
Fundamentals of Optical Sciences

WS 2015/2016

6. Exercise

23.11.2015

Lecture: Prof. Dr. Alejandro Saenz, Prof. Dr. Oliver Benson

Prepare your answers for the exercise on 30.11.2015.

Exercise 1

Consider a two-level atom and its interaction with the electric component of an electromagnetic field $\vec{E}(t) = E_0 \cos(\omega t) \vec{\varepsilon}$ in the semiclassical description. The full Hamilton operator \hat{H} is

$$\hat{H} = \hat{H}_0 + \hat{H}_1$$

in which \hat{H}_0 is the atomic Hamiltonian and $\hat{H}_1 = -e \hat{\underline{x}} \vec{E}(t)$ describes the Hamiltonian of the radiation field. The atom has two stationary, orthonormal eigenstates of \hat{H}_0 , $|g\rangle$ (“ground”) and $|e\rangle$ (“excited”) with the eigenvalues ω_g and ω_e respectively. Determine in the following the solution of the time-dependent Schrödinger equation

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H} |\psi(t)\rangle.$$

- a) Show that the Hamilton operator can be represented by the following expression

$$\hat{H} = \hbar\omega_g |g\rangle\langle g| + \hbar\omega_e |e\rangle\langle e| + E_0 \cos \omega t \left[d_{g,e} |g\rangle\langle e| + d_{e,g} |e\rangle\langle g| \right]$$

with the dipole moments $d_{i,j} = -e \langle i | \vec{\varepsilon} \cdot \hat{\underline{x}} | j \rangle$.

Hint: The basis $\{|g\rangle, |e\rangle\}$ is complete.

- b) Use the following *ansatz* for the wavefunction of the system

$$\begin{aligned} \Psi(\vec{r}, t) &= c_g(t) \Phi_g(\vec{r}) e^{i\omega_g t} + c_e(t) \Phi_e(\vec{r}) e^{i\omega_e t} \\ |\Psi(t)\rangle &= C_g(t) |g\rangle + C_e(t) |e\rangle. \end{aligned}$$

Show that the equations for the coefficients result in

$$\begin{aligned} i\hbar \frac{d}{dt} C_g(t) &= \hbar\omega_g C_g + E_0 \cos(\omega t) d_{g,e} C_e \\ i\hbar \frac{d}{dt} C_e(t) &= \hbar\omega_g C_e + E_0 \cos(\omega t) d_{e,g} C_g. \end{aligned}$$

- c) Now, introduce the coefficients

$$\begin{aligned} c_e(t) &= C_e(t) e^{i\omega_e t} \\ c_g(t) &= C_g(t) e^{i\omega_g t} \end{aligned}$$

as an *ansatz* to show that

$$\begin{aligned} i\hbar \frac{d}{dt} c_g(t) &= \frac{E_0}{2} d_{g,e} c_e(t) [e^{i(-\omega_{e,g}+\omega)t} + e^{-i(\omega_{e,g}+\omega)t}] \\ i\hbar \frac{d}{dt} c_e(t) &= \frac{E_0}{2} d_{e,g} c_g(t) [e^{i(\omega_{e,g}+\omega)t} + e^{i(\omega_{e,g}-\omega)t}]. \end{aligned}$$

Here: $\omega_{e,g} = \omega_e - \omega_g$.

d) Transform to a rotating coordinate system by introducing the amplitudes

$$\begin{aligned} \tilde{c}_g(t) &= e^{-i\frac{\delta t}{2}} c_g(t) \\ \tilde{c}_e(t) &= e^{+i\frac{\delta t}{2}} c_e(t), \end{aligned}$$

where the *detuning* δ is given by $\delta = \omega - \omega_{e,g}$. The commonly used *rotating-wave approximation* (RWA) neglects the faster oscillating terms with $\omega_{e,g} + \omega$ compared to $\omega_{e,g} - \omega$. Define the *Rabi frequency* $\Omega_0 = \frac{d_{e,g} E_0}{\hbar}$ and use the RWA to show that

$$\frac{d}{dt} \begin{pmatrix} \tilde{c}_g(t) \\ \tilde{c}_e(t) \end{pmatrix} \approx -\frac{i}{2} \begin{pmatrix} \delta & \Omega_0 \\ \Omega_0 & -\delta \end{pmatrix} \begin{pmatrix} \tilde{c}_g(t) \\ \tilde{c}_e(t) \end{pmatrix}.$$

e) Show that the system of equations decouples for the second derivatives,

$$\frac{d^2}{dt^2} \begin{pmatrix} \tilde{c}_g(t) \\ \tilde{c}_e(t) \end{pmatrix} \approx -\frac{1}{4} \begin{pmatrix} \delta^2 + \Omega_0^2 & 0 \\ 0 & \delta^2 + \Omega_0^2 \end{pmatrix} \begin{pmatrix} \tilde{c}_g(t) \\ \tilde{c}_e(t) \end{pmatrix}.$$

f) Show that the solutions of these differential equations are given by

$$\begin{aligned} \tilde{c}_g(t) &\approx -i\frac{\Omega_0}{\Omega} \sin\left(\frac{\Omega t}{2}\right) \tilde{c}_e(0) + \left[\cos\left(\frac{\Omega t}{2}\right) - i\frac{\delta}{\Omega} \sin\left(\frac{\Omega t}{2}\right) \right] \tilde{c}_g(0) \\ \tilde{c}_e(t) &\approx -i\frac{\Omega_0}{\Omega} \sin\left(\frac{\Omega t}{2}\right) \tilde{c}_g(0) + \left[\cos\left(\frac{\Omega t}{2}\right) + i\frac{\delta}{\Omega} \sin\left(\frac{\Omega t}{2}\right) \right] \tilde{c}_e(0) \end{aligned}$$

with the *generalized Rabi frequencies* $\Omega = \sqrt{\delta^2 + \Omega_0^2}$.

g) What are the consequences of the result in f) for the time-dependent population $P_e = |\tilde{c}_e(t)|^2$ of the excited state when the system is at the beginning in its ground state

($c_g(t=0) = 1, c_e(t=0) = 0$) ?

Exercise 2

Consider an atom interacting with an oscillating light field with constant amplitude and frequency ω . The dipole approximation is used for describing the atom-light interaction in the semiclassical treatment. Using perturbation theory the transition probability can be expressed in the form

$$P(t) \propto \frac{\sin^2\left(\frac{1}{2}(\omega_{m,k} - \omega)t\right)}{(\omega_{m,k} - \omega)^2} \quad (1)$$

where $\omega_{m,k} = \omega_m - \omega_k$ is the transition frequency between the states φ_m and φ_k which are the eigenfunctions of the unperturbed Hamiltonian. The behaviour of this function shall be investigated in the following.

- a) Plot function (1) over the detuning $\delta = \omega_{m,k} - \omega$ for (i) $t = 0.5$ s, (ii) $t = 1$ s, (iii) $t = 5$ s, (iv) $t = 10$ s using a plotting program of your choice and explain the behaviour.
- b) For nearly resonant light pulses ($\omega \approx \omega_{m,k}$) the model of an effective two-level atom is appropriate. Why?
- c) Consider the hydrogen atom interacting with a laser pulse with a wavelength of 630 nm. Which transitions in a hydrogen atom can be excited with a 1 fs laser pulse?

Hint: Balmer series!!!