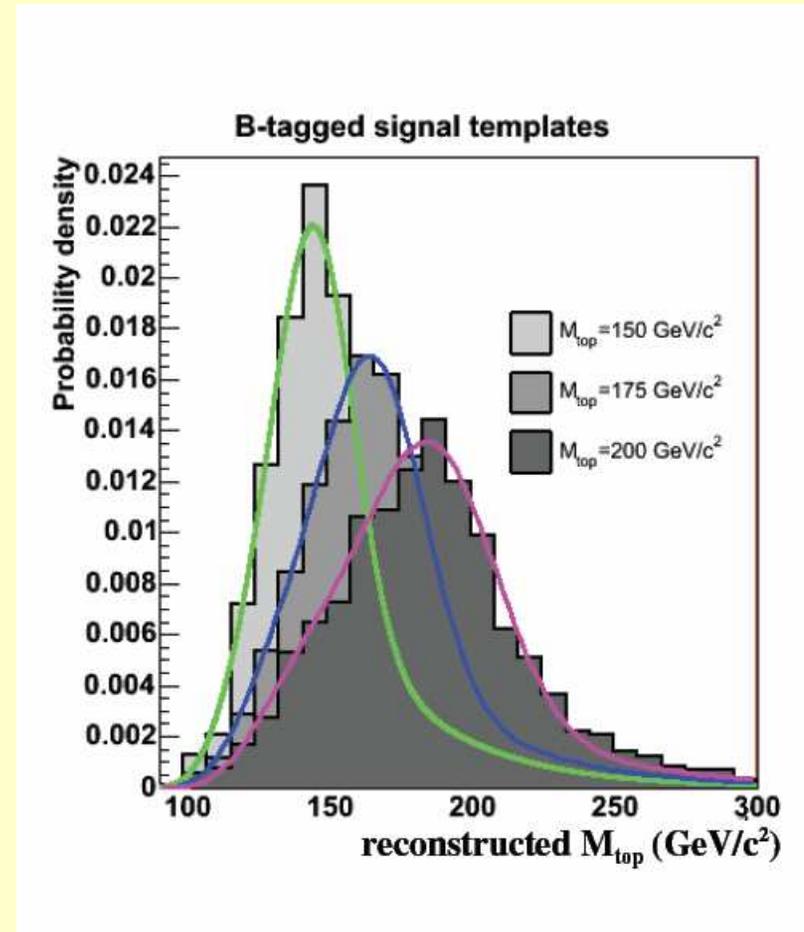
- Reduce jet energy scale systematic by using in-situ hadronic W mass in tt events
(simultaneous determination of m_t and energy scale)



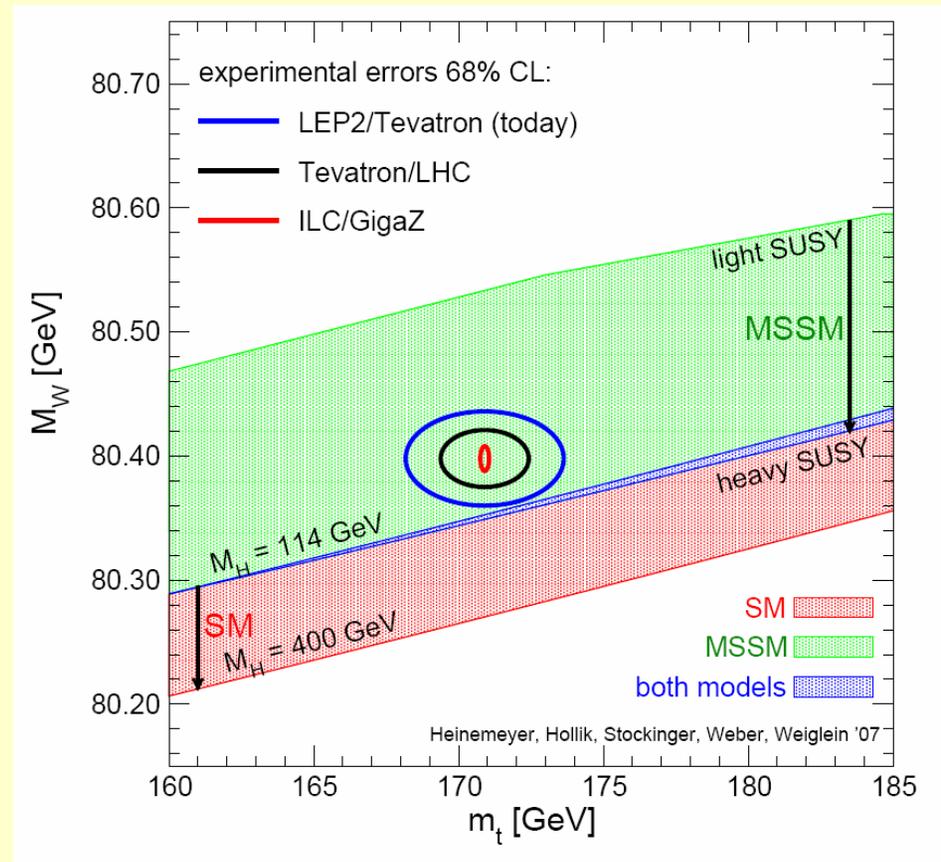
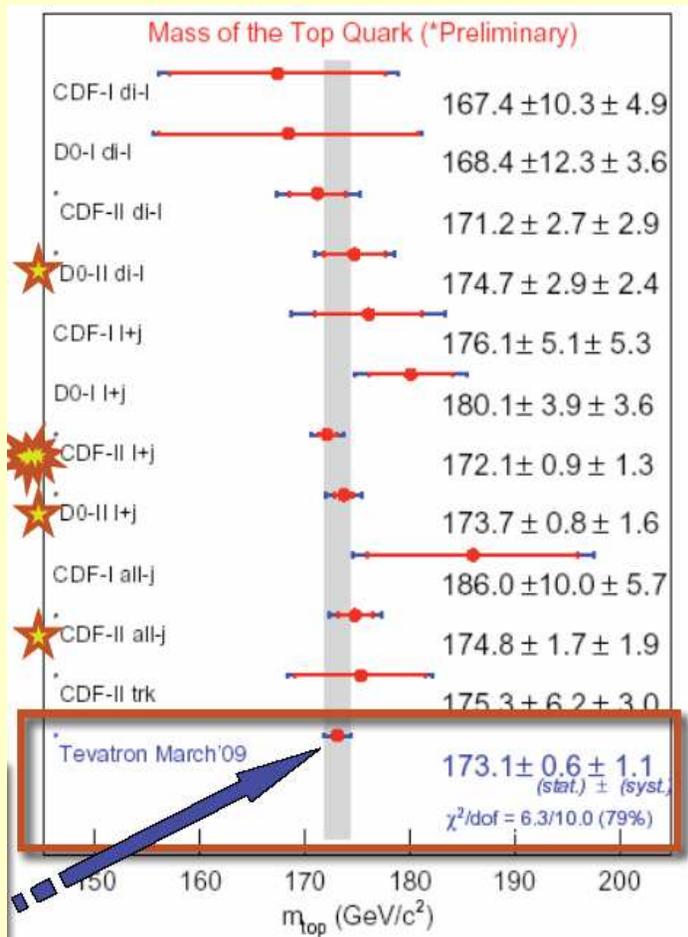
full hadronic channel

Example: template method

- Calculate a per-event observable that is sensitive to m_t
- Make templates from signal and background events
- Use pseudo-experiments (Monte Carlo) to check that method works
- Fit data to templates using maximum likelihood method



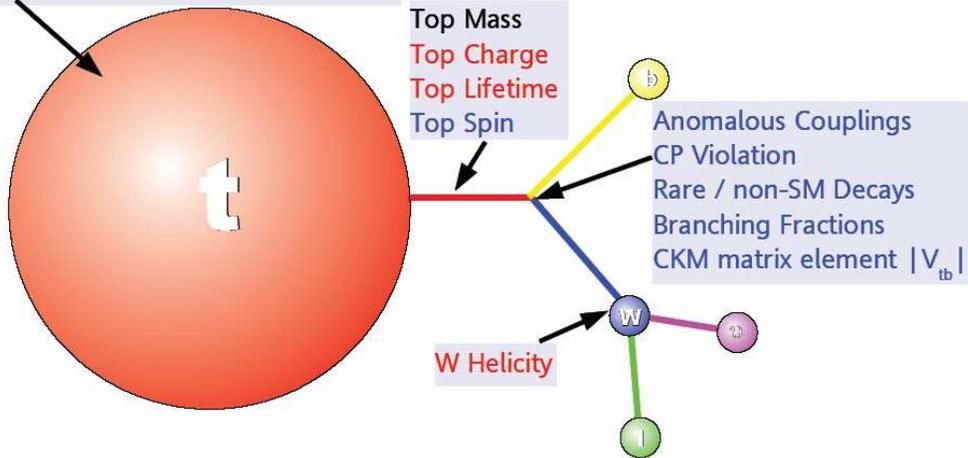
Future Prospects for the top quark mass measurement



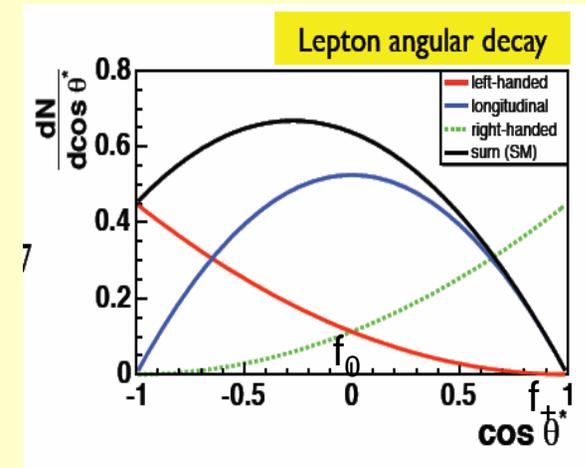
Expected LHC precision for 10 fb⁻¹:
(Combination of several methods, maybe somewhat conservative)

< ~ 1 GeV/c²

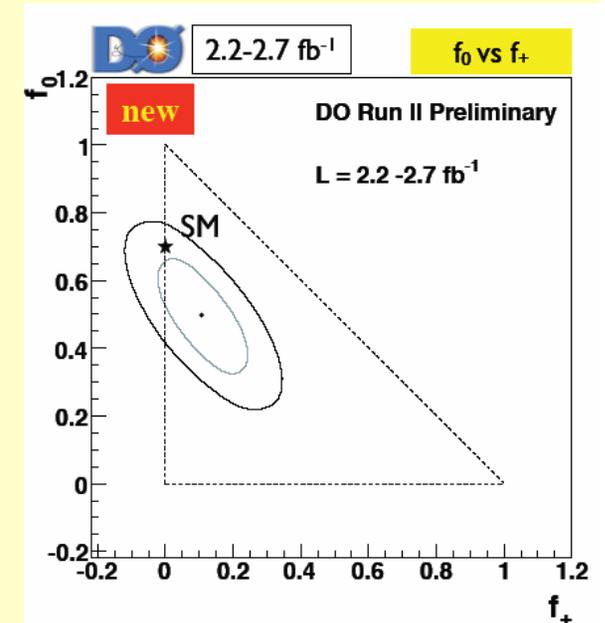
Production Cross-Section
 Production Kinematics
 Spin Polarization
 Production via interm. Resonances
 t' Production



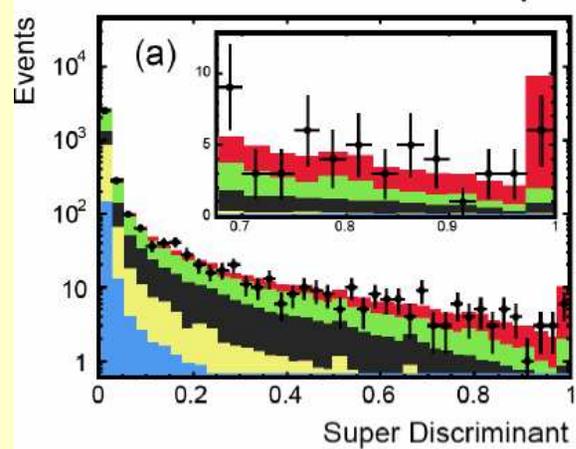
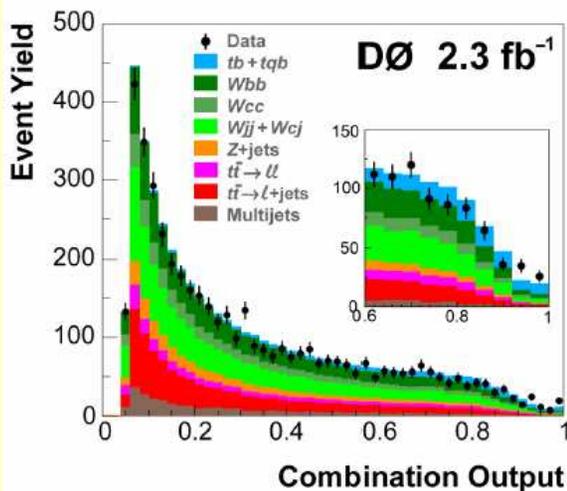
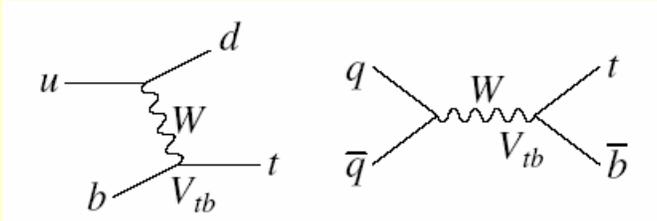
Other top properties



	Tevatron Result	luminosity (fb ⁻¹)
Mass	173.1 ± 1.3 GeV	~ 3.0
W helicity	CDF: f ₀ = 0.66 ± 0.16, f ₊ = -0.03 ± 0.07 DØ: f ₀ = 0.49 ± 0.14, f ₊ = 0.11 ± 0.08	1.9 2.2 – 2.7
Charge	rule out Q = +4/3 (90.% C.L.)	1.5
Lifetime	Γ _t < 13.1 GeV (95% C.L.)	
V _{tb}	V _{tb} > 0.89 (95% C.L.)	~ 1.0
BR(t → Wb) / BR(W → Wq)	R = 0.97 (+0.09) (-0.08)	0.9
BR (t → Zq)	< 3.7% (95% C.L.)	

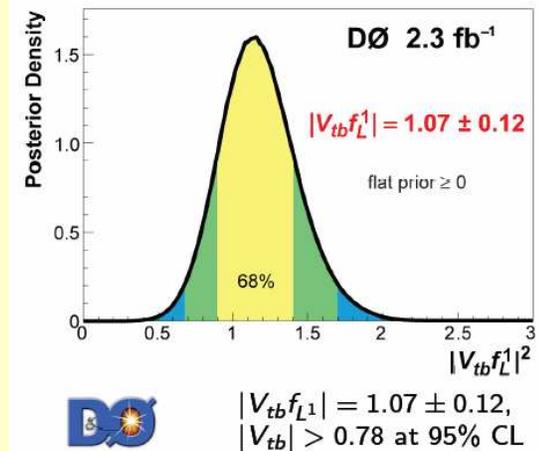
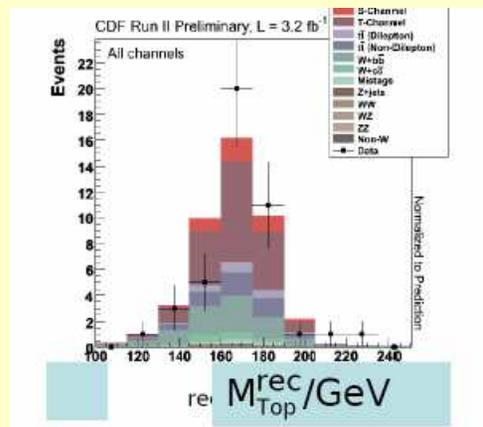
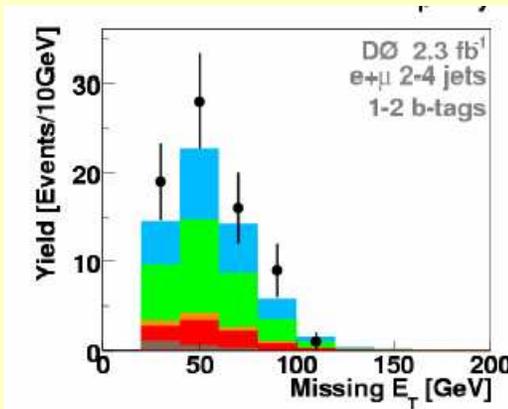


First observation of Single Top Production at the Tevatron



Combined Results

	\mathcal{L} [fb ⁻¹]	Significance Exp.	Obs.	σ_{s+t} [pb]
	2.3	4.5 σ	5.0 σ	3.9 ^{+0.9} _{-0.9}
	3.2	5.9 σ	5.0 σ	2.3 ^{+0.6} _{-0.5}



Summary of the 2. Lecture

- Hadron Colliders Tevatron and LHC play an important role in future tests of the Standard Model
- Predictions of Quantum Chromodynamics can be tested in
 - High p_T jet production
 - W/Z production
 - Top quark production
 -
- In addition, precise measurements of Standard Model parameters can be carried out.

Examples: W mass can be measured to $\sim 10 - 15$ MeV
Top-quark mass to better than ~ 1 GeV